## LECTURE NOTES ON CURVES AND SURFACES IN $\mathbb{R}^{3}$

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Preliminary version and in progress, April 2, 2003

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## 1. Introduction

To study a collection of geometric objects, we often proceed as follows:
(1) Specify a group of "symmetries" that acts on objects, and define two objects $a, b$ to be equivalent if and only if $b=g(a)$ for some $g$ in the symmetry group,
(2) Find a complete set of invariants and their relations. This means that we want to associate to each object $a$ a set of quantities $I(a)$ satisfying certain relations so that
(i) if $a$ and $b$ are equivalent, then $I(a)=I(b)$,
(ii) given a set $\mathcal{I}$ satisfies the relations then there exists an object $a$ unique up to equivalence so that $I(a)=\mathcal{I}$.
One of the simplest examples is the study of triangles in $\mathbb{R}^{3}$ with the group of rigid motions as symmetry group. Recall that a rigid motion is a rotation follows by a translation. It follows from plane geometry that the length of three sides $I(a)=\left\{s_{1}, s_{2}, s_{3}\right\}$ of a triangle $a$ forms a complete set of invariants and they satisfy the relations (triangle inequalities and length is positive):

$$
s_{1}+s_{2}>s_{3}, \quad s_{2}+s_{3}>s_{1}, \quad s_{1}+s_{3}>s_{2}, \quad s_{i}>0, \quad 1 \leq i \leq 3
$$

In this series of lectures, we use differential calculus, linear algebra, and differential equations to study curves and surfaces in $\mathbb{R}^{3}$ with the rigid motions as symmetry group. For example, for

- the study of curves in $\mathbb{R}^{2}$, the objects are immersed curves in $\mathbb{R}^{2}$, and the curvature function of a curve is a complete set of invariants,
- the study of curves in $\mathbb{R}^{3}$, the objects are immersed curves in $\mathbb{R}^{3}$, and the curvature $k$ and torsion $\tau$ of a curve form a complete set of invariants with condition $k>0$,
- the study of surfaces in $\mathbb{R}^{3}$, the objects are immersed surfaces in $\mathbb{R}^{3}$, and their first and second fundamental forms form a complete set of local invariants and they must satisfy the GaussCodazzi equations.
These invariants determine the local geometry of curves and surfaces uniquely up to rigid motions. One goal of these lectures is to explain the Fundamental theorems of curves and surfaces in $\mathbb{R}^{3}$. Another goal is to demonstrate by examples how to construct soliton equations from natural geometric curve flows and from the Gauss-Codazzi equations for surfaces whose numerical invariants satisfy certain natural conditions. The examples we will give are:
(1) the evolution of local invariants of space curves move in the direction of the binormal with curvature as its speed is the non-linear Schrödinger equation,
(2) the Gauss-Codazzi equation for surfaces in $\mathbb{R}^{3}$ with $K=-1$ is the sine-Gordon equation,
(3) the Gauss-Codazzi equation for surfaces in $\mathbb{R}^{3}$ admitting a conformal line of curvature coordinate system is a reduced 3 -wave equation.


## 2. Curves in $\mathbb{R}^{2}$ and $\mathbb{R}^{3}$

Let . denote the standard dot product in $\mathbb{R}^{n}$, and $\|v\|=\sqrt{v \cdot v}$ the length of $v \in \mathbb{R}^{n}$. In this section, a curve in $\mathbb{R}^{n}$ means a smooth map $\alpha:[a, b] \rightarrow \mathbb{R}^{n}$ so that $\alpha^{\prime}(t) \neq 0$ for all $t \in[a, b]$, i.e., $\alpha$ is an immersion. The arc length of $\alpha$ from $t_{0} \in[a, b]$ to $t$ is

$$
s(t)=\int_{t_{0}}^{t}\left\|\alpha^{\prime}(t)\right\| d t
$$

Since $\frac{d s}{d t}=\left\|\alpha^{\prime}\right\|$ and $\alpha^{\prime}(t) \neq 0$ for all $t \in[a, b]$, the function $s:[a, b] \rightarrow$ $[0, \ell]$ is a diffeomorphism, where $\ell$ is the arc length of $\alpha$. So we can make a change of coordinate $t=t(s)$ (the inverse function of $s=s(t)$ ). Moreover, by the chain rule we have

$$
\frac{d \alpha}{d s}=\frac{\frac{d \alpha}{d t}}{\frac{d t}{d s}}=\frac{\alpha^{\prime}(t)}{\left\|\alpha^{\prime}(t)\right\|}
$$

i.e., $\frac{d \alpha}{d s}$ is a unit vector tangent to $\alpha$.

A curve $\alpha:\left[c_{1}, c_{2}\right] \rightarrow \mathbb{R}^{n}$ is said to be parametrized by its arc length if $\left\|\alpha^{\prime}(t)\right\|=1$ for all $t \in\left[c_{1}, c_{2}\right]$. The above discussion says that any curve in $\mathbb{R}^{n}$ can be parametrized by its arc length.

### 2.1. Some results from ODE.

ca Proposition 2.1.1. Suppose $e_{1}, \cdots, e_{n}:[a, b] \rightarrow \mathbb{R}^{n}$ are smooth maps, and $\left\{e_{1}(t), \cdots, e_{n}(t)\right\}$ is an orthonormal (o.n.) basis of $\mathbb{R}^{n}$ for all $t \in[a, b]$. Then
(1) there exists a matrix valued function $A(t)=\left(a_{i j}(t)\right)$ so that

$$
\begin{equation*}
e_{i}^{\prime}(t)=\sum_{j=1}^{n} a_{j i}(t) e_{j}(t) \tag{2.1.1}
\end{equation*}
$$

(2) $a_{i j}=e_{j}^{\prime} \cdot e_{i}$.
(3) $a_{i j}+a_{j i}=0$, i.e., $A(t)$ is skew-symmetric

Proof. Since $e_{1}, \cdots, e_{n}$ form an o.n. basis of $\mathbb{R}^{n}, e_{i}^{\prime}(t)$ can be written as a linear combination of $e_{1}(t), \cdots, e_{n}(t)$. This proves (1). Dot (2.1.1) with $e_{j}$ to get (2). Because $e_{i} \cdot e_{j}=\delta_{i j}$, the derivative of this equation implies that $e_{i}^{\prime} \cdot e_{j}+e_{i} \cdot e_{j}^{\prime}=0$, which proves (3).

It follows from the definition of matrix multiplication that

$$
\left(v_{1}, v_{2}, \cdots, v_{n}\right)\left(\begin{array}{c}
x_{1} \\
x_{2} \\
\cdot \\
\cdot \\
x_{m}
\end{array}\right)=\sum_{k} x_{k} v_{k}
$$

where $v_{1}, \cdots, v_{n}$ are column vectors. So we can rewrite ( $(2.1 .1)$ in matrix form

$$
g^{\prime}(t)=g(t) A(t), \quad \text { where } g(t)=\left(e_{1}(t), \cdots, e_{n}(t)\right), \quad A(t)=\left(a_{i j}(t)\right)
$$

Note the $i$-th column of $A$ gives the coefficient of of $e_{i}^{\prime}$ with respect to the o.n. basis $e_{1}, \cdots, e_{n}$.

The following is a special case of the existence and uniqueness theorem for solutions of ODE:
ah Theorem 2.1.2. (Existence and Uniqueness of ODE) Let $I$ be an interval of $\mathbb{R}, U$ an open subset of $\mathbb{R}^{n}$, and $F: I \times U \rightarrow \mathbb{R}^{n}$ a smooth map. Then given $t_{0} \in I$ and $p_{0} \in U$, there exists $\delta>0$ and a unique smooth $x:\left(t_{0}-\delta, t_{0}+\delta\right) \rightarrow U$ so that $x$ is a solution of the following initial value problem:

$$
\left\{\begin{array}{l}
x^{\prime}(t)=F(t, x(t)),  \tag{2.1.2}\\
x\left(t_{0}\right)=p_{0}
\end{array}\right.
$$

Let $g l(n)$ denote the space of all $n \times n$ real matrices. An $n \times n$ matrix $C$ is called orthogonal if $C^{t} C=\mathrm{I}$. It follows from definition of matrix product that:

Proposition 2.1.3. The following statements are equivalent for a matrix $C \in \operatorname{gl}(n)$ :
(1) $C$ is orthogonal,
(2) the column vectors of $C$ are o.n.,
(3) the row vectors of $C$ are o.n..
as Proposition 2.1.4. Suppose $A(t)=\left(a_{i j}(t)\right)$ is skew symmetric for all $t$. If $C \in \operatorname{gl}(n)$ is an orthogonal matrix, $t_{0} \in(a, b)$, and $g:(a, b) \rightarrow$ $g l(n)$ is the solution of

$$
g^{\prime}(t)=g(t) A(t), \quad g\left(t_{0}\right)=C,
$$

then $g(t)$ is orthogonal for all $t \in(a, b)$.
Proof. Set $y=g^{t} g$. We compute the derivative $y^{\prime}$ to get

$$
y^{\prime}=\left(g^{t}\right)^{\prime} g+g^{t} g^{\prime}=\left(g^{\prime}\right)^{t} g+g^{t} g^{\prime}=(g A)^{t} g+g^{t}(g A)=A^{t} y+y A
$$

Since the constant function $y(t)=I$ is a solution with initial data $y_{h}\left(t_{0}\right)=\mathrm{I}$, it follows from the uniquness of solutions of ODE (Theorem 2.1.2) that $y(t)=\mathrm{I}$ for all $t \in(a, b)$.

## Exercise 2.1.5.

(1) For $X, Y \in \operatorname{gl}(n)$, let $(X, Y)=\operatorname{tr}\left(X^{t} Y\right)$. Show that $($,$) is an$ inner product on $g l(n)$, and $g l(n)$ is isometric to the Euclidean space $\mathbb{R}^{n^{2}}$ with the standard inner product.
(2) Given $A \in g l(n)$, prove that

$$
y(t)=\sum_{j=0}^{\infty} \frac{A^{n}}{n!} t^{n}
$$

converges uniformly on any finite interval $[-r, r]$. Hence $y(t)$ is continuous. Prove also that $y$ is differentiable and $y$ is the solution of the initial value problem:

$$
y^{\prime}(t)=A y(t), \quad y(0)=\mathrm{Id} .
$$

We will denote $y(t)$ as $e^{t A}$.
(3) Compute $e^{t A}$ for $A=\operatorname{diag}\left(a_{1}, \ldots, a_{n}\right)$ and for $A=\left(\begin{array}{cc}0 & 1 \\ -1 & 0\end{array}\right)$.
(4) Suppose $D$ is the diagonal matrix $\operatorname{diag}\left(d_{1}, \cdots, d_{n}\right), g$ is an invertible $n \times n$ matrix, and $A=g D g^{-1}$. Show that $e^{t A}=$ $g e^{t D} g^{-1}$.

### 2.2. Plane curves.

Let $\alpha(t)=(x(t), y(t))$ be a smooth curve parametrized by its arc length. Then $e_{1}(t)=\alpha^{\prime}(t)=\left(x^{\prime}(t), y^{\prime}(t)\right)$ is a unit vector tangent to $\alpha$ at $\alpha(t)$. Note that $\left(-x_{2}, x_{1}\right)$ is perpendicular to $\left(x_{1}, x_{2}\right)$. Let $e_{2}(t)=\left(-y^{\prime}(t), x^{\prime}(t)\right)$. Then $e_{1}(t), e_{2}(t)$ form an o.n. basis of $\mathbb{R}^{2}$. Such $\left(e_{1}(t), e_{2}(t)\right)$ is called a o.n. moving frame along $\alpha$. Since a $2 \times 2$ skew-symmetric matrix is of the form

$$
\left(\begin{array}{cc}
0 & -u \\
u & 0
\end{array}\right)
$$

It follows from Proposition ${ }_{2}^{\text {Ca }} 1.1 .1$ that there exists a smooth function $k(t)$ so that

$$
\left(e_{1}, e_{2}\right)^{\prime}=\left(e_{1}, e_{2}\right)\left(\begin{array}{cc}
0 & -k \\
k & 0
\end{array}\right)
$$

In fact, we have

$$
k=e_{1}^{\prime} \cdot e_{2},
$$

and

$$
\left\{\begin{array}{l}
e_{1}^{\prime}=k e_{2}, \\
e_{2}^{\prime}=-k e_{1}
\end{array}\right.
$$

The function $k$ is called the curvature of the curve $\alpha$.
Example 2.2.1. A straight line can be parametrized as $\alpha(t)=p_{0}+t u$, where $p_{0} \in \mathbb{R}^{2}$ and $u$ a constant unit vector in $\mathbb{R}^{2}$. Then $e_{1}=u$ a constant, and $e_{1}^{\prime}=0$. Hence the curvature of a straight line is identically 0 . Conversely, if the curvature of a curve $\beta$ parametrized by arc length is identially zero, then $\beta$ must be a straight line. To see this, note first that $e_{1}^{\prime}=\beta^{\prime \prime}=k e_{2}=0$. Hence $e_{1}$ is a constant vector, say $u$. But $(\beta(t)-t u)^{\prime}=e_{1}-u=0$ implies that $\beta-t u$ is a constant $p_{0}$. So $\beta(t)=p_{0}+t u$ is a straight line. So the curvature of a plane curve measures the deviation of the curve being a straight line.

Example 2.2.2. The circle of radius $r$ centered at $p_{0}$ is

$$
\alpha(t)=p_{0}+\left(r \cos \frac{t}{r}, r \sin \frac{t}{r}\right) .
$$

Since $\alpha^{\prime}(t)=\left(-\sin \frac{t}{r}, \cos \frac{t}{r}\right)$ has length $1, t$ is an arc length parameter. A direct computation shows that the curvature is the constant function $\frac{1}{r}$.

The following Proposition says that the curvature is the instantaneous rate of change of the polar angle of the tangent of the curve. So when $k$ is large, the curve winds around faster.

Proposition 2.2.3. Suppose $\alpha(t)=(x(t), y(t))$ is parametrized by its arc length, and $\theta(t)$ the angle from $(1,0)$ to $\alpha^{\prime}(t)$. Then $\theta^{\prime}(t)=k(t)$.
Proof. Since $e_{1}=\alpha^{\prime}=(\cos \theta(t), \sin \theta(t)), e_{2}(t)=(-\sin \theta(t), \cos \theta(t))$. A direct computation implies that $k=e_{1}^{\prime} \cdot e_{2}=\theta^{\prime}(t)$. Moreover,

$$
\left\{\begin{array}{l}
\alpha^{\prime}=(\cos \theta, \sin \theta)  \tag{2.2.1}\\
\theta^{\prime}=k
\end{array}\right.
$$

Note that the tangent line to $\alpha$ at $\alpha\left(t_{0}\right)$ gives the best linear approximation to $\alpha$ near $\alpha\left(t_{0}\right)$. The best quadratic approximation of $\alpha$ is the osculating circle. We explain this next.

Definition 2.2.4. The osculating circle at $\alpha\left(t_{0}\right)$ of a curve $\alpha$ is the circle of radius $r=\frac{1}{k\left(t_{0}\right)}$ centered at $\alpha\left(t_{0}\right)+r e_{2}\left(t_{0}\right)$.
Proposition 2.2.5. Suppose $\alpha(t)=(x(t), y(t))$ is a plane curve, and $t$ is the arc length parameter. Then the osculating circle $C$ at $\alpha\left(t_{0}\right)$ is the best second order approximation of $\alpha$ near $t_{0}$.

Proof. We may assume that $t_{0}=0, \alpha(0)=(0,0), \alpha^{\prime}(0)=(1,0)$, and $e_{2}(0)=(0,1)$. Let $r=\frac{1}{k(0)}$. Since $e_{1}^{\prime}=k e_{2}, \alpha^{\prime \prime}(0)=\frac{1}{r}(0,1)$. So the degree 2 Taylor series of $\alpha(t)=(x(t), y(t))$ at $t=0$ is

$$
\left(x(0)+x^{\prime}(0) t+\frac{1}{2} x^{\prime \prime}(0) t^{2}, y(0)+y^{\prime}(0) t+\frac{1}{2} y^{\prime \prime}(0) t^{2}\right)=\left(t, \frac{t^{2}}{2 r}\right) .
$$

The osculating circle at $\alpha(0)=(0,0)$ is the circle of radius $r$ centered at $(0,0)+r(0,1)=(0, r)$. So

$$
C(t)=(0, r)+r\left(\sin \frac{t}{r},-\cos \frac{t}{r}\right)
$$

gives an arc length parameter of the osculating circle at $\alpha(0)$. It is easy to check that the degree two Taylor series of $C(t)$ at $t=0$ is the same as the one for $\alpha(t)$.

## Theorem 2.2.6. (Fundamental Theorem of Curves in $\mathbb{R}^{2}$ )

(1) Given a smooth function $k:(a, b) \rightarrow \mathbb{R}, p_{0} \in \mathbb{R}^{2}, t_{0} \in(a, b)$, and $u_{1}, u_{2}$ a fixed o.n. basis of $\mathbb{R}^{2}$, there exists a unique curve $\alpha(t)$ defined in a small neighborhood of $t_{0}$ so that $t$ is an arc length parameter, $\alpha(0)=p_{0}, \alpha^{\prime}(0)=u_{1}$, and $k(t)$ is its curvature function.
(2) Let $\alpha, \tilde{\alpha}:(-\delta, \delta) \rightarrow \mathbb{R}^{2}$ be smooth curves parametrized by arc length. If $\alpha$ and $\tilde{\alpha}$ have the same curvature function, then there exists a rigid motion $T$ so that $\tilde{\alpha}(t)=T(\alpha(t))$ for all $t$.

Proof. (1) To construct the curve $\alpha$, we need to solve

$$
\left\{\begin{array}{l}
\alpha^{\prime}=e_{1},  \tag{2.2.2}\\
e_{1}^{\prime}=k e_{2}, \\
e_{2}^{\prime}=-k e_{1}
\end{array}\right.
$$

with initial data $\alpha(0)=p_{0}, e_{1}(0)=u_{1}, e_{2}(0)=u_{2}$. Let $F:(a, b) \times$ $\left(\mathbb{R}^{2}\right)^{3} \rightarrow\left(\mathbb{R}^{2}\right)^{3}$ be the smooth map defined by

$$
F\left(t, y_{0}, y_{1}, y_{2}\right)=\left(y_{1}, k(t) y_{2},-k(t) y_{3}\right) .
$$

It follows from Theorem ${ }^{\text {ah }} 2.1 .2$ that there exists a unique solution with initial data $\left(y_{0}, y_{1}, y_{2}\right)(0)=\left(p_{0}, u_{1}, u_{2}\right)$. Since $\left(\begin{array}{cc}0 & -k \\ k & 0\end{array}\right)$ is skew-symmetric and

$$
\left(y_{1}, y_{2}\right)^{\prime}=\left(y_{1}, y_{2}\right)\left(\begin{array}{cc}
0 & -k \\
k & 0
\end{array}\right)
$$

it follows from Proposition ${ }_{22.1 .4}$ that $y_{1}(t), y_{2}(t)$ form an o.n. basis for each $t$.
(2) Let $e_{1}=\alpha^{\prime}, e_{2}$ the unit normal for $\alpha$ so that $\operatorname{det}\left(e_{1}, e_{2}\right)=1$, and $\tilde{e}_{1}=\tilde{\alpha}^{\prime}, \tilde{e}_{2}$ the unit normal for $\tilde{\alpha}$ so that $\operatorname{det}\left(\tilde{e}_{1}, \tilde{e}_{2}\right)=1$. By assumption $\alpha$ and $\tilde{\alpha}$ have the same curvature function $k$, so both $\left(\alpha, e_{1}, e_{2}\right)$ and $\left(\tilde{\alpha}, \tilde{e}_{1}, \tilde{e}_{2}\right)$ satisfy the same differential equation

$$
\begin{equation*}
y_{0}^{\prime}=y_{1}, \quad y_{1}^{\prime}=k y_{2}, \quad y_{2}^{\prime}=-k y_{1} . \tag{2.2.3}
\end{equation*}
$$

at
Let $h$ denote the constant rotation

$$
h=\left(\tilde{e}_{1}(0), \tilde{e}_{2}(0)\right)\left(e_{1}(0), e_{2}(0)\right)^{-1}
$$

and $T$ the rigid motion

$$
T(v)=\tilde{\alpha}(0)-h(\alpha(0))+h v .
$$

Set $\beta(t)=T(\alpha(t))$. Then $\beta(0)=\tilde{\alpha}(0)$ and $\beta^{\prime}=h e_{1}$. Since $h$ is a rotation, $h e_{2}$ is the unit normal of $\beta$ and $\left(h e_{1}, h e_{2}\right)$ is a local o.n. frame along $\beta$. By definition of $h$, we have $h\left(e_{i}(0)\right)=\tilde{e}_{i}(0)$ for $i=1,2$. But

$$
\left\{\begin{array}{l}
\beta^{\prime}=h e_{1}, \\
\left(h e_{1}\right)^{\prime}=h\left(e_{1}^{\prime}\right)=h\left(k e_{2}\right)=k h\left(e_{2}\right), \\
\left(h\left(e_{2}\right)^{\prime}=h\left(e_{2}^{\prime}\right)=h\left(-k e_{1}\right)=-k h\left(e_{1}\right)\right.
\end{array}\right.
$$

So both $\tilde{a}$ and $\beta$ are solutions of the following initial value problem:

$$
\left\{\begin{array}{l}
y_{0}^{\prime}=y_{1} \\
y_{1}^{\prime}=k y_{2} \\
y_{2}^{\prime}=-k y_{1} \\
\left(y_{0}, y_{1}, y_{2}\right)(0)=\left(\tilde{\alpha}(0), \tilde{e}_{1}(0), \tilde{e}_{2}(0)\right)
\end{array}\right.
$$

By the uniqueness of ODE, $\beta=\tilde{\alpha}$, i.e., $\tilde{\alpha}=T(\alpha)$.
Remark. We can prove Theorem ${ }_{2}^{\text {av }} 2.2 .6$ by solve equation $\left(\frac{\mathrm{fh}}{2.2} \cdot 1\right)$. ff Note given two initial conditions $\alpha(0)=p_{0}$ and $\theta(0)=\theta_{0}$, equation (2.2.1) can be solved by integration. However, the benefit of the ODE proof given above is that the same method also works for the space curves.

Let $\alpha:[0, \ell] \rightarrow \mathbb{R}^{2}$ be a smooth curve parametrized by its arc length, and $e_{1}=\alpha^{\prime}$ and $e_{2}$ the unit normal. Below we give three classical ways to generate new curves from $\alpha$ :
(1) The involute of $\alpha$ starting at $\alpha\left(t_{0}\right)$ is the curve

$$
\beta_{t_{0}}(t)=\alpha\left(t+t_{0}\right)-t e_{1}\left(t+t_{0}\right) .
$$

Note that if $\alpha$ is a simple close curve, then the involute of $\alpha$ at $\alpha\left(t_{0}\right)$ can be viewed as pulling the thread from a spool of the shape of $\alpha$ from the point $\alpha\left(t_{0}\right)$.
(2) The evolute of a curve $\alpha$ is the locus of the centers of osculating circles of $\alpha$, i.e., the evolute is

$$
\beta(t)=\alpha(t)+\frac{1}{k(t)} e_{2}(t) .
$$

(3) The curve parallel to $\alpha$ with distance $r$ is

$$
\alpha(t)+r e_{2}(t)
$$

## Exercise 2.2.7.

(1) Recall that a map $T: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ is a rigid motion if $T$ is of the form $T(x)=g x+b$ for some $n \times n$ orthogonal matrix $g$ and $b \in \mathbb{R}^{n}$. Prove that the set of all rigid motions of $\mathbb{R}^{n}$ with composition as multiplication is a group.
(2) Let $\alpha$ be a curve in $\mathbb{R}^{2}$ parametrized by arc length, $k$ its curvature, and $T$ a rigid motion. Prove that the curvature of $\beta=T \circ \alpha$ is also $k$.
(3) Suppose $\alpha(t)$ is not necessarily parametrized by arc length. Write down the curvature as a function of $t$.
(4) Prove that the curve parallel to $\alpha$ with distance $r$ fails to be an immersed curve at $t_{0}$ if and only if the radius of osculating circle at $\alpha\left(t_{0}\right)$ is $\frac{1}{r}$. Use 3D-XplorMath to visualize this.
(5) Use the 3D-XplorMath Plane curve category to see that the involutes of a curve seem to be all parallel. Use mathematics to either prove or disprove what you see.
(6) Write down the curve obtained by tracing a fixed point on a circle of radius $a$ when the circle is rolling along the $x$-axis with a fixed speed $b$. This curve is called a cycloid.
(7) Show that the evolute of the cycloid is a translation of the same cycloid. Check this visually using the 3D-XplorMath.

### 2.3. Space curves.

Let $\alpha(t)$ be a curve in $\mathbb{R}^{3}$ parametrized by arc length. As for plane curves, we need to construct a moving o.n. frame $e_{1}(t), e_{2}(t), e_{3}(t)$ along $\alpha$ such that $e_{1}(t)=\alpha^{\prime}(t)$. Since the normal space at every point is two dimension, it is not clear which normal direction should be $e_{2}$. There are two natural ways to choose $e_{2}$. The first is the classical Frenet frame. Since $e_{1}=\alpha^{\prime}$ has unit length,

$$
\left(e_{1} \cdot e_{1}\right)^{\prime}=2 e_{1}^{\prime} \cdot e_{1}=0
$$

So $\alpha^{\prime \prime}=e_{1}^{\prime}$ is perpendicular to $e_{1}$. Assume that $\alpha^{\prime \prime}(t)$ never vanishes in an interval $(a, b)$. Let $e_{2}(t)$ denote the unit direction along $\alpha^{\prime \prime}(t)$ and $k(t)$ the length of $\alpha^{\prime \prime}(t)$, i.e., $e_{1}^{\prime}=k e_{2}$. Let $e_{3} \overline{\bar{c}} e_{1} \times e_{2}$. So $\left\{e_{1}, e_{2}, e_{3}\right\}$ is an o.n. basis of $\mathbb{R}^{3}$ for all $t$. By Proposition 2.1.1, there exists a skew symmetric $3 \times 3$ matrix valued map $A(t)$ so that

$$
\left(e_{1}, e_{2}, e_{3}\right)^{\prime}=\left(e_{1}, e_{2}, e_{3}\right) A
$$

Since $3 \times 3$ skew-symmetric matrix must be of the form

$$
\left(\begin{array}{ccc}
0 & -f_{1} & -f_{2} \\
f_{1} & 0 & -f_{3} \\
f_{2} & f_{3} & 0
\end{array}\right)
$$

and $e_{1}^{\prime}=k e_{2}$, the first column of $A$ must be $(0, k, 0)^{t}$ and there exists $\tau$ so that

$$
\left(e_{1}, e_{2}, e_{3}\right)^{\prime}=\left(e_{1}, e_{2}, e_{3}\right)\left(\begin{array}{ccc}
0 & -k & 0  \tag{2.3.1}\\
k & 0 & -\tau \\
0 & \tau & 0
\end{array}\right)
$$

Note that $\tau=e_{2}^{\prime} \cdot e_{3}$. The functions $k$ and $\tau$ are called the curvature and torsion respectively, $e_{2}$ the normal, $e_{3}$ the binormal, and $\left(e_{1}, e_{2}, e_{3}\right)$ is
the Frenet frame along $\alpha$. Equation (2.3.1) is called the Frenet-Serret equation, which can also be written as

$$
\left\{\begin{array}{l}
e_{1}^{\prime}=k e_{2}  \tag{2.3.2}\\
e_{2}^{\prime}=-k e_{1}+\tau e_{2} \\
e_{3}^{\prime}=-\tau e_{2}
\end{array}\right.
$$

Proposition 2.3.1. If the torsion of $\alpha$ is identically zero, then $\alpha$ must lie in a plane.
Proof. Since $\tau=0$, the third equation of ( ${ }^{\mathrm{ek}} 3.2$ ) implies that $e_{3}^{\prime}=$ $-\tau e_{2}=0$. Hence $e_{3}=b$ is a constant vector. The derivative of $\alpha(t) \cdot b$ is $e_{1} \cdot b=e_{1} \cdot e_{3}=0$. So $\alpha(t) \cdot b=c_{0}$ a constant.

So the torsion of a space curve measures the deviation of $\alpha$ from being a plane curve.

## Theorem 2.3.2. (Fundamental Theorem of curves in $\mathbb{R}^{3}$ )

(1) Given smooth functions $k, \tau:(a, b) \rightarrow \mathbb{R}$ so that $k(t)>0$, $t_{0} \in(a, b), p_{0} \in \mathbb{R}^{3}$ and $\left(u_{1}, u_{2}, u_{3}\right)$ a fixed o.n. basis of $\mathbb{R}^{3}$, there exists $\delta>0$ and a unique curve $\alpha:\left(t_{0}-\delta, t_{0}+\delta\right) \rightarrow \mathbb{R}^{3}$ parametrized by arc length such that $\alpha(0)=p_{0}$, and $\left(u_{1}, u_{2}, u_{3}\right)$ is the Frenet frame of $\alpha$ at $t=t_{0}$.
(2) Suppose $\alpha, \tilde{\alpha}:(a, b) \rightarrow \mathbb{R}^{3}$ are curves parametrized by arc length, and $\alpha, \tilde{\alpha}$ have the same curvature function $k$ and torsion function $\tau$. Then there exists a rigid motion $T$ so that $\tilde{\alpha}=T(\alpha)$.

Proof. This theorem can be proved in exactly the same way as Theorem 2.2 .6 by applying the existence and uniqueness Theorem 2.1 .2 of solution of ODE to $F:(a, b) \times\left(\mathbb{R}^{3}\right)^{3} \rightarrow\left(\mathbb{R}^{3}\right)^{3}$ defined by

$$
F\left(t, y_{0}, y_{1}, y_{2}\right)=\left(y_{1}, k(t) y_{2},-k(t) y_{1}+\tau(t) y_{3},-\tau(t) y_{2}\right) .
$$

Next we describe parallel frames. Let $\left(e_{1}, \xi_{1}, \xi_{2}\right)$ be a smooth o.n. frame along $\alpha$ such that $e_{1}=\alpha^{\prime}$. Then by Proposition [2.1.1, there exist three smooth functions $f_{1}, f_{2}, f_{3}$ so that

$$
\left(e_{1}, \xi_{1}, \xi_{2}\right)^{\prime}=\left(e_{1}, \xi_{1}, \xi_{2}\right)\left(\begin{array}{ccc}
0 & -f_{1} & -f_{2} \\
f_{1} & 0 & -f_{3} \\
f_{2} & f_{3} & 0
\end{array}\right)
$$

We want to change the o.n. normal frame $\left(\xi_{1}, \xi_{2}\right)$ to $\left(v_{1}, v_{2}\right)$ so that the 23 entry of the coefficient matrix of $\left(e_{1}, v_{1}, v_{2}\right)^{\prime}$ is zero. To do this, we
rotate the normal frame $\xi_{1}(t), \xi_{2}(t)$ by an angle $\beta(t)$ :

$$
\left\{\begin{array}{l}
v_{1}(t)=\cos \beta(t) \xi_{1}+\sin \beta(t) \xi_{2} \\
v_{2}(t)=-\sin \beta(t) \xi_{1}+\cos \beta(t) \xi_{2}
\end{array}\right.
$$

We want to choose $\beta$ so that $v_{1}^{\prime} \cdot v_{2}=0$. But

$$
\begin{aligned}
v_{1}^{\prime} \cdot v_{2}= & \beta^{\prime}\left(-\sin \beta \xi_{1}+\cos \beta \xi_{2}\right) \cdot\left(-\sin \beta \xi_{1}+\cos \beta \xi_{2}\right) \\
& +\left(\cos \beta \xi_{1}^{\prime}+\sin \beta \xi_{2}^{\prime}\right) \cdot\left(-\sin \beta \xi_{1}+\cos \beta \xi_{2}\right) \\
= & \beta^{\prime}+\left(\cos ^{2} \beta \xi_{1}^{\prime} \cdot \xi_{2}-\sin ^{2} \beta \xi_{2}^{\prime} \cdot \xi_{1}\right) \\
= & \beta^{\prime}+f_{3} .
\end{aligned}
$$

So if we choose $\beta(t)$ so that

$$
\begin{equation*}
\beta^{\prime}=-f_{3}, \tag{2.3.3}
\end{equation*}
$$

cb
then the new o.n. frame $\left(e_{1}, v_{1}, v_{2}\right)$ satisfies

$$
\left(e_{1}, v_{1}, v_{2}\right)^{\prime}=\left(e_{1}, v_{1}, v_{2}\right)\left(\begin{array}{ccc}
0 & -k_{1} & -k_{2}  \tag{2.3.4}\\
k_{1} & 0 & 0 \\
k_{2} & 0 & 0
\end{array}\right)
$$

where

$$
\left\{\begin{array}{l}
k_{1}=e_{1}^{\prime} \cdot v_{1}=e_{1}^{\prime} \cdot\left(\cos \beta \xi_{1}+\sin \beta \xi_{2}\right)=f_{1} \cos \beta+f_{2} \sin \beta  \tag{2.3.5}\\
k_{2}=e_{1}^{\prime} \cdot v_{2}=-f_{1} \sin \beta+f_{2} \cos \beta
\end{array}\right.
$$

The o.n. frame $\left(e_{1}, v_{1}, v_{2}\right)$ is called a parallel frame along $\alpha$. The reason for the name parallel is because the normal component of the derivative of the normal vector field $v_{i}$ is zero. The functions $k_{1}, k_{2}$ are called the prinicipal curvatures along $v_{1}, v_{2}$ respectively. However, the choice of parallel frame is not unique because $\beta$ only need to satisfy ( 2.3 .3 ), so we can replace $\beta$ by $\beta$ plus a constant $\theta_{0}$. If we rotate $\left(v_{1}, v_{2}\right)$ by a constant rotation $\theta_{0}$ to ( $\tilde{v}_{1}, \tilde{v}_{2}$ ), then $\left(e_{1}, \tilde{v}_{1}, \tilde{v}_{2}\right)$ is again parallel. But the principal curvature $\tilde{k}_{1}$ along $\tilde{v}_{1}$ is

$$
\tilde{k}_{1}=e_{1}^{\prime} \cdot \tilde{v}_{1}=e_{1}^{\prime} \cdot\left(\cos \theta_{0} v_{1}+\sin \theta_{0} v_{2}\right)=\cos \theta_{0} k_{1}+\sin \theta_{0} k_{2} .
$$

Similar computation gives

$$
\tilde{k}_{2}=-\sin \theta_{0} k_{1}+\cos \theta_{0} k_{2}
$$

Note that

$$
\alpha^{\prime \prime}=e_{1}^{\prime}=k_{1} v_{1}+k_{2} v_{2}=\tilde{k}_{1} \tilde{v}_{1}+\tilde{k}_{2} \tilde{v}_{2} .
$$

So we have
cf Proposition 2.3.3. Let $\alpha$ be a curve in $\mathbb{R}^{3}$ parametrized by arc length, $\left(v_{1}, v_{2}\right)$ a parallel normal frame along $\alpha$, and $k_{1}, k_{2}$ principal curvature with respect to $v_{1}, v_{2}$ respectively. Then $\alpha^{\prime \prime}=k_{1} v_{1}+k_{2} v_{2}$.
el Corollary 2.3.4. Let $\left(v_{1}, v_{2}\right)$ be a parallel normal frame for $\alpha, k_{1}, k_{2}$ the corresponding principal curvatures, $\left(e_{1}, e_{2}, e_{3}\right)$ the Frenet frame, and $k, \tau$ the curvature and torsion of $\alpha$. Then
(1) there exists $\beta(t)$ so that $v_{1}=\cos \beta e_{2}+\sin \beta e_{3}, v_{2}=-\sin \beta e_{2}+$ $\cos \beta e_{3}$, and $\beta^{\prime}=-\tau$,
(2) $k_{1}=k \cos \beta, k_{2}=-k \sin \beta$.

Proof. Use $\xi_{1}=e_{2}, \xi_{2}=e_{3}$ in the computation before Proposition $\frac{\text { cf }}{\text { L. }}$ 2.3.3. Since $f_{1}=k, f_{2}=0$ and $f_{3}=\tau$, equations (2.3.3) and (2.3.5) imply the Corollary.

## Exercise 2.3.5.

(1) Find all curves in $\mathbb{R}^{3}$ with constant curvature and torsion.
(2) Find all curves in $\mathbb{R}^{3}$ with constant principal curvatures.
(3) Let $\alpha$ be a curve in $\mathbb{R}^{n}$ parametrized by arc length. Prove that for generic set of curves $\alpha$, there exists an o.n. moving frame $\left(e_{1}, \cdots, e_{n}\right)$ such that

$$
\left(e_{1}, \cdots, e_{n}\right)^{\prime}=\left(e_{1}, \cdots, e_{n}\right) P
$$

where

$$
P=\left(\begin{array}{cccccc}
0 & -k_{1} & 0 & \cdots & & \\
k_{1} & 0 & -k_{2} & & & \\
0 & k_{2} & 0 & -k_{3} & & \\
& & & & 0 & -k_{n-1} \\
0 & & & & k_{n-1} & 0
\end{array}\right)
$$

for some smooth functions $k_{1}, \cdots, k_{n-1}$. (So the coefficient matrix for $\left(e_{1}, \cdots, e_{n}\right)^{\prime}$ is tri-diagonal). This is an analogue of the Frenet frame for curves in $\mathbb{R}^{n}$.
(4) Formulate and prove the fundamental theorem for curves in $\mathbb{R}^{n}$ with the frame given in the above exercise.
(5) Let $\alpha$ be a curve in $\mathbb{R}^{n}$ parametrized by arc length.
(i) Prove that there exists a o.n. moving frame $\left(e_{1}, v_{1}, \cdots, v_{n-1}\right)$ such that

$$
\left(e_{1}, v_{1}, \cdots, v_{n-1}\right)^{\prime}=\left(e_{1}, v_{1}, \cdots, v_{n-1}\right) P
$$

where

$$
P=\left(\begin{array}{ccccc}
0 & -k_{1} & -k_{2} & \cdots & -k_{n-1} \\
k_{1} & 0 & \cdots & & \\
k_{2} & 0 & \cdots & & \\
\cdot & & & & \\
k_{n-1} & 0 & & &
\end{array}\right)
$$

for some smooth function $k_{1}, \cdots, k_{n-1}$. This is the parallel frame for curves in $\mathbb{R}^{n}$.
(ii) Formulate and prove the fundamental theorem for curves in $\mathbb{R}^{n}$ with the parallel frame.
(iii) Find all curves in $\mathbb{R}^{n}$ with constant principal curvatures.

## xj

### 2.4. The Frobenius Theorem.

Let $\mathcal{O}$ be an open subset of $\mathbb{R}^{2}$, and $A, B ; \mathcal{O} \times \mathbb{R} \rightarrow \mathbb{R}$ smooth maps. Consider the following first order PDE system for $u: \mathcal{O} \rightarrow \mathbb{R}$ :

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial x}=A(x, y, u(x, y))  \tag{2.4.1}\\
\frac{\partial u}{\partial y}=B(x, y, u(x, y))
\end{array}\right.
$$

The Frobenius Theorem gives a necessary and sufficient condition for the first order PDE system ( ${ }^{2} \mathrm{E} .4 .1$ ) to be solvable. We need to use this theorem extensively in the study of curve evolutions and of surfaces in $\mathbb{R}^{3}$. We will see from the proof of this theorem that, although we are dealing with a PDE, the algorithm to construct solutions of this PDE is to solve one ODE system in $x$ variable and the solve a family of ODE systems in $y$ variables. The condition given in the Frobenius Theorem guarantees that this process produces a solution of the first order PDE.

If the system $\left(\frac{\mathrm{ex}}{2.4 .1)}\right.$ admits a smooth solution $u(x, y)$, then the mixed derivatives $u_{x y}=u_{y x}$, i.e.,

$$
\begin{aligned}
u_{x y} & =\left(u_{x}\right)_{y}=\left(A(x, y, u(x, y))_{y}=A_{y}+A_{u} u_{y}=A_{y}+A_{u} B\right. \\
& =\left(u_{y}\right)_{x}=\left(B(x, y, u(x, y))_{x}=B_{x}+B_{u} u_{x}=B_{x}+B_{u} A .\right.
\end{aligned}
$$

So $A, B$ must satisfy the compatibility condition:

$$
A_{y}+A_{u} B=B_{x}+B_{u} A .
$$

 solvable:

Theorem 2.4.1. (Frobenius Theorem) Let $U_{1} \subset \mathbb{R}^{2}$ and $U_{2} \subset \mathbb{R}^{n}$ be open subsets, $A=\left(A_{1}, \ldots, A_{n}\right), B=\left(B_{1}, \ldots, B_{n}\right): U_{1} \times U_{2} \rightarrow \mathbb{R}^{n}$ smooth maps, $\left(x_{0}, y_{0}\right) \in U_{1}$, and $p_{0} \in U_{2}$. Then the following first order system

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial x}=A(x, y, u(x, y)),  \tag{2.4.2}\\
\frac{\partial u}{\partial y}=B(x, y, u(x, y)) \\
u\left(x_{0}, y_{0}\right)=p_{0}
\end{array}\right.
$$

bb
has a smooth solution for $u$ in a neighborhood of $\left(x_{0}, y_{0}\right)$ if and only if

$$
\begin{equation*}
\left(A_{i}\right)_{y}+\sum_{j=1}^{n} \frac{\partial A_{i}}{\partial u_{j}} B_{j}=\left(B_{i}\right)_{x}+\sum_{j=1}^{n} \frac{\partial B_{i}}{\partial u_{j}} A_{j}, \quad 1 \leq i \leq n . \tag{2.4.3}
\end{equation*}
$$

Proof. If $u$ is a smooth solution of $\left(\frac{\mathrm{ba}}{\mathrm{bb}} \mathrm{l}_{2} .4 .2\right)$, then $\left(u_{x}\right)_{y}=\left(u_{y}\right)_{x}$ implies that the compatibility condition $\left(2_{2}^{2} 43\right)$ must hold.

Conversely, assume $A, B$ satisfy (2.4.3). To solve ( $\left.{ }^{(2.4 .2}\right)$, we proceed as follows: The existence and uniqueness of solutions of ODE, Theorem 2.1.2, implies that there exist $\delta>0$ and $\alpha:\left(x_{0}-\delta, x_{0}+\delta\right) \rightarrow U_{2}$ satisfying

$$
\left\{\begin{array}{l}
\frac{d \alpha}{d x}=A\left(x, y_{0}, \alpha(x)\right),  \tag{2.4.4}\\
\alpha\left(x_{0}\right)=p_{0} .
\end{array}\right.
$$

For each $x \in\left(x_{0}-\delta, x_{0}+\delta\right)$, let $\beta^{x}(y)$ denote the unique solution of

$$
\left\{\begin{array}{l}
\frac{d \beta^{x}}{d y}=B\left(x, y, \beta^{x}(y)\right)  \tag{2.4.5}\\
\beta^{x}\left(y_{0}\right)=\alpha(x)
\end{array}\right.
$$

Set $\mu_{0}^{u}(x, y)=\beta^{x}(y)$. By construction, $u$ satisfies the second equation of (2.4.2) $\operatorname{and}_{00} u\left(x_{0}, y_{0}\right)=p_{0}$. It remains to prove $u$ satisfies the first equation of (2.4.2). To prove this, we compute the $y$-derivative of $z(x, y)=u_{x}-A(x, y, u(x, y))$ to get

$$
\begin{aligned}
z_{y} & =\left(u_{x}-A(x, y, u)\right)_{y}=u_{x y}-A_{y}-A_{u} u_{y}=\left(u_{y}\right)_{x}-\left(A_{y}+A_{u} B\right) \\
& =(B(x, y, u))_{x}-\left(A_{y}+A_{u} B\right)=B_{x}+B_{u} u_{x}-\left(A_{y}+A_{u} B\right) \\
& =B_{x}+B_{u} u_{x}-\left(B_{x}+B_{u} A\right)=B_{u}\left(u_{x}-A\right)=B_{u}(x, y, u(x, y)) z
\end{aligned}
$$

This proves $z$ is a solution of the following differential equation:

$$
\begin{equation*}
z_{y}(x, y)=B_{u}(x, y, z) z(x, y) \tag{2.4.6}
\end{equation*}
$$

But

$$
z\left(x, y_{0}\right)=u_{x}\left(x, y_{0}\right)-A\left(x, y_{0}, u\left(x, y_{0}\right)\right)=\alpha^{\prime}(x)-A\left(x, y_{0}, \alpha(x)\right)=0
$$

Since the constant function zero is also a solution of $\binom{\text { be }}{2.4 .6}_{\mathrm{h}}$ with 0 initial data, by the uniqueness of solutions of ODE (Theorem 2.1.2) we have $z=0$, i.e., $u$ satisfies the second equation of (2.4.2).

Remark The proof of $\mathrm{T}_{\mathrm{b}}$ heorem ${ }^{\text {bf }} \mathbf{2 . 4} 4$ gives an algorithm to construct numerical solution of (2.4.2). The algorithm is as follows: First solve the ODE (2.4.4) on the horizontal line $y=y_{0}$ by a numerical method (for example Runge-Kutta) to get $u\left(x_{k}, y_{0}\right)$ for $x_{k}=x_{0}+k \epsilon$ where $\epsilon$ is the step size in the numerical method. Then we solve the ODE system (2.4.5) on the vertical line $x=x_{n}$ for each $n$ to get the value $u\left(x_{n}, y_{m}\right)$.
bc Corollary 2.4.2. Let $U_{0}$ be an open subset of $\mathbb{R}^{2},\left(x_{0}, y_{0}\right) \in U_{0}, C \in$ $g l(n)$, and $P, Q: U_{0} \rightarrow g l(n)$ smooth maps. Then the following initial value problem

$$
\left\{\begin{array}{l}
g_{x}=g P  \tag{2.4.7}\\
g_{y}=g Q \\
g\left(x_{0}, y_{0}\right)=C
\end{array}\right.
$$

has a gl(n)-valued smooth solution $g$ defined in a neighborhood of $\left(x_{0}, y_{0}\right)$ if and only if

$$
\begin{equation*}
P_{y}-Q_{x}=[P, Q]=P Q-Q P \tag{2.4.8}
\end{equation*}
$$

Proof. This Corollary follows from Theorem $\frac{\mathrm{bf}}{2.4 .1}$ with $A(x, y, u)=$ $u P(x, y)$ and $B(x, y, u)=u Q(x, y)$. We can also compute the mixed derivatives directly as follows:

$$
\begin{aligned}
\left(g_{x}\right)_{y} & =(g P)_{y}=g_{y} P+g P_{y}=g Q P+g P_{y}=g\left(Q P+P_{y}\right), \\
=\left(g_{y}\right)_{x} & =(g Q)_{x}=g_{x} Q+g Q_{x}=g P Q+g Q_{x}=g\left(P Q+Q_{x}\right) .
\end{aligned}
$$

So the compatibility condition is

$$
Q P+P_{y}=P Q+Q_{x}
$$

which is ( ${ }^{(2 . \mathrm{E}} 4.8$ ).
Next we give an application of the Frobenius Theorem ${ }^{\text {bf }} 2.4 .1$ :
ff Theorem 2.4.3. Given a smooth function $q: \mathbb{R}^{2} \rightarrow \mathbb{R}$ and a non-zero real constant $r$, the following system of first order PDE is solvable for $q^{*}: \mathbb{R}^{2} \rightarrow \mathbb{R}:$

$$
\left\{\begin{array}{l}
\left(q^{*}\right)_{s}=-q_{s}+r \sin \left(q^{*}-q\right),  \tag{2.4.9}\\
\left(q^{*}\right)_{t}=q_{t}+\frac{1}{r} \sin \left(q^{*}+q\right)
\end{array}\right.
$$

dl
if and only if $q$ satisfies

$$
q_{s t}=\sin q \cos q
$$

SGE.
Moreover, if $q$ is a solution of the $S G E$, then the solution $q^{*}$ of $\left(\frac{\mathrm{dl}}{6.3 .12)}\right.$ is again a solution of the SGE.

Proof. If $\left(\frac{\mathrm{dl}}{6.3 .12}\right)$ has a $C^{2}$ solution $q^{*}$, then the mixed derivatives must be equal. Compute directly to see that

$$
\begin{aligned}
\left(\left(q^{*}\right)_{s}\right)_{t} & =-q_{s t}+r \cos \left(q^{*}-q\right)\left(q^{*}-q\right)_{t} \\
& =-q_{s t}+r \cos \left(q^{*}-q\right)\left(\frac{1}{r} \sin \left(q^{*}+q\right)\right),
\end{aligned}
$$

so we get

$$
\begin{equation*}
\left(\left(q^{*}\right)_{s}\right)_{t}=-q_{s t}+\cos \left(q^{*}-q\right) \sin \left(q^{*}+q\right) \tag{2.4.10}
\end{equation*}
$$

A similar computation implies that

$$
\begin{equation*}
\left(\left(q^{*}\right)_{t}\right)_{s}=-q_{t s}+\cos \left(q^{*}+q\right) \sin \left(q^{*}-q\right) . \tag{2.4.11}
\end{equation*}
$$

Since both $q$ and $q^{*}$ are $C^{2}, q_{s t}=q_{t s}$ and $q_{s t}^{*}=q_{t s}^{*}$. Subtract $\left(\frac{\left.{ }^{\frac{f}{2} \mathrm{e}} \mathrm{Z} .4 .11\right)}{2}\right)$ from (2.4.10) to get

$$
2 q_{s t}=\sin \left(q^{*}+q\right) \cos \left(q^{*}-q\right)-\sin \left(q^{*}-q\right) \cos \left(q^{*}+q\right)=\sin (2 q) .
$$



$$
2 q_{s t}^{*}=\sin \left(q^{*}+q\right) \cos \left(q^{*}-q\right)+\sin \left(q^{*}-q\right) \cos \left(q^{*}+q\right)=\sin (2 q *)
$$

In the above computation we use the addition formulas

$$
\sin (A \pm B)=\sin A \cos B \pm \cos A \sin B
$$

This completes proof of the theorem.
This theorem says that if we know one solution of the non-linear wave equation SGE then we can solve the first order system (6.3.12), which is just solving two ODEs, to construct a family of solutions of the SGE (one for each real constant $r$ ). Note that $q=0$ is a trivial solution of the SGE. Theorem $\overline{2} .4 .3$ implies that system ( $\overline{2} .4 .10$ ) can be solved for $q^{*}$ with $q=0$ :

$$
\left\{\begin{array}{l}
q_{s}^{*}=r \sin q^{*}  \tag{2.4.12}\\
q_{t}^{*}=\frac{1}{r} \sin q^{*}
\end{array}\right.
$$

To solve this, we note that the ODE $d y / d x=\sin y$ can be solved easily using integration:

$$
\int \frac{d y}{\sin y}=\int d x
$$

which implies

$$
y(x)=2 \arctan \left(e^{x}\right)
$$

Hence we find solutions for $\left(\frac{\mathrm{fg}}{2.4 .12)}\right.$ ) and the SGE:

$$
q^{*}(s, t)=2 \arctan \left(e^{r s+\frac{1}{r} t}\right)
$$

The SGE is a so called soliton equation. A special feature of soliton equations is the existence of first order systems that can generate new solutions from an old one. System (6.3.12) was first constructed by Bäcklund in []. Thus a transformation from a given solution to a new solution is called a Bäcklund transformation. System ( $\overline{6.3 .12 \text { ) seems to }}$ come out from nowhere. But we will explain its origin from a natural geometric theorem in a later section.

### 2.5. The Smoke-ring equation and the Non-linear Schrödinger equation.

In 1906, a graduate student of Levi-Civita, da Rios, wrote a master degree thesis, in which he modeled the movement of a thin vortex in a vicous fluid by the motion of a curve propagating in $\mathbb{R}^{3}$ according to:

$$
\begin{equation*}
\gamma_{t}=\gamma_{x} \times \gamma_{x x} \tag{2.5.1}
\end{equation*}
$$

This is called the vortex filament equation or smoke ring equation, and it can viewed as a dynamical system on the space of curves in $\mathbb{R}^{3}$.

Proposition 2.5.1. If $\gamma(x, t)$ is a solution of the smoke ring equation (2.5.1) and $\left\|\gamma_{x}(x, 0)\right\|=1$ for all $x$, then $\left\|\gamma_{x}(x, t)\right\|=1$ for all $(x, t)$. In other words, if $\gamma(\cdot, 0)$ is parametrized by arc length, then so is $\gamma(\cdot, t)$ for all $t$.

Proof. It suffices to prove that $\left(\gamma_{x} \cdot \gamma_{x}\right)_{t}=0$. To see this, we compute directly to get

$$
\begin{aligned}
\left(\gamma_{x} \cdot \gamma_{x}\right)_{t} & =2 \gamma_{x t} \cdot \gamma_{x}=2\left(\gamma_{t}\right)_{x} \cdot \gamma_{x}=2\left(\gamma_{x} \times \gamma_{x x}\right)_{x} \cdot \gamma_{x} \\
& =2\left(\gamma_{x x} \times \gamma_{x x}+\gamma_{x} \times \gamma_{x x x}\right) \cdot \gamma_{x}=0+0=0 .
\end{aligned}
$$

So for a solution $\gamma$ of $\left(\frac{\mathrm{l} \text { W. }}{2 \mathrm{~W}} .1\right)$, we may assume that $\gamma(\cdot, t)$ is parametrized by arc length for all $t$. Next we explain the geometric meaning of the evolution equation $(2.5 .1)$ on the space of curves in $\mathbb{R}^{3}$. Let $\left(e_{1}, e_{2}, e_{3}\right)(\cdot, t)$ denote the Frenet frame of the curve $\gamma(\cdot, t)$. Since $\gamma_{x}=e_{1}$ and $\gamma_{x x}=$ $\left(e_{1}\right)_{x}=k e_{2}$, the curve flow (2.5.1) becomes

$$
\gamma_{t}=k e_{1} \times e_{2}=k e_{3} .
$$

In other words, the curve flow (2.5.1) moves in the direction of binormal with curvature ${ }_{1}$ as its speed. If we use a parallel frame $\left(e_{1}, v_{1}, v_{2}\right)$, then by Corollary 2.3.4

$$
k e_{3}=k\left(\sin \beta v_{1}+\cos \beta v_{2}\right)=-k_{2} v_{1}+k_{1} v_{2}
$$

So the curve flow (2.5.1) becomes

$$
\begin{equation*}
\gamma_{t}=-k_{2} v_{1}+k_{1} v_{2} \tag{2.5.2}
\end{equation*}
$$

By Proposition ${ }_{2 c f}^{\text {2.3.3 }}, \alpha^{\prime \prime}=k_{1} v_{1}+k_{2} v_{2}$, so the above equation can be viewed as

$$
\begin{equation*}
\gamma_{t}=J_{t}\left(\alpha^{\prime \prime}\right)=J_{t}\left(k_{1} v_{1}+k_{2} v_{2}\right) \tag{2.5.3}
\end{equation*}
$$

where $J_{t}$ is the rotation of $\frac{\pi}{2}$ on the normal plane at $\alpha(t)$. Note that the parallel frame $\left(v_{1}, v_{2}\right)$ is not unique. Since $\alpha^{\prime \prime}=k_{1} v_{1}+k_{2} v_{2}$ and the rotation of $\frac{\pi}{2}$ in the normal plane is independent of the choice of the parallel frame, the right hand side of equation (2.5.3) is independent of the choice of parallel frame.

Next we want to derive the evolution equation of principal curvatures for $\gamma(\cdot, t)$ for a solution of the smoke-ring equation. We will show that they evolve according to the non-linear Schrödinger equation (NLS), which is the same equation that models the motion of a wave envelope travelling in an optic fiber:

$$
\begin{equation*}
q_{t}=i\left(q_{x x}+\frac{1}{2}|q|^{2} q\right) . \tag{NLS}
\end{equation*}
$$


#### Abstract

$\square$


We want to compute the evolution of $\left(e_{1}, v_{1}, v_{2}\right)$, i.e., $\left(e_{1}, v_{1}, v_{2}\right)_{t}$. We first compute

$$
\begin{aligned}
\left(e_{1}\right)_{t} & =\left(\gamma_{x}\right)_{t}=\left(\gamma_{t}\right)_{x}=\left(-k_{2} v_{1}+k_{1} v_{2}\right)_{x} \\
& =-\left(k_{2}\right)_{x} v_{1}-k_{2}\left(v_{1}\right)_{x}+\left(k_{1}\right)_{x} v_{2}+k_{1}\left(v_{2}\right)_{x}, \quad \text { by }(2.5 .5) \\
& =-\left(k_{2}\right)_{x} v_{1}-k_{2}\left(-k_{1} e_{1}\right)+\left(k_{1}\right)_{x} v_{2}+k_{1}\left(-k_{2} e_{1}\right) \\
& =-\left(k_{2}\right)_{x} v_{1}+\left(k_{1}\right)_{x} v_{2} .
\end{aligned}
$$

Since $\left(e_{1}, v_{1}, v_{2}\right)$ is o.n., by Proposition ${ }^{\text {2.1. }}$. 1 there exists a function $u$ so that

$$
\left(e_{1}, v_{1}, v_{2}\right)_{t}=\left(e_{1}, v_{1}, v_{2}\right)\left(\begin{array}{ccc}
0 & \left(k_{2}\right)_{x} & -\left(k_{1}\right)_{x} \\
-\left(k_{2}\right)_{x} & 0 & -u \\
\left(k_{1}\right)_{x} & u & 0
\end{array}\right)
$$

or equivalently,

$$
\left\{\begin{array}{l}
\left(e_{1}\right)_{t}=-\left(k_{2}\right)_{x} v_{1}+\left(k_{1}\right)_{x} v_{2} \\
\left(v_{1}\right)_{t}=\left(k_{2}\right)_{x} e_{1}+u v_{2} \\
\left(v_{2}\right)_{t}=-\left(k_{1}\right)_{x} e_{1}-u v_{1}
\end{array}\right.
$$

We claim that

$$
u_{x}=-\frac{1}{2}\left(k_{1}^{2}+k_{2}^{2}\right)_{x} .
$$

To prove this, we compute

$$
\begin{aligned}
\left(v_{1}\right)_{x t} \cdot v_{2} & =\left(-k_{1} e_{1}\right)_{t} \cdot v_{2}=-k_{1}\left(e_{1}\right)_{t} \cdot v_{2} \\
& =-k_{1}\left(-\left(k_{2}\right)_{x} v_{1}+\left(k_{1}\right)_{x} v_{2}\right) \cdot v_{2}=-k_{1}\left(k_{1}\right)_{x}, \\
=\left(v_{1}\right)_{t x} \cdot v_{2} & =\left(\left(k_{2}\right)_{x} e_{1}+u v_{2}\right)_{x} \cdot v_{2}=\left(k_{2}\right)_{x}\left(e_{1}\right)_{x} \cdot v_{2}+u_{x} \\
& =\left(k_{2}\right)_{x}\left(k_{1} v_{1}+k_{2} v_{2}\right) \cdot v_{2}+u_{x}=\left(k_{2}\right)_{x} k_{2}+u_{x} .
\end{aligned}
$$

This proves the claim. Hence

$$
u(x, t)=-\frac{1}{2}\left(k_{1}^{2}+k_{2}^{2}\right)+c(t)
$$

for some smooth function $c(t)$. Remember that for each fixed $t$, we can rotate $\left(v_{1}, v_{2}\right)(\cdot, t)$ by a constant angle $\theta(t)$ to another parallel normal frame $\left(\tilde{v}_{1}, \tilde{v}_{2}\right)$ of $\gamma(\cdot, t)$, i.e.,

$$
\left\{\begin{array}{l}
\tilde{v}_{1}(x, t)=\cos \theta(t) v_{1}(x, t)+\sin \theta(t) v_{2}(x, t) \\
\tilde{v}_{2}(x, t)=-\sin \theta(t) v_{1}(x, t)+\cos \theta(t) v_{2}(x, t)
\end{array}\right.
$$

A direct computation shows that

$$
\tilde{u}=\left(\tilde{v}_{1}\right)_{t} \cdot \tilde{v}_{2}=u+\frac{d \theta}{d t}=-\frac{1}{2}\left(k_{1}^{2}+k_{2}^{2}\right)+c(t)+\theta^{\prime}(t)
$$

If we choose $\theta$ so that $\theta^{\prime}=-c$, then the new parallel frame satisfies

$$
\left(e_{1}, \tilde{v}_{1}, \tilde{v}_{2}\right)_{t}=\left(e_{1}, \tilde{v}_{1}, \tilde{v}_{2}\right)\left(\begin{array}{ccc}
0 & \left(\tilde{k}_{2}\right)_{x} & -\left(\tilde{k}_{1}\right)_{x} \\
-\left(\tilde{k}_{2}\right)_{x} & 0 & \frac{\tilde{k}_{1}^{2}+\tilde{k}_{2}^{2}}{2} \\
\left(\tilde{k}_{1}\right)_{x} & -\frac{\tilde{k}_{1}^{2}+\tilde{k}_{2}^{2}}{2} & 0
\end{array}\right) .
$$

So we have proved the first part of the following Theorem:
ba Theorem 2.5.2. Suppose $\gamma(x, t)$ is a solution of the smoke ring equation $(2.5 .1)$ and $\left\|\gamma_{x}(x, 0)\right\|=1$ for all $x$. Then
(1) there exists a parallel normal frame $\left(v_{1}, v_{2}\right)(\cdot, t)$ for each curve $\gamma(\cdot, t)$ so that

$$
\left\{\begin{array}{l}
g_{x}=g\left(\begin{array}{ccc}
0 & -k_{1} & -k_{2} \\
k_{1} & 0 & 0 \\
k_{2} & 0 & 0
\end{array}\right)  \tag{2.5.6}\\
g_{t}=g\left(\begin{array}{ccc}
0 & \left(k_{2}\right)_{x} & -\left(k_{1}\right)_{x} \\
-\left(k_{2}\right)_{x} & 0 & \frac{k_{1}^{2}+k_{2}^{2}}{2} \\
\left(k_{1}\right)_{x} & -\frac{k_{1}^{2}+k_{2}^{2}}{2} & 0
\end{array}\right)
\end{array}\right.
$$

where $g=\left(e_{1}, v_{1}, v_{2}\right)$ and $k_{1}(\cdot, t)$ and $k_{2}(\cdot, t)$ are the principal curvatures of $\gamma(\cdot, t)$ with respect to $v_{1}(\cdot \cdot, t), v_{2}(\cdot, t)$ respectively,
(2) $q=k_{1}+i k_{2}$ is a soluton of the NLS (2.5.4).

Proof. We have proved (1). For (2), we use ( $(2.5 .6)$ to compute the evolution of $\left(k_{1}\right)_{t}$ and $\left(k_{2}\right)_{t}$ :

$$
\begin{aligned}
\left(k_{1}\right)_{t} & =\left(\left(e_{1}\right)_{x} \cdot v_{1}\right)_{t}=\left(e_{1}\right)_{x t} \cdot v_{1}+\left(e_{1}\right)_{x} \cdot\left(v_{1}\right)_{t} \\
& =\left(e_{1}\right)_{t x} \cdot v_{1}+\left(k_{1} v_{1}+k_{2} v_{2}\right) \cdot\left(\left(k_{2}\right)_{x} e_{1}-\left(\frac{k_{1}^{2}+k_{2}^{2}}{2}\right) v_{2}\right) \\
& =\left(-\left(k_{2}\right)_{x} v_{1}+\left(k_{1}\right)_{x} v_{2}\right)_{x} \cdot v_{1}-\frac{k_{2}\left(k_{1}^{2}+k_{2}^{2}\right)}{2} \\
& =-\left(k_{2}\right)_{x x}-\frac{k_{2}\left(k_{1}^{2}+k_{2}^{2}\right)}{2} .
\end{aligned}
$$

Similar computation gives

$$
\left(k_{2}\right)_{t}=\left(k_{1}\right)_{x x}+\frac{k_{1}\left(k_{1}^{2}+k_{2}^{2}\right)}{2} .
$$

Let $q=k_{1}+i k_{2}$. Then we have

$$
q_{t}=i\left(q_{x x}+\frac{|q|^{2}}{2} q\right)
$$

The converse of Theorem ${ }^{\text {ba }} 2.5 .2$ is also true. Given a solution $q=$ $k_{1}+i k_{2}$ of the NLS (2.5.4), we need to first solve the first order PDE system (2.5.6) to get $\left(e_{1}, v_{1}, v_{2}\right)$, then solve another first order PDE system

$$
\left\{\begin{array}{l}
\gamma_{x}=e_{1}  \tag{2.5.7}\\
\gamma_{t}=-k_{2} v_{1}+k_{1} v_{2}
\end{array}\right.
$$

Then $\gamma(x, t)$ satisfies the smoke-ring equation (2.5.5.1). Now we are ready to prove the converse of Theorem 2.5.2. Let $q=k_{1}+i k_{2}$ be a solution of NLS. To construct a solution $\gamma(x, t)$ of the smoke ring equation (2.5.1), we need to solve two first order PDE systems ( ${ }_{\text {az }}^{2 z .5} 5$.6)
 the compatibility condition is (2.4.8), i.e.,

$$
\begin{equation*}
P_{t}-Q_{x}=[P, Q], \tag{2.5.8}
\end{equation*}
$$

where

$$
P=\left(\begin{array}{ccc}
0 & -k_{1} & -k_{2} \\
k_{1} & 0 & 0 \\
k_{2} & 0 & 0
\end{array}\right), \quad Q=\left(\begin{array}{ccc}
0 & \left(k_{2}\right)_{x} & -\left(k_{1}\right)_{x} \\
-\left(k_{2}\right)_{x} & 0 & \frac{k_{1}^{2}+k_{2}^{2}}{2} \\
\left(k_{1}\right)_{x} & -\frac{k_{1}^{2}+k_{2}^{2}}{2} & 0
\end{array}\right)
$$

Since $P, Q$ are skew-symmetric,

$$
[P, Q]^{t}=(P Q-Q P)^{t}=\left(Q^{t} P^{t}-P^{t} Q^{t}\right)=(Q P-P Q)=-[P, Q]
$$

Hence $[P, Q]$ is is skew-symmetric. So we only need to check the 21 , 31 and 32 entries of ( 2.5 .8 ). A direct computation shows that ( ${ }^{[6] 5.8 \text { ) }}$ holds if and only if

$$
\left\{\begin{array}{l}
\left(k_{1}\right)_{t}+\left(k_{2}\right)_{x x}=-\frac{k_{2}\left(k_{1}^{2}+k_{2}^{2}\right)}{2} \\
\left(k_{2}\right)_{t}-\left(k_{1}\right)_{x x}=\frac{k_{1}\left(k_{1}^{2}+k_{2}^{2}\right)}{2}
\end{array}\right.
$$

which is the NLS equation for $q=k_{1}+i k_{2}$. So by Corollary ${ }_{2}^{\mathrm{bc}} 2.4 .2$, given a constant orthogonal matrix $\left(u_{1}^{0}, u_{2}^{0}, u_{3}^{0}\right)$, there exists a unique solution $g=\left(e_{1}, v_{1}, v_{2}\right)$ for system $(2.5 .6)$ with initial data $g(0)=\left(u_{1}^{0}, u_{2}^{0}, u_{3}^{0}\right)$. To construct the solution $\gamma$ for the smoke ring ${ }_{i}$ equation, we need to solve ( $\left(\frac{D_{1}}{2.5 .7}\right)$. The compatibility condition for $\left(\frac{B_{1}}{2.5} .7\right)$ is given by $\gamma_{x t}=$ $\gamma_{t x}$, which is

$$
\left(e_{1}\right)_{t}=\left(-k_{2} v_{1}+k_{1} v_{2}\right)_{x} .
$$

This holds because $\left(e_{1}, v_{1}, v_{2}\right)$ satisfies $\left(\frac{\text { az }}{2.5} .6\right)$. We carry out this direct computation:

$$
\begin{aligned}
\left(e_{1}\right)_{t} & =-\left(k_{2}\right)_{x} v_{1}+\left(k_{1}\right)_{x} v_{2}, \\
\left(-k_{2} v_{1}+k_{1} v_{2}\right)_{x} & =-\left(k_{2}\right)_{x} v_{1}-k_{2}\left(v_{1}\right)_{x}+\left(k_{1}\right)_{x} v_{2}+k_{1}\left(v_{2}\right)_{x} \\
& =-\left(k_{2}\right)_{x} v_{1}+k_{2} k_{1} e_{1}+\left(k_{1}\right)_{x} v_{2}-k_{1}\left(k_{2} e_{1}\right) \\
& =-\left(k_{2}\right)_{x} v_{1}+\left(k_{1}\right)_{x} v_{2} .
\end{aligned}
$$

In other word, we have proved
ce Theorem 2.5.3. Let $q=k_{1}+i k_{2}$ be a solution of $N L S, p_{0} \in \mathbb{R}^{3}$, and $\left(u_{1}, u_{2}, u_{3}\right)$ a constant orthogonal matrix. Then there exist a solution $\gamma$ of the smoke ring equation and a parallel frame $\left(e_{1}, v_{1}, v_{2}\right)(\cdot, t)$ for $\gamma(\cdot, t)$ so that $k_{1}(\cdot, t), k_{2}(\cdot, t)$ are the principal curvatures along $v_{1}(\cdot, t), v_{2}(\cdot, t)$ respectively, $\gamma(0,0)=p_{0}$, and $\left(e_{1}, v_{1}, v_{2}\right)(0,0)=\left(u_{1}, u_{2}, u_{3}\right)$.

Note that given a solution $q=k_{1}+i k_{2}$ of NLS, the proof of Theroem 2.5.3 gives an algorithm to construct numerical solution of the smokering equation.

## 3. Surfaces in $\mathbb{R}^{3}$

### 3.1. Some linear algebra.

Let $V$ be a vector space, and $($,$) an inner product on V$. A linear $\operatorname{map} T: V \rightarrow V$ is self-adjoint if

$$
\left(A\left(v_{1}\right), v_{2}\right)=\left(v_{1}, A\left(v_{2}\right)\right) \quad \forall v_{1}, v_{2} \in V .
$$

The following is the Spectral Theorem for self-adjoint operators:
aa Theorem 3.1.1. Suppose $V$ is an $n$-dimensional vector space equipped with an inner product (,). If $T: V \rightarrow V$ is a linear self-adjoint operator, then
(1) eigenvalues of $T$ are real,
(2) there is an orthonormal basis $v_{1}, \cdots, v_{n}$ of $V$ that are eigenvectors of $T$, i.e., $T\left(v_{i}\right)=\lambda_{i} v_{i}$ for $1 \leq i \leq n$.

A bilinear functional $b: V \times V \rightarrow \mathbb{R}$ is symmetric if

$$
b(v, w)=b(v, w) \quad \text { for all } v, w \in V
$$

A symmetric bilinear form $b$ is positive definite if $b(v, v)>0$ for all non-zero $v \in V$.

We associate to each self-adjoint linear map $T$ on $V$ a symmetric bilinear functional:

Theorem 3.1.2. Let $T: V \rightarrow V$ be a self-adjoint linear map, and $b: V \times V \rightarrow \mathbb{R}$ the map defined by

$$
b\left(v_{1}, v_{2}\right)=\left(T\left(v_{1}\right), v_{2}\right)
$$

Then
(1) $b$ is a symmetric bilinear form on $V$,
(2) all eigenvalues of $T$ are positive if and only if $b$ is positive definite.

Given a linear map $T: V \rightarrow V$ and a basis $v_{1}, \cdots, v_{n}$ of $T$, we can associate to $T$ a matrix $A$ so that $T$ is represented by the multiplication of $A$. This is because each $T\left(v_{i}\right)$ can be written as a linear combination of the basis:

$$
T\left(v_{i}\right)=\sum_{j} a_{j i} v_{j}
$$

Write $v=\sum_{j} x_{j} v_{j}$, and $T(v)=\sum_{j} y_{j} v_{j}$. A direct computation gives $y=A x$, where $A=\left(a_{i j}\right), x=\left(x_{1}, \cdots, x_{n}\right)^{t}$, and $y=\left(y_{1}, \cdots, y_{n}\right)^{t}$. In other words, after choosing a basis of $V$, the operator $T$ looks like the map $x \mapsto A x$. We call $A$ the matrix of $T$ associated to the basis
$v_{1}, \cdots, v_{n}$. If we change basis, the corresponding matrix of $T$ change by a conjugation:

Proposition 3.1.3. Let $T: V \rightarrow V$ be a linear map, and $A$ and $B$ the matrices of $T$ associated to bases $\left\{v_{1}, \cdots, v_{n}\right\}$ and $\left\{u_{1}, \cdots, u_{n}\right\}$ of $V$ respectively. Write $u_{i}=\sum_{j} g_{j i} v_{j}$. Then $B=g^{-1} A g$, where $g=\left(g_{i j}\right)$.

Below we list some properties of determinants:

$$
\begin{aligned}
& \operatorname{det}(A B)=\operatorname{det}(A) \operatorname{det}(B), \\
& \operatorname{det}\left(A^{t}\right)=\operatorname{det}(A), \\
& \operatorname{det}\left(g A g^{-1}\right)=\operatorname{det}(A) .
\end{aligned}
$$

Let $A$ denote the matrix associated to the linear operator $T: V \rightarrow V$ with respect to the basis $v_{1}, \cdots, v_{n}$. Define

$$
\operatorname{tr}(T)=\operatorname{tr}(A), \quad \operatorname{det}(T)=\operatorname{det}(A)
$$

where $A$ is the matrix of $T$ associated to the basis $v_{1}, \cdots, v_{n}$. Since

$$
\operatorname{tr}\left(g A g^{-1}\right)=\operatorname{tr}(A), \quad \operatorname{det}\left(g A g^{-1}\right)=\operatorname{det}(A)
$$

$\operatorname{tr}(T)$ and $\operatorname{det}(T)$ are independent of the choice of basis. Moreover, if $T$ has an eignbasis with eigenvalues $\lambda_{1}, \cdots, \lambda_{n}$, then

$$
\operatorname{tr}(T)=\sum_{j=1}^{n} \lambda_{j}, \quad \operatorname{det}(T)=\prod_{j=1}^{n} \lambda_{j} .
$$

Let $v_{1}, \cdots, v_{n}$ be a basis of $V, b: V \times V \rightarrow \mathbb{R}$ a symmetric bilinear form, and $b_{i j}=b\left(v_{i}, v_{j}\right)$. Given $\xi=\sum_{i} x_{i} v_{i}$ and $\eta=\sum_{i} y_{i} v_{i}$, we have

$$
b(\xi, \eta)=\sum_{i j} b_{i j} x_{i} y_{j}=X^{t} B Y
$$

where $X^{t}=\left(x_{1}, \ldots, x_{n}\right), Y^{t}=\left(y_{1}, \ldots, y_{n}\right)$, and $B=\left(b_{i j}\right)$. We call $B=\left(b_{i j}\right)$ the matrix associated to $b$ with respect to basis $v_{1}, \cdots, v_{n}$.
ac Proposition 3.1.4. Let $(V,()$,$) be an inner product space, T: V \rightarrow V$ a self-adjoint operator, $b(X, Y)=(T(X), Y)$ the symmetric bilinear form associated to $T$, and $A=\left(a_{i j}\right), B=\left(b_{i j}\right)$ the matrices of $T$ and $b$ associated to basis $v_{1}, \cdots, v_{n}$ respectively. Let $g_{i j}=\left(v_{i}, v_{j}\right), G=\left(g_{i j}\right)$, and $G^{-1}=\left(g^{i j}\right)$ the inverse of $G$. Then
(1) $B=A^{t} G$,
(2) $A=G^{-1} B$,
(3) $\operatorname{det}(A)=\frac{\operatorname{det}(B)}{\operatorname{det}(G)}, \operatorname{tr}(A)=\sum_{i, j} b_{i j} g^{i j}$.

Proof. Compute directly to get

$$
b_{i j}=\left(T\left(v_{i}\right), v_{j}\right)=\left(\sum_{k} a_{k i} v_{k}, v_{j}\right)=\sum_{k} a_{k i} g_{k j}=\left(A^{t} G\right)_{i j} .
$$

This proves (1). Statement (2) and (3) are consequences of (1).
Suppose $b: V \times V \rightarrow \mathbb{R}$ is the symmetric bilinear form associated to the self-adjoint operator $T: V \rightarrow V$, and $v_{1}, \cdots, v_{n}$ are orthonormal eigenbasis of $T$ with eigenvalues $\lambda_{1}, \cdots, \lambda_{n}$. Then the diagonal matrix $\operatorname{diag}\left(\lambda_{1}, \cdots, \lambda_{n}\right)$ is the matrix associated to both $T$ and $b$ with respect to the basis $v_{1}, \cdots, v_{n}$. Moreover, given $v \in V$, write $v=\sum_{i=1}^{n} x_{i} v_{i}$, then

$$
b(v, v)=b\left(\sum_{i=1}^{n} x_{i} v_{i}, \sum_{j=1}^{n} x_{j} v_{j}\right)=\left(\sum_{i=1}^{n} \lambda_{i} x_{i} v_{i}, \sum_{j=1}^{n} x_{j} v_{j}\right)=\sum_{i=1}^{n} \lambda_{i} x_{i}^{2} .
$$

In other words, the quadratic form $Q(v)=b(v, v)$ associated to $b$ is diagonalized. We also see that the maximum (minimum resp.) of $b(v, v)$ taken among all $\|v\|=1$ is the largest (smallest resp.) eigenvalue of $T$.

## Tensor algebra

Let $V$ be a finite dimensional vector space, and the dual $V^{*}=$ $L(V, \mathbb{R})$ the space of all linear functional from $V$ to $\mathbb{R}$. Note that $V$ is naturally isomorphic to $V^{* *}=\left(V^{*}\right)^{*}=L\left(V^{*}, \mathbb{R}\right)$ via the map

$$
\phi(v)(\ell)=\ell(v), \quad v \in V, \quad \ell \in V^{*}
$$

The tensor product $V_{1}^{*} \otimes \cdots V_{k}^{*}$ is the vector space of all multilinear functionals $b: V_{1} \times \cdots \times V_{k} \rightarrow \mathbb{R}$, and the tensor product

$$
\mathcal{T}_{m}^{k}(V)=\otimes^{k} V^{*} \otimes^{m} V
$$

is the space of all multilinear functionals from $V \times \cdots V \times V^{*} \times \cdots \times V^{*}$ to $\mathbb{R}$ (here we use $k$ copies of $V$ and $m$ copies of $V^{*}$ ).

Given $f_{i} \in V_{i}^{*}$ and $\xi_{i+k} \in V_{i+k}$, let $f_{1} \otimes \cdots \otimes f_{k} \otimes \xi_{k+1} \otimes \cdots \otimes \xi_{k+m}$ be the element in $V_{1}^{*} \otimes V_{k}^{*} \otimes V_{k+1} \otimes \cdots \otimes V_{k+m}$ defined by

$$
\begin{array}{r}
\left(f_{1} \otimes \cdots \otimes f_{k} \otimes \xi_{k+1} \otimes \cdots \otimes \xi_{k+m}\right)\left(u_{1}, \cdots, u_{k}, h_{k+1}, \cdots, h_{k+m}\right) \\
=f_{1}\left(u_{1}\right) \cdots f_{k}\left(u_{k}\right) h_{k+1}\left(\xi_{k+1}\right) \cdots h_{k+m}\left(\xi_{k+m}\right)
\end{array}
$$

for $u_{i} \in V_{i}$ and $h_{k+i} \in V_{k+i}^{*}$.
Proposition 3.1.5. Let $v_{1}, \cdots, v_{n}$ be a basis of $V$, and $\ell_{1}, \cdots, \ell_{n} \in V^{*}$ the dual basis, i.e., $\ell_{i}\left(v_{j}\right)=\delta_{i j}$. Let $b: V \times V \rightarrow \mathbb{R}$ be a bilinear functional, and $b_{i j}=b\left(v_{i}, v_{j}\right)$. Then

$$
b=\sum_{i j} b_{i j} \ell_{i} \otimes \ell_{j} .
$$

Let $L\left(V_{1}, V_{2}\right)$ denote the space of all linear maps from $V_{1}$ to $V_{2}$. Given $\ell \in V_{1}^{*}$ and $v \in V_{2}$, the tensor product $\ell \otimes v$ can be identified as the linear map from $V_{1}$ to $V_{2}$ defined by $u \mapsto \ell(u) v$. It is easy to check that

$$
L\left(V_{1}, V_{2}\right) \simeq V_{1}^{*} \otimes V_{2}
$$

In fact,
Proposition 3.1.6. Let $\left\{v_{1}, \cdots, v_{n}\right\}$ and $\left\{w_{1}, \cdots, w_{m}\right\}$ be basis of $V_{1}$ and $V_{2}$ respectively, and $\ell_{1}, \cdots, \ell_{n}$ the dual basis of $V_{1}^{*}$. Suppose $T$ : $V_{1} \rightarrow V_{2}$ is a linear map, and $T\left(v_{i}\right)=\sum_{j=1}^{m} a_{j i} w_{j}$ for $1 \leq i \leq n$. Then

$$
T=\sum_{i j} a_{j i} \ell_{i} \otimes w_{j}
$$

dp Corollary 3.1.7. Suppose $\left\{v_{1}, \cdots, v_{n}\right\}$ is a basis of $V$ and $\left\{\ell_{1}, \cdots, \ell_{n}\right\}$ is the dual basis of $V^{*}$. Then the identity map $\mathrm{Id}=\sum_{i=1}^{n} \ell_{i} \otimes e_{i}$.

A multilinear functional $u \in \otimes^{k}\left(V^{*}\right)$ is called alternating if

$$
u\left(\xi_{\tau(1)}, \cdots, \xi_{\tau(k)}\right)=\operatorname{sgn}(\tau) u\left(\xi_{1}, \cdots, \xi_{k}\right)
$$

for all $\xi_{1}, \cdots, \xi_{k} \in V$, where $\tau$ a permutation of $\{1, \cdots, k\}$ and $\operatorname{sgn}(\tau)$ is the sign of $\tau$.

For $\theta_{1}, \theta_{2} \in V^{*}$, let $\theta_{1} \wedge \theta_{2}$ denote the bilinear functional on $V$ defined by

$$
\left(\theta_{1} \wedge \theta_{2}\right)\left(\xi_{1}, \xi_{2}\right)=\theta_{1}\left(\xi_{1}\right) \theta_{2}\left(\xi_{2}\right)-\theta_{2}\left(\xi_{1}\right) \theta_{1}\left(\xi_{2}\right)
$$

By definition, $\theta_{1} \wedge \theta_{2}$ is alternating. So $\theta_{1} \wedge \theta_{2}$ is the anti-symmetrization of $\theta_{1} \otimes \theta_{2}$.

Let $\wedge^{k}\left(V^{*}\right)$ denote the space of all alternating functionals in $\otimes^{k}\left(V^{*}\right)$. It follows from definitions that:

Proposition 3.1.8. (1) If $\theta \in V^{*}$, then $\theta \wedge \theta=0$.
(2) If $\theta_{1}, \theta_{2} \in V^{*}$, then $\theta_{1} \wedge \theta_{2}=-\theta_{2} \wedge \theta_{1}$.
(3) If $\ell_{1}, \cdots, \ell_{n}$ is a basis of $V^{*}$, then

$$
\left\{\ell_{i_{1}} \wedge \cdots \wedge \ell_{i_{k}} \mid 1 \leq i_{1}<\cdots<i_{k} \leq n\right\}
$$

is a basis of $\wedge^{k}\left(V^{*}\right)$. In particular, $\operatorname{dim}\left(\wedge^{n}\left(V^{*}\right)\right)=1$ and $\operatorname{dim}\left(\wedge^{k}\left(V^{*}\right)\right)=0$ if $k>n$.

### 3.2. Embedded surfaces in $\mathbb{R}^{3}$.

Let $\mathcal{O}$ be an open subset of $\mathbb{R}^{2}$. A smooth map $f: \mathcal{O} \rightarrow \mathbb{R}^{3}$ is called an immersion if $f_{x_{1}}(p), f_{x_{2}}(p)$ are linearly independent for all $p \in \mathcal{O}$.

Let $M$ be a subset of $\mathbb{R}^{3}$. A subset $U$ of $M$ is open in the induced topology if there exists an open subset $U_{0}$ of $\mathbb{R}^{n}$ so that $U=U_{0} \cap M$.

Definition 3.2.1. A subset $M$ of $\mathbb{R}^{3}$, equipped with the induced topology, is called an embedded surface if there exist an open cover $\left\{U_{\alpha} \mid \alpha \in I\right\}$ of $M$ and a collection of homeomorphisms $\phi_{\alpha}: \mathcal{O}_{\alpha} \rightarrow U_{\alpha}$ from open subset $\mathcal{O}_{\alpha}$ of $\mathbb{R}^{2}$ onto $U_{\alpha}$ such that
(1) $\phi_{\alpha}: \mathcal{O}_{\alpha} \rightarrow U_{\alpha} \subset \mathbb{R}^{3}$ is an immersion,
(2) $\phi_{\beta}^{-1} \phi_{\alpha}: \phi_{\alpha}^{-1}\left(U_{\alpha} \cap U_{\beta}\right) \rightarrow \phi_{\beta}^{-1}\left(U_{\alpha} \cap U_{\beta}\right)$ is a smooth diffomorphism for any $\alpha, \beta \in I$.
The collection $\mathcal{A}=\left\{\phi_{\alpha}: \mathcal{O}_{\alpha} \rightarrow U_{\alpha} \mid \alpha \in I\right\}$ is called an atlas of $M$. A homeomorphism $\phi: \mathcal{O} \rightarrow U \subset M$ is said to be compatible to $\mathcal{A}$ if

$$
\phi_{\alpha}^{-1} \circ \phi: \phi^{-1}\left(U \cap U_{\alpha}\right) \rightarrow \phi_{\alpha}^{-1}(U \cap U \alpha)
$$

is a diffeomorphism for all $\alpha \in I$. Such $\phi$ is called a local coordinate system or a local chart on $M$.

## Example 3.2.2. Examples of embedded surfaces in $\mathbb{R}^{3}$

(1) Graph.

Let $\mathcal{O}$ be an open subset of $\mathbb{R}^{2}$, and $u: \mathcal{O} \rightarrow \mathbb{R}$ a smooth proper map. Here a map is proper if the preimage of any compact subset is compact. The graph

$$
M=\left\{\left(x_{1}, x_{2}, u\left(x_{1}, x_{2}\right)\right) \mid x_{1}, x_{2} \in \mathcal{O}\right\}
$$

is an embedded surface in $\mathbb{R}^{3}$ with one chart defined by $\phi\left(x_{1}, x_{2}\right)=$ $\left(x_{1}, x_{2}, u\left(x_{1}, x_{2}\right)\right)$.
(2) Surface of revolution.

Let $\alpha(s)=(y(s), z(s)):(a, b) \rightarrow \mathbb{R}^{2}$ be the parametrization of a smooth curve $C$ in the $y z$-plane, and $M$ the set obtained by rotating the curve $C$ along the $y$-axis, i.e.,
$M=\{(z(s) \cos t, y(s), z(s) \sin t) \mid s \in(a, b), t \in[0,2 \pi]\}$.
It is easy to check that $M$ is an embedded surface in $\mathbb{R}^{3}$ (we need two charts to cover $M$ ).
(3) Let $\mathcal{O}$ be an open subset of $\mathbb{R}^{3}$, and $f: \mathcal{O} \rightarrow \mathbb{R}^{3}$ be a smooth map such that $f_{x_{1}}, f_{x_{2}}$ are linearly independent at every point of $\mathcal{O}$ and $f$ is a homeomorphism from $\mathcal{O}$ to $f(\mathcal{O})$, where $f(\mathcal{O})$ is equipped with the induced topology. Then $M=f(\mathcal{O})$ is an embedded surface of $\mathbb{R}^{3}$.
(4) If $M$ is an embedded surface of $\mathbb{R}^{3}$ and $\phi_{\alpha}: \mathcal{O}_{\alpha} \rightarrow U_{\alpha}$ is a local coordinate system on $M$, then $U_{\alpha}$ is also an embedded surface in $\mathbb{R}^{3}$.

The following Proposition is a consequence of the Implicit Function Theorem:

Proposition 3.2.3. Let $f: \mathbb{R}^{3} \rightarrow \mathbb{R}$ be a smooth function. If $c \in \mathbb{R}$ such that for all $p \in f^{-1}(c)$ the gradient

$$
\nabla f(p)=\left(f_{x_{1}}(p), f_{x_{2}}(p), f_{x_{3}}(p)\right) \neq 0
$$

then $M=f^{-1}(c)$ is an embedded surface.
Applying the above Proposition to the function

$$
f(x)=a_{1} x_{1}^{2}+a_{2} x_{2}^{2}+a_{3} x_{3}^{2},
$$

we see that the spheres, ellipsoids, hyperbolids are embedded surfaces. Apply to

$$
f\left(x_{1}, x_{2}, x_{3}\right)=x_{1}^{2}+x_{2}^{2}
$$

to see that the cylinder is an embedded surface.

## Definition 3.2.4.

(1) Let $M$ be an embedded surface in $\mathbb{R}^{3}$. A map $\alpha:(a, b) \rightarrow M$ is called a smooth curve if $\pi_{i} \circ \alpha:(a, b) \rightarrow \mathbb{R}$ is smooth for all $1 \leq i \leq 3$, where $\pi_{i}: \mathbb{R}^{3} \rightarrow \mathbb{R}$ is the projection of $\mathbb{R}^{3}$ to the $i$-th component.
(2) Let $M$ and $N$ be embedded surfaces in $\mathbb{R}^{3}$. A map $f: M \rightarrow N$ is differentiable at $p_{0} \in M$ if there exist a local chart $\phi_{\alpha}: \mathcal{O}_{\alpha} \rightarrow$ $\mathcal{U}_{\alpha}$ on $M$, a local chart $\psi_{\beta}: \mathcal{O}_{\beta} \rightarrow \mathcal{V}_{\beta}$ on $N$ such that $p_{0} \in \mathcal{U}_{\alpha}$, $f\left(p_{0}\right) \in \mathcal{V}_{\beta}$, and $\psi_{\beta} \circ f \circ \phi_{\alpha}^{-1}$ is differentiable at $x_{0}=\phi_{\alpha}^{-1}\left(p_{0}\right)$.

## Tangent plane

Let $M$ be an embedded surface in $\mathbb{R}^{3}$, and $p_{0} \in M$. If $\alpha:(-\delta, \delta) \rightarrow$ $M$ is a smooth curve with $\alpha(0)=p_{0}$, then $\alpha$ can be viewed as a map from $(-\delta, \delta)$ to $\mathbb{R}^{3}$. So we can compute $\alpha^{\prime}(0)$ as in calculus. The space

$$
T M_{p_{0}}=\left\{\alpha^{\prime}(0) \mid \alpha:(-\delta, \delta) \rightarrow M \text { is smooth with } \alpha(0)=p_{0}\right\}
$$

is a vector space, and is called the tangent plane of $M$ at $p_{0}$.

## Exercise 3.2.5.

If $f: \mathcal{O} \rightarrow U \subset M$ is a local coordinate system of $M$, then $f_{x_{1}}(q), f_{x_{2}}(q)$ spans $T M_{f(q)}$ and $\frac{f_{x_{1}} \times f_{x_{2}}}{\left\|f f_{x_{1}} \times f_{x_{2}}\right\|}(q)$ is a unit vector normal to $T M_{f(q)}$.
Prove that the definition of $f: M \rightarrow N$ being differentiable at $p_{0}$ is independent of the choice of local charts.

## The differential of a smooth map

Suppose $M$ and $N$ are two embedded surfaces in $\mathbb{R}^{3}$, and $f: M \rightarrow N$ is a smooth map. The differential of $f$ at $p, d f_{p}$, is the linear map from $T M_{p}$ to $T N_{f(p)}$ defined by

$$
d f_{p}\left(\alpha^{\prime}(0)\right)=(f \circ \alpha)^{\prime}(0),
$$

where $\alpha:(-\delta, \delta) \rightarrow M$ is a smooth curve on $M$ with $\alpha(0)=p$. It follows from the Chain Rule that $d f_{p}$ is well-defined, i.e., if $\alpha, \beta$ are two curves on $M$ so that $\alpha(0)=\beta(0)=p$ and $\alpha^{\prime}(0)=\beta^{\prime}(0)$, then $(f \circ \alpha)^{\prime}(0)=(f \circ \beta)^{\prime}(0)$.

## Smooth vector fields

Let $U$ be an open subset of $M$. A map $\xi: U \rightarrow \mathbb{R}^{3}$ is called a smooth vector field of $M$ defined on $M$ if
(1) $\xi(p) \in T M_{p}$ for each $p \in U$,
(2) given any local coordinate system $\phi_{\alpha}: \mathcal{O}_{\alpha} \rightarrow U_{\alpha}$ of $M$ there exist smooth function $\xi_{1}, \xi_{2}$ on $\phi_{\alpha}^{-1}\left(U \cap U_{\alpha}\right)$ so that

$$
\xi(f(q))=\xi_{1}(q)\left(\phi_{\alpha}\right)_{x_{1}}(q)+\xi_{2}(q)\left(\phi_{\alpha}\right)_{x_{2}}(q)
$$

for all $q \in U \cap U_{\alpha}$.
We will sometimes write a shorthand for $\xi$ in local coordinate system $\phi_{\alpha}$ by

$$
\xi=\xi_{1}\left(\phi_{\alpha}\right)_{x_{1}}+\xi_{2}\left(\phi_{\alpha}\right)_{x_{2}} .
$$

If $f: \mathcal{O} \rightarrow U \subset M$ is a local coordinate system, then we often let $\frac{\partial}{\partial x_{i}}$ to denote the vector field $f_{x_{i}} \circ f^{-1}$, i.e.,

$$
\frac{\partial}{\partial x_{i}}(p):=f_{x_{i}}\left(f^{-1}(p)\right) .
$$

Then the above vector field $\xi$ can be written as $\xi_{1} \frac{\partial}{\partial x_{1}}+\xi_{2} \frac{\partial}{\partial x_{2}}$.

## Smooth differential 1-forms

Let $f: \mathcal{O} \rightarrow U \subset M$ be a local coordinate system on the embedded surface $M$ of $\mathbb{R}^{3}$. Let $\left(d x_{1}\right)_{p},\left(d x_{2}\right)_{p}$ denote the base of $T M_{p}^{*}$ dual to the base $f_{x_{1}}(p), f_{x_{2}}(p)$ of $T M_{p}$, i.e., $d x_{i}\left(f_{x_{j}}\right)=\delta_{i j}$. Let $\left(d x_{1} \wedge d x_{2}\right)_{p}=$ $\left(d x_{1}\right)_{p} \wedge\left(d x_{2}\right)_{p}$.

A smooth differential 1-form $\theta$ on an open subset $U$ of $M$ is an assignment $\theta(p) \in T M_{p}^{*}$ for each $p \in U$ satisfying the following condition:
given any local coordinate system $\phi_{\alpha}: \mathcal{O}_{\alpha} \rightarrow U_{\alpha} \subset M$, there exist smooth $A_{1}, A_{2}: \phi_{\alpha}^{-1}\left(U \cap U_{\alpha}\right) \rightarrow \mathbb{R}$ so that

$$
\theta_{f(q)}=A_{1}(q)\left(d x_{1}\right)_{f(q)}+A_{2}(q)\left(d x_{2}\right)_{f(q)}
$$

for all $q \in \phi_{\alpha}^{-1}\left(U \cap U_{\alpha}\right)$. We often write the above equation in local coordinate system $\phi_{\alpha}$ as

$$
\theta=A_{1} d x_{1}+A_{2} d x_{2}
$$

If $h: M \rightarrow \mathbb{R}$ is a smooth map, then $d h_{p}: T M_{p} \rightarrow T \mathbb{R}_{h(p)}=\mathbb{R}$. So we can identify $d h_{p} \in T M_{p}^{*}$ and $d h$ as a smooth differential 1-form.

## Differential 2-forms

A smooth differential 2 -form $\omega$ on an open subset $U$ of $M$ is an assignment $p \in U \mapsto \omega_{p} \in \wedge^{2} T M_{p}^{*}$ satisfying the condition that given any local coordinate system $f: \mathcal{O} \rightarrow U_{0} \subset M$ there exists a smooth function $h: f^{-1}\left(U \cap U_{0}\right) \rightarrow \mathbb{R}$ such that $\omega$ can be written as

$$
\omega(f(q))=h(q)\left(d x_{1}\right)_{f(q)} \wedge\left(d x_{2}\right)_{f(q)}
$$

for all $q \in f^{-1}\left(U \cap U_{0}\right)$. We will use the notation

$$
\omega=h d x_{1} \wedge d x_{2}
$$

on $U \cap U_{0}$ with respect to the local coordinate system $f$.

## Wedge products

Let $\theta_{1}, \theta_{2}$ be two smooth 1 -forms on $M$. The wedge product $\theta_{1} \wedge \theta_{2}$ is defined by

$$
\theta_{1} \wedge \theta_{2}\left(v_{1}, v_{2}\right)=\theta_{1}\left(v_{1}\right) \theta_{2}\left(v_{2}\right)-\theta_{2}\left(v_{1}\right) \theta_{1}\left(v_{2}\right)
$$

for all $v_{1}, v_{2} \in T M_{p}$ and $p \in M$. It follows from the definition that if $\theta_{1}, \theta_{2}$ are 1 -forms, then

$$
\theta_{1} \wedge \theta_{2}=-\theta_{2} \wedge \theta_{1}
$$

## Exterior differentiation

If 1 -form $\theta$ in local coordinates is $\theta=h_{1} d x_{1}+h_{2} d x_{2}$, then the exterior differentiation of $\theta$ is defined to be

$$
d \theta=\left(-\left(h_{1}\right)_{x_{2}}+\left(h_{2}\right)_{x_{1}}\right) d x_{1} \wedge d x_{2} .
$$

A 1-form $\theta$ is called exact if there exists a smooth function $h$ on $M$ so that $\theta=d h$. A 1-form $\theta$ is called closed if $d \theta=0$.

Next we prove the Poincaré Lemma.

Proposition 3.2.6. Let $\theta$ be a 1 -form on an embedded surface $M$ in $\mathbb{R}^{3}$. Then
(i) $\theta$ is exact implies $\theta$ is closed,
(ii) if $\theta$ is closed, then given any $p \in M$, there exist an open subset $U$ containing $p$ and a smooth function $h: U \rightarrow \mathbb{R}$ such that $\theta=d h$ on $U$, i.e., $\theta$ is locally exact.

Proof. Let $h: M \rightarrow \mathbb{R}$ be a smooth function, then by definition of $d$,

$$
d(d h)=d\left(h_{x_{1}} d x_{1}+h_{x_{2}} d x_{2}\right)=\left(-\left(h_{x_{1}}\right)_{x_{2}}+\left(h_{x_{2}}\right)_{x_{1}}\right) d x_{1} \wedge d x_{2} .
$$

So $d h$ is closed. This proves (i).
Let $\theta$ be a smooth 1-form. Given a local coordinate system $f: \mathcal{O} \rightarrow$ $U_{0}$ around $p=f\left(x_{1}^{0}, x_{2}^{0}\right)$, we can write $\theta=A\left(x_{1}, x_{2}\right) d x_{1}+B\left(x_{1}, x_{2}\right) d x_{2}$ with some smooth function $A, B: \mathcal{O} \rightarrow \mathbb{R}$. The closeness of $\theta$ means $A_{x_{2}}=B_{x_{1}}$. This implies that there exists a small open subset $\mathcal{O}_{0}$ of $\mathcal{O}$ containing $\left(x_{1}^{0}, x_{2}^{0}\right)$ so that

$$
h_{x_{1}}=A, \quad h_{x_{2}}=B
$$

So $\theta=d h$ on $U_{0}=f\left(\mathcal{O}_{0}\right)$.
It follows from definitions of $d$ and wedge product that:
Proposition 3.2.7. If $\theta$ is a 1 -form and $h: M \rightarrow \mathbb{R}$ smooth, then

$$
d(h \theta)=d h \wedge \theta+h d \theta=d(\theta h)=h d \theta-\theta \wedge d h .
$$

## Exercise 3.2.8.

The definition of $d$ is independent of the choice of local coordiante systems. In other words, if $\theta=h_{1} d x_{1}+h_{2} d x_{2}$ and $\theta=b_{1} d y_{1}+b_{2} d y_{2}$ with respect to local coordinate systems $\phi_{\alpha}$ and $\phi_{\beta}$ respectively, then

$$
\left(-\left(h_{1}\right)_{x_{2}}+\left(h_{2}\right)_{x_{1}}\right) d x_{1} \wedge d x_{2}=\left(-\left(b_{1}\right)_{y_{2}}+\left(b_{2}\right)_{y_{1}}\right) d y_{1} \wedge d y_{2} .
$$

## Smooth bilinear forms

A smooth bilinear form $b$ on an open subset $U$ of $M$ is a collection of bilinear forms $b_{p}: T M_{p} \times T M_{p} \rightarrow \mathbb{R}$ with $p \in U$ satisfying the condition that given any local coordinate system $f: \mathcal{O} \rightarrow U_{0}$ on $M$ there exist smooth functions $b_{i j}$ on $f^{-1}\left(U \cap U_{0}\right)$ so that

$$
b_{f(q)}=\sum_{i, j=1}^{2} b_{i j}(q)\left(d x_{i}\right)_{f(q)} \otimes\left(d x_{j}\right)_{f(q)}
$$

We will write $b$ locally as $\sum_{i, j=1}^{2} b_{i j}(x) d x_{i} \otimes d x_{j}$. If $b_{p}$ is symmetric bilinear for each $p \in U$, then $b_{i j}=b_{j i}$. When $b$ is symmetric, we will skip the $\otimes$ and simply write $b$ as $\sum_{i, j=1}^{2} b_{i j}(x) d x_{i} d x_{j}$.

### 3.3. The first and second fundamental forms.

Let $M$ be an embedded surface in $\mathbb{R}^{3}$, and $f: \mathcal{O} \rightarrow U$ a local coordinate system on $M$. Then $f_{x_{1}}, f_{x_{2}}$ form a base of $T M_{f(x)}$ for all $x \in \mathcal{O}$ and

$$
N=\frac{f_{x_{1}} \times f_{x_{2}}}{\sqrt{\left\|f_{x_{1}} \times f_{x_{2}}\right\|}} .
$$

is the unit vector that is normal to $T M_{f\left(x_{1}, x_{2}\right)}$.
We will assciate to each $p \in M$ two smooth symmetric bilinear forms, the first and second fundamental forms of $M$. The first fundamental form of $M$ at $p \in M$ denoted by

$$
\mathrm{I}_{p}: T M_{p} \times T M_{p} \rightarrow \mathbb{R}
$$

is the bilinear functional defined by $\mathrm{I}_{p}(u, v)=u \cdot v$, the dot product of $u, v$ in $\mathbb{R}^{3}$. Then $I_{p}$ defines an inner product on $T M_{p}$. Let

$$
g_{11}=f_{x_{1}} \cdot f_{x_{1}}, \quad g_{12}=g_{21}=f_{x_{1}} \cdot f_{x_{2}}, \quad g_{21}=f_{x_{2}} \cdot f_{x_{2}} .
$$

Then

$$
\begin{aligned}
\mathrm{I} & =g_{11} d x_{1} \otimes d x_{2}+g_{12}\left(d x_{1} \otimes d x_{2}+d x_{2} \otimes d x_{1}\right)+g_{22} d x_{2} \otimes d x_{2} \\
& =g_{11} d x_{1}^{2}+2 g_{12} d x_{1} d x_{2}+g_{22} d x_{2}^{2} .
\end{aligned}
$$

Since $g_{i j}$ are smooth functions, I is a smooth symmetric bilinear form on $M$.

The unit normal vector field $N$ is a map from $M$ to $S^{2}$, which is called the Gauss map of $M$. Recall that the differential of $N$ at $p$ is the linear map from $T M_{p}$ to $T\left(S^{2}\right)_{N(p)}$ defined by

$$
d N_{p}(u)=(N \circ \alpha)^{\prime}(0),
$$

where $\alpha:(-\epsilon, \epsilon) \rightarrow M$ is a smooth curve with $\alpha(0)=p$ and $\alpha^{\prime}(0)=u$. It follows from the definition of $d N_{p}$ that

$$
d N\left(f_{x_{1}}\right)=N_{x_{1}}, \quad d N\left(f_{x_{2}}\right)=N_{x_{2}} .
$$

Since the tangent plane of $S^{2}$ at a point $q \in S^{2}$ is the plane perpendicular to $q$. We see that both $T S_{N(p)}^{2}$ and $T M_{p}$ are perpendicular to $N_{p}$, so they are the same plane. Hence $d N_{p}$ can be viewed as a linear operator from $T M_{p}$ to $T M_{p}$.
ab Proposition 3.3.1. The linear map $d N_{p}: T M_{p} \rightarrow T M_{p}$ is self-adjoint.

Proof. Use $f_{x_{i}} \cdot N=0$ to compute

$$
d N_{p}\left(f_{x_{i}}\right) \cdot f_{x_{j}}=N_{x_{i}} \cdot f_{x_{j}}=\left(N \cdot f_{x_{j}}\right)_{x_{i}}-N \cdot f_{x_{j} x_{i}}=-N \cdot f_{x_{i} x_{j}}
$$

Similarly,

$$
f_{x_{i}} \cdot d N_{p}\left(f_{x_{j}}\right)=-N \cdot f_{x_{i} x_{j}}
$$

This porves that $d N_{p}\left(f_{x_{i}}\right) \cdot f_{x_{j}}=f_{x_{i}} \cdot d N_{p}\left(f_{x_{j}}\right)=-N \cdot f_{x_{i} x_{j}}$.
The self-adjoint operator $-d N_{p}: T M_{p} \rightarrow T M_{p}$ is called the shape operator of the surface $M$ at $p$. The symmetric bilinear form $\mathrm{II}_{p}$ associated to the self-adjoint map $-d N_{p}$ is called the second fundamental form of the surface, i.e.,

$$
\operatorname{II}\left(f_{x_{i}}, f_{x_{j}}\right)=-d N\left(f_{x_{i}}\right) \cdot f_{x_{j}}=N \cdot f_{x_{i} x_{j}}
$$

We can write the second fundmental form II in terms of the dual basis $d x_{1}, d x_{2}$ :

$$
\mathrm{II}=\ell_{11} d x_{1}^{2}+2 \ell_{12} d x_{1} d x_{2}+\ell_{22} d x_{2}^{2},
$$

where

$$
\ell_{i j}=-N_{x_{i}} \cdot f_{x_{j}}=N \cdot f_{x_{i} x_{j}}
$$

We claim that the first and second fundamental forms are invariant under rigid motions. To see this, let $C(x)=A x+v_{0}$ be a rigid motion of $\mathbb{R}^{3}$, i.e., $A$ is an orthogonal transformation and $v_{0}$ is a constant vector in $\mathbb{R}^{3}$. Let $f: \mathcal{O} \rightarrow U \subset M$ is a local coordinate system of the embedded surface $M$ in $\mathbb{R}^{3}$. Then $\hat{f}=C \circ f$ is a local coordinate system for the new surface $C(M), \hat{f}_{x_{i}}=A f_{x_{i}}$, and $\hat{N}=A N$ is normal to $C(M)$. Since $A$ is an orthogonal transformation,

$$
\left(A v_{1}\right) \cdot\left(A v_{2}\right)=v_{1} \cdot v_{2}
$$

Let $\hat{I}$ and III denote the fundamental forms of $\hat{M}=C(M)$. Then

$$
\begin{aligned}
\hat{g}_{i j} & =\hat{\mathrm{I}}\left(\hat{f}_{x_{i}}, \hat{f}_{x_{j}}\right)=\hat{f}_{x_{i}} \cdot \hat{f}_{x_{j}}=A f_{x_{i}} \cdot A f_{x_{j}} \\
& =f_{x_{i}} \cdot f_{x_{j}}=\mathrm{I}\left(f_{x_{i}}, f_{x_{j}}\right)=g_{i j} .
\end{aligned}
$$

So $\hat{I}=$ I. Similarly, $\hat{I} I=I$. This proves our claim.

## bo Example 3.3.2. Fundamental forms of a graph

Let $u: \mathcal{O} \rightarrow \mathbb{R}$ be a smooth function, and $f: \mathcal{O} \rightarrow \mathbb{R}^{3}$ the graph of $u$, i.e., $f$ is defined by

$$
f\left(x_{1}, x_{2}\right)=\left(x_{1}, x_{2}, u\left(x_{1}, x_{2}\right)\right),
$$

i.e., $M=f(\mathcal{O})$ is the graph of $u$. Then

$$
\begin{aligned}
& f_{x_{1}}=\left(1,0, u_{x_{1}}\right), \quad f_{x_{2}}=\left(0,1, u_{x_{2}}\right), \\
& f_{x_{1} x_{1}}=\left(0,0, u_{x_{1} x_{1}}\right), \quad f_{x_{1} x_{2}}=\left(0,0, u_{x_{1} x_{2}}\right), \quad f_{x_{2} x_{2}}=\left(0,0, u_{x_{2} x_{2}}\right), \\
& N=\frac{\left(-u_{x_{1}},-u_{x_{2}}, 1\right)}{\sqrt{u_{x_{1}}^{2}+u_{x_{2}}^{2}+1}}
\end{aligned}
$$

A direct computation gives

$$
\begin{aligned}
& g_{11}=f_{x_{1}} \cdot f_{x_{1}}=1+u_{x_{1}}^{2}, \\
& g_{12}=f_{x_{1}} \cdot f_{x_{2}}=u_{x_{1}} u_{x_{2}}, \\
& g_{22}=f_{x_{2}} \cdot f_{x_{2}}=1+u_{x_{2}}^{2}, \\
& h_{11}=N \cdot f_{x_{1} x_{1}}=\frac{u_{x_{1} x_{1}}}{\sqrt{u_{x_{1}}^{2}+u_{x_{2}}^{2}+1}}, \\
& h_{12}=N \cdot f_{x_{1} x_{2}}=\frac{u_{x_{1} x_{2}}}{\sqrt{u_{x_{1}}^{2}+u_{x_{2}}^{2}+1}}, \\
& h_{22}=N \cdot f_{x_{2} x_{2}}=\frac{u_{x_{2} x_{2}}}{\sqrt{u_{x_{1}}^{2}+u_{x_{2}}^{2}+1}} .
\end{aligned}
$$

So the two fundamental forms for $f(\mathcal{O})$ are

$$
\begin{aligned}
& \mathrm{I}=\left(1+u_{x_{1}}^{2}\right) d x_{1}^{2}+2 u_{x_{1}} u_{x_{2}} d x_{1} d x_{2}+\left(1+u_{x_{2}}^{2}\right) d x_{2}^{2}, \\
& \mathrm{II}=\frac{1}{\sqrt{1+u_{x_{1}}^{2}+u_{x_{2}}^{2}}}\left(u_{x_{1} x_{1}} d x_{1}^{2}+2 u_{x_{1} x_{2}} d x_{1} d x_{2}+u_{x_{2} x_{2}} d x_{2}^{2}\right) .
\end{aligned}
$$

Proposition 3.3.3. Locally a surface is a graph.
Proof. Let $f=\left(f_{1}, f_{2}, f_{3}\right): U \rightarrow \mathbb{R}^{3}$ be a surface. Given $p=\left(x_{0}, y_{0}\right) \in$ $U$, since $f_{x_{1}}, f_{x_{2}}$ are linearly independent,

$$
\operatorname{det}\left(\begin{array}{ccc}
i & j & k \\
\left(f_{1}\right) x_{1} & \left(f_{2}\right) x_{1} & \left(f_{3}\right) x_{1} \\
\left(f_{1}\right)_{x_{2}} & \left(f_{2}\right) x_{2} & \left(f_{3}\right) x_{3}
\end{array}\right) \quad(p) \neq 0 .
$$

So one of the following $2 \times 2$ determinants at $p$ must be non-zero:

$$
\left(\begin{array}{ll}
\left(f_{1}\right)_{x_{1}} & \left(f_{2}\right)_{x_{1}} \\
\left(f_{1}\right)_{x_{2}} & \left(f_{2}\right)_{x_{2}}
\end{array}\right), \quad\left(\begin{array}{ll}
\left(f_{1}\right)_{x_{1}} & \left(f_{3}\right)_{x_{1}} \\
\left(f_{1}\right)_{x_{2}} & \left(f_{3}\right)_{x_{2}}
\end{array}\right), \quad\left(\begin{array}{ll}
\left(f_{2}\right)_{x_{1}} & \left(f_{3}\right)_{x_{1}} \\
\left(f_{2}\right)_{x_{2}} & \left(f_{3}\right)_{x_{2}}
\end{array}\right) .
$$

Suppose

$$
\operatorname{det}\left(\begin{array}{ll}
\left(f_{1}\right)_{x_{1}} & \left(f_{2}\right)_{x_{1}} \\
\left(f_{1}\right)_{x_{2}} & \left(f_{2}\right)_{x_{2}}
\end{array}\right)(p) \neq 0
$$

Let $\pi_{12}: \mathbb{R}^{3} \rightarrow \mathbb{R}^{2}$ denote the projection of $\mathbb{R}^{3}$ to the first two coordinates. Then the differential of the map $\pi_{12} \circ f: U \rightarrow \mathbb{R}^{2}$ at $p$ is given by the matrix

$$
\left(\begin{array}{ll}
\left(f_{1}\right)_{x_{1}} & \left(f_{2}\right)_{x_{1}} \\
\left(f_{1}\right)_{x_{2}} & \left(f_{2}\right)_{x_{2}}
\end{array}\right)
$$

at $p$, which is non-singular. By the Inverse Function Theorem, there exist open subsets $\mathcal{O}_{1}, \mathcal{O}_{2}$ containing $p$ and $\left(f_{1}(p), f_{2}(p)\right)$ respectively so that $\pi_{12} \circ f$ maps $\mathcal{O}_{1}$ diffeomorphically to $\mathcal{O}_{2}$. Let $h: \mathcal{O}_{2} \rightarrow \mathcal{O}_{1}$ denote the inverse of $\pi_{12} \circ f: \mathcal{O}_{1} \rightarrow \mathcal{O}_{2}$, and $\tilde{f}\left(y_{1}, y_{2}\right)=f \circ h\left(y_{1}, y_{2}\right)$. In other words, we have changed coordinates $\left(x_{1}, x_{2}\right)$ to $\left(y_{1}, y_{2}\right)$ and expressed $\left(x_{1}, x_{2}\right) \in \mathcal{O}_{1}$ in terms of $\left(y_{1}, y_{2}\right) \in \mathcal{O}_{2}$. Hence $f_{3}\left(x_{1}, x_{2}\right)=f_{3}\left(h\left(y_{1}, y_{2}\right)\right)$ is a function of $y_{1}$ and $y_{2}$. So we see that

$$
\tilde{f}\left(y_{1}, y_{2}\right)=\left(y_{1}, y_{2}, f_{3}\left(h\left(y_{1}, y_{2}\right)\right)\right)
$$

is a graph of $u=f_{3} \circ h$ and $\tilde{f}$ is a local parametrization of $M$.

## Angle and arc length

Since $\mathrm{I}_{p}$ gives an inner product on $T M_{p}$ for all $p \in M$, we can use I to compute the arc length of curves on $M$ and the angle between two curves on $M$ :
(1) if $\alpha_{1}, \alpha_{2}:(-\epsilon, \epsilon) \rightarrow M$ are two smooth curves such that $\alpha_{1}(0)=$ $\alpha_{2}(0)$, then the angle between $\alpha_{1}$ and $\alpha_{2}$ is defined to be the angle between $\alpha_{1}^{\prime}(0)$ and $\alpha_{2}^{\prime}(0)$,
(2) the arc length of a curve on $M$ given by $\alpha(t)=f\left(x_{1}(t), x_{2}(t)\right)$ with $t \in[a, b]$ is

$$
\begin{aligned}
\int_{a}^{b}\left\|\alpha^{\prime}(t)\right\| d t & =\int_{a}^{b}\left\|f_{x_{1}} x_{1}^{\prime}+f_{x_{2}} x_{2}^{\prime}\right\| d t \\
& =\int_{a}^{b}\left(\sum_{i, j=1}^{2}\left(f_{x_{i}} \cdot f_{x_{j}}\right) x_{i}^{\prime} x_{j}^{\prime}\right)^{\frac{1}{2}} d t \\
& =\int_{a}^{b}\left(\sum_{i, j=1}^{2} g_{i j}\left(x_{1}(t), x_{2}(t)\right) x_{i}^{\prime}(t) x_{j}^{\prime}(t)\right)^{\frac{1}{2}} d t
\end{aligned}
$$

Recall that the area of the parallelgram spanned by $v_{1}, v_{2}$ in $\mathbb{R}^{3}$ is equal to

$$
\left\|v_{1}\right\|\left\|v_{1}\right\||\cos \theta|=\left\|v_{1} \times v_{2}\right\|=\left(\operatorname{det}\left(\begin{array}{ll}
v_{1} \cdot v_{1} & v_{1} \cdot v_{2} \\
v_{1} \cdot v_{2} & v_{2} \cdot v_{2}
\end{array}\right)\right)^{\frac{1}{2}}
$$

where $\theta$ is the angle between $v_{1}$ and $v_{2}$.
bm Proposition 3.3.4. Let $f: \mathcal{O} \rightarrow U \subset M$ be a local coordinate system of an embedded surface, and $g_{i j}=f_{x_{i}} \cdots f_{x_{j}}$ the coefficients of the first fundamental form. Then the area of $U$ is

$$
\operatorname{Area}(U)=\int_{\mathcal{O}}\left\|f_{x_{1}} \cdot f_{x_{2}}\right\| d x_{1} d x_{2}=\int_{\mathcal{O}} \sqrt{\operatorname{det}\left(g_{i j}\right)} d x_{1} d x_{2}
$$

The second fundamental form has a very simple geometric meaning:
dc Proposition 3.3.5. Let $v$ be a unit tangent vector of $M=f(\mathcal{O})$ at $p \in M=f(\mathcal{O})$, and $\alpha$ the intersection curve of $M$ and the plane through $p$ spanned by $v$ and the normal vector $N_{p}$ of $M$ at $p$. Then $\mathrm{II}(v, v)$ is the curvature of the plane curve $\alpha$ at $p$.

Proof. Write $\alpha(s)=f\left(x_{1}(s), x_{2}(s)\right)$ with arc length parameter. Then

$$
\alpha^{\prime \prime}(s)=\sum_{i j=1}^{2} f_{x_{i} x_{j}} x_{i}^{\prime} x_{j}^{\prime} .
$$

Since $N(p)$ is normal to $M$ at $p$ and $\alpha^{\prime}(0)=v, N(p)$ is normal to $\alpha^{\prime}(0)$. But the curvature of the plane curve $\alpha$ at $s=0$ is $\alpha^{\prime \prime}(0) \cdot N(p)$, and

$$
\alpha^{\prime \prime}(0) \cdot N=\sum_{i, j=1}^{2}\left(f_{x_{i} x_{j}} \cdot N\right) x_{i}^{\prime} x_{j}^{\prime}=\sum_{i j=1}^{2} h_{i j} x_{i}^{\prime} x_{j}^{\prime}=\operatorname{II}\left(\alpha^{\prime}(0), \alpha^{\prime}(0)\right)
$$

### 3.4. The Gaussian and mean curvature.

Definition 3.4.1. The eigenvalues $\lambda_{1}, \lambda_{2}$ of the shape operator $-d N_{p}$ are called the principal curvatures of $M$ at $p$, the unit eigenvectors of $-d N_{p}$ are called the principal directions,

$$
K(p)=\operatorname{det}\left(-d N_{p}\right)=\lambda_{1} \lambda_{2}, \quad H(p)=\frac{1}{2} \operatorname{tr}\left(-d N_{p}\right)=\frac{1}{2}\left(\lambda_{1}+\lambda_{2}\right),
$$

are called the Gaussian, and the mean curvature of $M$ at $p$ respectively.
do Proposition 3.4.2. Let $f: \mathcal{O} \rightarrow U$ be a local coordinate system of an embedded surface $M$ in $\mathbb{R}^{3}$, and $g_{i j}$, $\ell_{i j}$ the coefficients of I, II on $U$ respectively. Then

$$
\begin{aligned}
K & =\frac{\operatorname{det}\left(\ell_{i j}\right)}{\operatorname{det}\left(g_{i j}\right)}=\frac{\ell_{11} \ell_{22}-\ell_{12}^{2}}{g_{11} g_{22}-g_{12}^{2}}, \\
H & =\frac{g_{22} \ell_{11}-2 g_{12} \ell_{12}+g_{11} \ell_{22}}{g_{11} g_{22}-g_{12}^{2}}
\end{aligned}
$$

Proof. By Proposition ªc $_{3.1 .4,}^{\text {ac }} K=\operatorname{det}\left(G^{-1} L\right)=\frac{\operatorname{det}(G)}{\operatorname{det}(L)}$, and $H=$ $\operatorname{tr}\left(G^{-1} L\right)$. The formula for $H$ follows from

$$
G^{-1}=\frac{1}{\operatorname{det}\left(g_{i j}\right)}\left(\begin{array}{cc}
g_{22} & -g_{12} \\
-g_{12} & g_{11}
\end{array}\right)
$$

Suppose $v_{1}, v_{2}$ are unit eigenvectors of $-d N_{p}$ with eigenvalues $\lambda_{1}>$ $\lambda_{2}$, then $\mathrm{II}_{p}(u, u)$ assumes minimum value $\lambda_{2}$ on the unit circle of $T M_{p}$ at $v_{2}$ and maximum value $\lambda_{1}$ at $v_{1}$.

Below we compute the $K$ and $H$ of a plane, a sphere, a cylinder, and a graph.

Cs Example 3.4.3. Let $M$ be a plane. Let us parametrize it as a graph $f\left(x_{1}, x_{2}\right)=\left(x_{1}, x_{2}, a_{1} x_{1}+a_{2} x_{2}\right)$. Then

$$
f_{x_{1}}=\left(1,0, a_{1}\right), \quad f_{x_{2}}=\left(0,1, a_{2}\right), \quad N=\frac{\left(a_{1}, a_{2},-1\right)}{\sqrt{a_{1}^{2}+a_{2}^{2}+1}}
$$

Since $N$ is a constant vector, $N_{x_{i}}=0$. So the shape operator $d N_{p}=0$. Hence $K=0$ and $H=0$. The first and second fundamental forms are

$$
\mathrm{I}=\left(1+a_{1}^{2}\right) d x_{1}^{2}+\left(1+a_{2}^{2}\right) d x_{2}^{2}, \quad \mathrm{II}=0
$$

Example 3.4.4. Let $M$ be the sphere of radius $r$. Let $f\left(x_{1}, x_{2}\right)$ be a parametrization of a piece of the sphere. Note that the unit normal $N$ of $M$ at $f\left(x_{1}, x_{2}\right)$ is $\frac{1}{r} f\left(x_{1}, x_{2}\right)$. Hence $d N\left(f_{x}\right)=\frac{1}{r} f_{x}$ and $d N\left(f_{y}\right)=\frac{1}{r} f_{y}$. This means that $d N_{p}=\frac{1}{r}$ Id. So $K=\frac{1}{r^{2}}$ and $H=-\frac{2}{r}$.
ct Example 3.4.5. Let $M$ be the cylinder parametrized as

$$
f(x, y)=(r \cos x, r \sin x, y)
$$

Then

$$
\begin{aligned}
& f_{x}=(-r \sin x, r \cos x, 0), \\
& f_{y}=(0,0,1) \\
& N=(\cos x, \sin x, 0) .
\end{aligned}
$$

So $N_{x}=\frac{1}{r} f_{x}$ and $N_{y}=0$. This means that $d N_{p}$ has two eigenvalues $\frac{1}{r}$ and 0 . Hence $K=0$ and $H=-\frac{1}{r}$. The first and second fundamental forms are

$$
\mathrm{I}=r^{2} d x_{1}^{2}+d y^{2}, \quad \mathrm{II}=r d x^{2}
$$

Note that both a plane and a cylinder have Gaussian curvature 0 .
ee Example 3.4.6. $K$ and $H$ for a graph
Let $f\left(x_{1}, x_{2}\right)=\left(x_{1}, x_{2}, u\left(x_{1}, x_{2}\right)\right)$ denote the graph of a $u: \mathcal{O} \rightarrow \mathbb{R}$. We have computed the first and second fundamental form in Example 3.3.2. By Proposition 3.4.2, we have

$$
\begin{align*}
& K=\frac{u_{x_{1} x_{1}} u_{x_{2} x_{2}}-u_{x_{1} x_{2}}^{2}}{\left(1+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right)^{2}}  \tag{3.4.1}\\
& H=\frac{\left(1+u_{x_{1}}^{2}\right) u_{x_{2} x_{2}}-2 u_{x_{1}} u_{x_{2}} u_{x_{1} x_{2}}+\left(1+u_{x_{2}}^{2}\right) u_{x_{1} x_{1}}}{\left(1+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right)^{\frac{3}{2}}} \tag{3.4.2}
\end{align*}
$$

aas:1
aas:2

Definition 3.4.7. A point $p$ in a surface $M \subset \mathbb{R}^{3}$ is called an umbilic point if the shape operator at $p$ is equal to $\lambda \mathrm{Id}$ for some scalar $\lambda$, or equivalently, the two prinicipal curvatures are equal at $p$.

Definition 3.4.8. A unit tangent direction $u$ of an embedded surface $M$ in $\mathbb{R}^{3}$ is called asymptotic if $\operatorname{II}(u, u)=0$.

Exercise 3.4.9. (1) If $K\left(p_{0}\right)<0$, then there exist exactly two asymptotic directions.
(2) If $p$ is not an umbilic point of $M$, then there exists a local smooth o.n. frame $v_{1}, v_{2}$ on $M$ near $p$ so that $v_{1}(q), v_{2}(q)$ are principal directions of $M$ at $q$.

Next we give some geometric interpretation of the Gaussian curvature $K$. The first is a consequence of the Taylor expansion:
Proposition 3.4.10. Let $M$ be an embedded surface in $\mathbb{R}^{3}$, and $p \in M$.
(1) If $K(p)>0$, then there is a small neighborhood $\mathcal{O}_{p}$ in $M$ of $p$ such that $\mathcal{O}_{p}$ lies on one side of the tangent plane of $M$ at $p$.
(2) If $K(p)<0$, then given any small open subset of $M$ containing $p, M$ lies on both sides of the tangent plane of $M$ at $p$.
Proof. After rotation and translation, we may assume that $p=(0,0,0)$, $M$ is parametrized as a graph

$$
f\left(x_{1}, x_{2}\right)=\left(x_{1}, x_{2}, u\left(x_{1}, x_{2}\right)\right)
$$

in a neighborhood of $p$, and the unit normal $N_{p}=(0,0,1)$. This means that we assume

$$
u(0,0)=0, \quad u_{x_{i}}(0,0)=0
$$

Assume $K(p)>0$, then $\lambda_{1}(p) \lambda_{2}(p)>0$. So either both $\lambda_{1}, \lambda_{2}$ are bositive or both are negative. So $\operatorname{II}(0,0)$ is definite. Use Example 3.3.2 formula (3.4.1) to conclude that $\left(u_{x_{i} x_{j}}(0,0)\right)$ is positive definite or negative definite. Suppose it is positive definite. Then the quadratic
form $\sum_{i j} u_{x_{i} x_{j}}(0,0) y_{i} y_{j}>0$ for all $y \in \mathbb{R}^{2}$. On the other hand, the Taylor expanison of $u$ implies that

$$
\begin{aligned}
& u\left(x_{1}, x_{2}\right) \\
& =u(0,0)+u_{x_{1}}(0,0) x_{1}+u_{x_{2}}(0,0) x_{2}+\frac{1}{2} \sum_{i j} u_{x_{i} x_{j}}(0,0) x_{i} x_{j}+o\left(r^{2}\right) \\
& =\frac{1}{2} \sum_{i j} u_{x_{i} x_{j}}(0,0) x_{i} x_{j}+o\left(r^{2}\right)
\end{aligned}
$$

where $r^{2}=\sum_{i} x_{i}^{2}$. So there exists $\delta>0$ so that if $0<\|x\|<\delta$, then $u\left(x_{1}, x_{2}\right)>0$. This proves (1). Statement (2) can be proved in a similar manner.

Proposition 3.4.11. Let $f: \mathcal{O} \rightarrow U \subset M$ be a local coordinate system of an embedded surface $M$ in $\mathbb{R}^{3}$, $N$ the unit normal field, and $p=$ $f(q) \in f(\mathcal{O})$. Then

$$
K(p)=\lim _{\Omega \rightarrow q} \frac{\operatorname{Area}(N(\Omega))}{\operatorname{Area}(f(\Omega))}
$$

Proof. By Proposition ${ }^{\mathrm{bm}} 3.3 .4$ we have

$$
\begin{aligned}
\operatorname{Area}(f(\Omega)) & =\int_{W}\left\|f_{x} \times f_{y}\right\| d x d y \\
\operatorname{Area}(N(\Omega)) & =\int_{W}\left\|N_{x} \times N_{y}\right\| d x d y
\end{aligned}
$$

Since $d N\left(f_{x}\right)=N_{x}$ and $d N\left(f_{y}\right)=N_{y}$, we have

$$
\left\|N_{x} \times N_{y}\right\|=|\operatorname{det}(d N)|\left\|f_{x} \times f_{y}\right\|=|K|\left\|f_{x} \times f_{y}\right\|
$$

Proposition follows from the fact that the Gaussian curvature $K$ is a smooth function.

To understand the meaning of mean curvature $H$ we need calculus of variations.

## 4. Calculus of Variations

Let $\Omega$ be an open subset of $\mathbb{R}^{n}$ with smooth boundary such that the closure $\bar{\Omega}$ is compact, and $\partial \Omega$ the boundary of $\Omega$. Let $\alpha: \partial \Omega \rightarrow \mathbb{R}^{m}$ be a smooth map, and $C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$ the space of smooth maps $u$ from $\Omega$ to $\mathbb{R}^{m}$ so that $u \mid \partial \Omega=\alpha$. Let $L: \Omega \times\left(\mathbb{R}^{m}\right)^{n+1} \rightarrow \mathbb{R}$ be a smooth map, and $J: C_{\alpha}^{\infty}(\Omega, \mathbb{R}) \rightarrow \mathbb{R}$ the function defined by

$$
J(u)=\int_{\Omega} L\left(x, u, u_{x_{1}}, \ldots, u_{x_{n}}\right) d x_{1} \cdots d x_{n}
$$

The function $L$ is called a Lagrangian, and $J$ the functional defined by $L$. The calculus of variations studies the condition on $u$ when $u$ is a "critical point" of $L$. Since $C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$ is an infinite dimensional vector space, we need to extend the concept of critical point to this function space.

Recall that the following statements are equivalent for a smooth function $f: \mathbb{R}^{n} \rightarrow \mathbb{R}$ :
(1) $p$ is a critical point of $f$,
(2) $f_{x_{1}}(p)=\cdots=f_{x_{n}}(p)=0$,
(3) $\left.\frac{d}{d t}\right|_{t=0} f(\alpha(t))=0$ for all smooth curve $\alpha$ through $p$.

We will use (3) to define critical points for the functional $J$ defined by the Lagrangian $L$. Let $u \in C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$.

Definition 4.0.12. A smooth variation of $u \in C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$ is a curve $\beta:(-\epsilon, \epsilon) \rightarrow C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$ so that $\beta(0)=u$ and $(s, x) \mapsto \beta(s)(x)$ is smooth. Note that $v(x)=\left.\frac{\partial}{\partial s}\right|_{s=0} \beta(s)(x)$ is a smooth map from $\Omega$ to $\mathbb{R}^{m}$ that vanishes on $\partial \Omega$.

Definition 4.0.13. $u \in C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$ is a critical point of $J$ if

$$
(J \circ \beta)^{\prime}(0)=0
$$

for all smooth variations $\beta$ of $u$ in $C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$.
ed Proposition 4.0.14. If $u \in C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$ assumes maximum (or minimum) of $J$, then $u$ is a critical point of $J$.

Proof. Let $\beta:(-\epsilon, \epsilon) \rightarrow C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$ be a smooth variation of $u$ in $C_{\alpha}^{\infty}\left(\Omega, \mathbb{R}^{m}\right)$. Then $s=0$ is an exteme point of $J \circ \beta:(-\epsilon, \epsilon) \rightarrow \mathbb{R}$ and $\left.\frac{d}{d s}\right|_{s=0}(J \circ \beta)=0$. This implies that extreme point of $J$ is a critical point of $J$.

To derive the condition for $u$ to be a critical point of the functional $J$, we need a Lemma:
ag Lemma 4.0.15. Given a smooth map $\left(g_{1}, \cdots, g_{m}\right): \Omega \rightarrow \mathbb{R}^{m}$, if

$$
\int_{\Omega} \sum_{j=1}^{n} g_{j}(x) h_{j}(x) d x_{1} \cdots d x_{n}=0
$$

for all smooth $\left(h_{1}, \cdots, h_{n}\right): \Omega \rightarrow \mathbb{R}^{n}$ with $h \mid \partial \Omega=0$, then $g_{j} \equiv 0$ for all $1 \leq j \leq n$.

Proof. We will prove this Lemma when $m=1$ and $n=1$. For general $m, n$, we use the Stoke's Theorem instead of the Fundamental theorem
of calculus. Since in this notes, we only need to deal with the case when $n=1$ or 2 , we will give a proof when $n=2$ in more detail later.

If $g$ is not identically zero, then we may assume there is an interval $\left(\delta_{1}, \delta_{2}\right) \subset[0,1]$ so that $g$ is positive on $\left(\delta_{1}, \delta_{2}\right)$. Choose $\epsilon>0$ so that $\epsilon<\frac{\left(\delta_{2}-\delta_{1}\right)}{2}$. Now choose $h$ to be a non-negative function, which vanishes outside ( $\delta_{1}, \delta_{2}$ ) and is $h>0$ in $\left(\delta_{1}+\epsilon, \delta_{2}-\epsilon\right)$. Then $g h>0$ on $\left(\delta_{1}+\epsilon, \delta_{2}-\epsilon\right.$ ) and is non-negative on $[0,1]$. So $\int_{0}^{1} g(t) h(t)>0$, a contradiction. Hence $g \equiv 0$.

We will derive the condition for critical points of $L$ in the later sections.

### 4.1. Calculus of Variations of one variable.

In this section, we derive the condition for a curve on an embedded surface $M$ in $\mathbb{R}^{3}$ joining $p, q \in M$ that has shortest arc length.

Given $p, q \in \mathbb{R}^{n}$, let $\mathcal{M}_{p, q}$ denote the space of smooth curves $x$ : $[0,1] \rightarrow \mathbb{R}^{n}$ so that $x(0)=p$ and $x(1)=q$. Let $L:[0,1] \times \mathbb{R}^{n} \times \mathbb{R}^{n} \rightarrow \mathbb{R}$ be a smooth function, and $J: \mathcal{M}_{p, q} \rightarrow \mathbb{R}$ the functional defined by

$$
\begin{equation*}
J(x)=\int_{0}^{1} L\left(t, x(t), x^{\prime}(t)\right) d t \tag{4.1.1}
\end{equation*}
$$

Suppose $\gamma$ is a critical point of $J$. Let $\beta$ be a variation of $\gamma$, and $\beta^{\prime}(0)=h$. Let

$$
b(s, t)=\beta(s)(t)
$$

We have

$$
b(0, t)=\gamma(t), \quad h(t)=\frac{\partial b}{\partial s}(0, t), \quad h(0)=h(1)=0 .
$$

By definition of critical point and integration by part, we get

$$
\begin{aligned}
& 0=\left.\frac{d}{d s}\right|_{s=0} J(\beta(s))=\left.\frac{d}{d s}\right|_{s=0} \int_{0}^{1} L\left(t, b(s, t), \frac{\partial b}{\partial t}\right) d t \\
&= \int_{0}^{1} \sum_{j} \frac{\partial L}{\partial x_{j}} \frac{\partial b_{j}}{\partial s}+\left.\frac{\partial L}{\partial y_{j}} \frac{\partial^{2} b}{\partial t \partial s}\right|_{s=0} d t \\
&= \int_{0}^{1} \sum_{j}\left(\frac{\partial L}{\partial x_{j}} h_{j}+\sum_{j} \frac{\partial L}{\partial y_{j}} h_{j}^{\prime}\right) d t, \quad \text { here }^{\prime} \text { means } \frac{\mathrm{d}}{\mathrm{dt}}, \\
&= \int_{0}^{1} \sum_{j} \frac{\partial L}{\partial x_{j}} h_{j} d t+\left.\left(\sum_{j} \frac{\partial L}{\partial y_{j}} h_{j}\right)\right|_{0} ^{1}-\int_{0}^{1}\left(\sum_{j} \frac{\partial L}{\partial y_{j}}\right)^{\prime} h_{j} d t \\
&= \int_{0}^{1} \sum_{j}\left(\frac{\partial L}{\partial x_{j}}\left(t, \gamma(t), \gamma^{\prime}(t)\right)-\left(\frac{\partial L}{\partial y_{j}}\left(t, \gamma(t), \gamma^{\prime}(t)\right)\right)^{\prime}\right) h_{j} d t \\
& \quad+\left.\sum_{j} \frac{\partial L}{\partial y_{j}} h_{j}\right|_{0} ^{1}
\end{aligned}
$$

Since $h(0)=h(1)=0$, the boundary term is zero. Hence we have

$$
\int_{0}^{1} \sum_{j}\left(\frac{\partial L}{\partial x_{j}}-\left(\frac{\partial L}{\partial y_{j}}\right)^{\prime}\right) h_{j} d t=0
$$

for all smooth $h:[0,1] \rightarrow \mathbb{R}^{n}$ with $h(0)=h(1)=0$. By Lemma ${ }^{\text {ag. }} 4.0 .15$, we get

$$
\frac{\partial L}{\partial x_{j}}\left(t, \gamma(t), \gamma^{\prime}(t)\right)-\left(\frac{\partial L}{\partial y_{j}}\left(t, \gamma(t), \gamma^{\prime}(t)\right)\right)^{\prime}=0
$$

So we have proved
Theorem 4.1.1. Let $J: \mathcal{M}_{p, q} \rightarrow \mathbb{R}$ be the functional defined by $L(t, x, y)$, i.e., $J(\gamma)=\int_{0}^{1} L\left(t, \gamma(t), \gamma^{\prime}(t)\right) d t$. Then $\gamma \in \mathcal{M}_{p, q}$ is a critical point of $J$ if and only if

$$
\begin{equation*}
\frac{\partial L}{\partial x_{j}}\left(t, \gamma, \gamma^{\prime}\right)=\left(\frac{\partial L}{\partial y_{j}}\left(t, \gamma, \gamma^{\prime}\right)\right)^{\prime} . \tag{4.1.2}
\end{equation*}
$$

Equation (4.1.2) is called the Euler-Lagrange equation for the functional $J$.
Example 4.1.2. Let $L(t, x, y)=\sum_{i=1}^{n} y_{i}^{2}$, and $E: \mathcal{M}_{p, q} \rightarrow \mathbb{R}$ the functional defined by $L$, i.e.,

$$
E(\gamma)=\int_{0}^{1} \sum_{i=1}^{n}\left(\gamma_{i}^{\prime}(t)\right)^{2} d t
$$

The functional $E$ is called the energy functional. A direct computation implies that Euler-Lagrangian equation for $E$ is $y^{\prime \prime}(t)=0$. So $y(t)$ is linear. Since $y(0)=p$ and $y(1)=q$, we have $y(t)=p+t(q-p)$ a straight line.
xq

### 4.2. Geodesics.

Let $M$ be an embedded surface in $\mathbb{R}^{3}$, and $f: \mathcal{O} \rightarrow U \subset M$ a coordinate system. A smooth curve $\alpha:[0,1] \rightarrow U \subset M$ is given by a smooth curve $\left(x_{1}(t), x_{2}(t)\right)$ in $\mathcal{O}$ such that

$$
\alpha(t)=f\left(x_{1}(t), x_{2}(t)\right)
$$

So $\alpha^{\prime}(t)=\sum_{i} f_{x_{i}} x_{i}^{\prime}$. Since the arc length of $\alpha$ is approximated by

$$
\sum\left\|\alpha^{\prime}(t)\right\| \Delta t
$$

the arc length of $\alpha$ is given $\int_{0}^{1}\left\|\alpha^{\prime}(t)\right\| d t$. But

$$
\begin{aligned}
\left\|\alpha^{\prime}\right\|^{2} & =\alpha^{\prime} \cdot \alpha^{\prime}=\left(\sum_{i} f_{x_{i}} x_{i}^{\prime}\right) \cdot\left(\sum_{j} f_{x_{j}} x_{j}^{\prime}\right)=\sum_{i j} f_{x_{i}} \cdot f_{x_{j}} x_{i}^{\prime} x_{j}^{\prime} \\
& =\sum_{i j} g_{i j}(x) x_{i}^{\prime} x_{j}^{\prime}=\mathrm{I}\left(\alpha^{\prime}, \alpha^{\prime}\right) .
\end{aligned}
$$

So the arc length of $\alpha$ is

$$
\mathcal{L}(x)=\int_{0}^{1} \sqrt{\sum_{i j} g_{i j}(x) x_{i}^{\prime} x_{j}^{\prime}} d t
$$

We want to calculate the Euler Lagrange equation for the energy and arc length functionals on the surface $M=f(\mathcal{O})$ in $\mathbb{R}^{3}$. Let $p=\left(x_{1}^{0}, x_{2}^{0}\right)$, $q=\left(x_{1}^{1}, x_{2}^{1}\right)$, and $\mathcal{M}_{p, q}$ the space of all smooth $x:[0,1] \rightarrow \mathcal{O}$ such that $x(0)=p$ and $x(1)=q$. Suppose the first fundamental form of $M$ is $\mathrm{I}=\sum_{i j} g_{i j}(x) d x_{i} d x_{j}$. If $x:[0,1] \rightarrow \mathcal{O}$ lies in $\mathcal{M}_{p, q}$, then $\gamma=f \circ x$ is a curve on $M$ joining $f(p)$ to $f(q)$ and

$$
\left\|\gamma^{\prime}(t)\right\|^{2}=\sum_{i j} g_{i j}(x(t)) x_{i}^{\prime}(t) x_{j}^{\prime}(t)
$$

Let

$$
\begin{equation*}
L(x, y)=\frac{1}{2} \sum_{i j} g_{i j}(x) y_{i} y_{j} \tag{4.2.1}
\end{equation*}
$$

and $E$ and $\mathcal{L}$ denote the energy and arc length functionals

$$
\begin{aligned}
& E(x)=\int_{0}^{1} L\left(x(t), x^{\prime}(t)\right) d t \\
& \mathcal{L}(x)=\int_{0}^{1} \sqrt{L\left(x(t), x^{\prime}(t)\right)} d t
\end{aligned}
$$

The Euler-Lagrange equation for the functional $E$ and $\mathcal{L}$ are

$$
\begin{gather*}
\frac{\partial L}{\partial x_{k}}=\left(\frac{\partial L}{\partial x_{k}^{\prime}}\right)^{\prime} . \quad k=1,2  \tag{4.2.2}\\
\frac{\frac{\partial L}{\partial x_{k}}}{2 \sqrt{L}}=\left(\frac{\frac{\partial L}{\partial x_{k}^{\prime}}}{2 \sqrt{L}}\right)^{\prime} \tag{4.2.3}
\end{gather*}
$$

respectively.
A direct computation gives

$$
\begin{aligned}
\frac{\partial L}{\partial x_{k}} & =\frac{1}{2} \sum_{i j} g_{i j, k} x_{i}^{\prime} x_{j}^{\prime}, \quad \text { where } g_{i j, k}=\frac{\partial g_{i j}}{\partial x_{k}} \\
& =\left(\frac{\partial L}{\partial x_{k}^{\prime}}\right)^{\prime}=\left(\sum_{i} g_{i k} x_{i}^{\prime}\right)^{\prime}=\sum_{i, m} g_{i k, m} x_{i}^{\prime} x_{m}^{\prime}+\sum_{i} g_{i k} x_{i}^{\prime \prime}
\end{aligned}
$$

Playing with indices to get

$$
\sum_{i, m} g_{i k, m} x_{i}^{\prime} x_{m}^{\prime}=\sum_{i j} g_{i k, j} x_{i}^{\prime} x_{j}^{\prime}=\sum_{j i} g_{j k, i} x_{j}^{\prime} x_{i}^{\prime}
$$

Note that the first equality is true because we just let $m=j$, the second equality is true because we interchange $i, j$ in the summand. Hence

$$
\sum_{i, m} g_{i k, m} x_{i}^{\prime} x_{m}^{\prime}=\frac{1}{2}\left(\sum_{i j} g_{i k, j} x_{i}^{\prime} x_{j}^{\prime}+\sum_{j i} g_{j k, i} x_{j}^{\prime} x_{i}^{\prime}\right)
$$

This implies that the Euler-Lagrange equation of the energy functional $E$ is

$$
\begin{equation*}
\sum_{i} g_{i k} x_{i}^{\prime \prime}+\frac{1}{2} \sum_{i j}\left(g_{k i, j}+g_{j k, i}-g_{i j, k}\right) x_{i}^{\prime} x_{j}^{\prime}=0 \tag{4.2.4}
\end{equation*}
$$

Let $g^{i j}$ denote the $i j$-th entry of the inverse of the matrix of $\left(g_{i j}\right)$. Multiply ( ${ }^{[3} .2 .4$ ) by $g^{k m}$ and sum over $k$ to get

$$
\sum_{i, k} g^{k m} g_{i k} x_{i}^{\prime \prime}+\frac{1}{2} \sum_{i j k} g^{k m}\left(g_{k i, j}+g_{j k, i}-g_{i j, k}\right) x_{i}^{\prime} x_{j}^{\prime}=0
$$

But $\sum_{k} g^{i k} g_{k m}=\delta_{i m}$ and $\left(g_{i j}\right)$ is symmetric. So we get

$$
\begin{equation*}
x_{m}^{\prime \prime}+\frac{1}{2} \sum_{i j k} g^{k m}\left(g_{k i, j}+g_{j k, i}-g_{i j, k}\right) x_{i}^{\prime} x_{j}^{\prime}=0 \tag{4.2.5}
\end{equation*}
$$

So we have proved the first part of the following theorem:
aq Theorem 4.2.1. If $x(t)$ is a critical point of the energy functional $E$, then
(1) $x$ is a solution of

$$
\begin{equation*}
x_{m}^{\prime \prime}+\sum_{i j} \Gamma_{i j}^{m} x_{i}^{\prime} x_{j}^{\prime}=0 \tag{4.2.6}
\end{equation*}
$$

where $\Gamma_{i j}^{m}=\frac{1}{2} \sum_{k} g^{k m}\left(g_{k i, j}+g_{j k, i}-g_{i j, k}\right)$.
(2) $\left\|x^{\prime}(t)\right\|$ is independent of $t$.

Proof. It remains to prove (2). We compute as follows:

$$
\begin{aligned}
& \left(\sum_{i j} g_{i j}(x) x_{i}^{\prime} x_{j}^{\prime}\right)^{\prime}=\sum_{i j k} g_{i j, k} x_{i}^{\prime} x_{j}^{\prime} x_{k}^{\prime}+\sum_{i j} g_{i j} x_{i}^{\prime \prime} x_{j}^{\prime}+g_{i j} x_{i}^{\prime} x_{j}^{\prime \prime} \\
& =\sum_{i j k} g_{i j, k} x_{i}^{\prime} x_{j}^{\prime} x_{k}^{\prime}-\left(\sum_{i j m} g_{i j} \Gamma_{m k}^{i} x_{m}^{\prime} x_{k}^{\prime} x_{j}^{\prime}+g_{i j} x_{i}^{\prime} \Gamma_{k m}^{j} x_{k}^{\prime} x_{m}^{\prime}\right) \\
& =\sum_{i j k} g_{i j, k} x_{i}^{\prime} x_{j}^{\prime} x_{k}^{\prime}-\sum_{i j m} g_{m j} \Gamma_{i k}^{m} x_{i}^{\prime} x_{k}^{\prime} x_{j}^{\prime}-\sum_{i j m} g_{i m} x_{i}^{\prime} \Gamma_{k j}^{m} x_{k}^{\prime} x_{j}^{\prime} \\
& =\sum_{i j k} g_{i j, k} x_{i}^{\prime} x_{j}^{\prime} x_{k}^{\prime}-\frac{1}{2} \sum_{i j k m r} g_{m j} g^{m r}[i k, r] x_{i}^{\prime} x_{k}^{\prime} x_{j}^{\prime}-\frac{1}{2} g_{i m} g^{m r}[k j, r] x_{i}^{\prime} x_{k}^{\prime} x_{j}^{\prime} \\
& =\sum_{i j k}\left(g_{i j, k}-\frac{1}{2} \sum_{r}\left(\delta_{j r}[i k, r]+\delta_{i r}[k j, r]\right)\right) x_{i}^{\prime} x_{k}^{\prime} x_{j}^{\prime} \\
& =\sum_{i j k}\left(g_{i j, k}-\frac{1}{2}[i k, j]-\frac{1}{2}[k j, i]\right) x_{i}^{\prime} x_{j}^{\prime} x_{k}^{\prime} \\
& =\sum_{i j k}\left(g_{i j, k}-\frac{1}{2}\left(g_{k j, i}+g_{i j, k}-g_{i k, j}\right)-\frac{1}{2}\left(g_{k i, j}+g_{i j, k}-g_{k j, i}\right)\right) x_{i}^{\prime} x_{j}^{\prime} x_{k}^{\prime} \\
& =0
\end{aligned}
$$

where $[i j, k]=g_{i k, j}+g_{j k, i}-g_{i j, k}$. (In the above computation, the second term of the third line is obtained by interchanging $i$ and $m$, and the third term is obtained by interchanging $j$ and $m$ ). This proves $\left\|x^{\prime}(t)\right\|^{2}=\sum_{i j} g_{i j}(x(t)) x_{i}^{\prime}(t) x_{j}^{\prime}(t)$ is constant.
Proposition 4.2.2. If $k:[a, b] \rightarrow[0,1]$ is a diffeomorphism, then
(1) $\mathcal{L}(x \circ k)=\mathcal{L}(x)$, i.e, the arc length of a curve on $M$ does not depend on the parametrization of the curve.
(2) if $x$ is a critical point of $J$, then so is $x \circ k$.

Proof. Suppose $s=k(r)$ and $y(r)=x(k(r))$. Then $\frac{d y}{d r}=\frac{d x}{d s} \frac{d s}{d r}$. Both statements follow from the change of variable formula in integration.

Proposition 4.2.3. If $x$ is a critical point of the arc length functional $\mathcal{L}$ and is parametrized proportional to its arc length, then $x$ is a critical point of the energy functional $E$. Conversely, if $x$ is a critical point of $E$, then $x$ is parametrized proportional to its arc length and is a critical point of $\mathcal{L}$.

Proof. If $x$ is a critical point of $\mathcal{L}$, then we may assume that $\alpha(t)=$ $f\left(x_{1}(t), x_{2}(t)\right)$ is parametrized by arc length parameter. This means we may assume $\left\|\alpha^{\prime}(t)\right\| \equiv 1$. But

$$
\left\|\alpha^{\prime}(t)\right\|^{2}=\sum_{i j} g_{i j} x_{i}^{\prime} x_{j}^{\prime}=L\left(x, x^{\prime}\right) \equiv 1
$$

Since $x$ is a critical point of $\mathcal{L}, x$ is a solution of $\left(\frac{\mathrm{bk}}{4.2} .3\right)$. But $L\left(x, x^{\prime}\right) \equiv 1$ implies that $x$ is a solution of (4.2.2), i.e., $x$ is a critical point for $E$.

Conversely, if $x$ is a critical point of $E$, then by Theorem $4.2 .1_{1}$ (ii), $L\left(x, x^{\prime}\right)$ is constant. If $L\left(x(t), x^{\prime}(t)\right)$ is constant, then $\left(\frac{1}{4.2 .3)}\right.$ and(4.2.2) are the same. This proves that if $x$ is a critical point of $E$, then $x$ is a critical point of $\mathcal{L}$.

A critical point of the arch length functional $\mathcal{L}$ on $\mathcal{M}_{p, q}$ is called a geodesic. Equation (4.2.6) is called the geodesic equation.

It follows from Proposition 4.0.14 that if $\gamma(t)=f\left(x_{1}(t), x_{2}(t)\right)$ is a curve on $M$ joining $p, q$, which is parametrized by arc length, and has shortest arc length, then $\gamma$ is a geodesic and $\left(x_{1}(t), x_{2}(t)\right)$ is a solution of the geodesic equation.

### 4.3. Calculus of variations of two variables.

Suppose $\Omega$ is an open subset of $\mathbb{R}^{2}$ such that the boundary $\partial \Omega$ is a simple closed curve. Fix a function $\gamma: \partial \Omega \rightarrow \mathbb{R}$, let $\mathcal{C}_{\gamma}(\Omega, \mathbb{R})$ denote the set of all smooth maps $u: \Omega \rightarrow \mathbb{R}$ so that $u \mid \partial \Omega=\gamma$. Let

$$
L: \Omega \times\left(\mathbb{R}^{3}\right)^{3} \rightarrow \mathbb{R}
$$

be a smooth function, and $J: \mathcal{C}_{\gamma}(\Omega, \mathbb{R}) \rightarrow \mathbb{R}$ the functional defined by

$$
\begin{equation*}
J(u)=\int_{\Omega} L\left(x, y, u(x, y), u_{x}(x, y), u_{y}(x, y)\right) d x d y \tag{4.3.1}
\end{equation*}
$$

ak

We want to derive the condition for $u$ being a critical point of $J$. This can be carried out in a similar manner as for one variable case except that we need to use the two dimensional version of The Fundamental Theorem of Calculus, i.e., the Green's Formula or the Stoke's theorem for dimension 2.

First, we recall the definition of the line integral,

$$
\oint_{\partial \Omega} P(x, y) d x+Q(x, y) d y .
$$

Let $(x, y):[a, b] \rightarrow \mathbb{R}^{2}$ be a parametrization of the boundary curve $\partial \Omega$. Then

$$
\oint_{\partial \Omega} P d x+Q d y=\int_{a}^{b} P(x(t), y(t)) x^{\prime}(t)+Q(x(t), y(t)) y^{\prime}(t) d t
$$

The Green's formula is
Theorem 4.3.1. Let $P, Q: \Omega \rightarrow \mathbb{R}$ be smooth functions. Then

$$
\begin{equation*}
\oint_{\partial \Omega} P(x, y) d x+Q(x, y) d y=\iint_{\Omega}\left(-P_{y}+Q_{x}\right) d x d y \tag{4.3.2}
\end{equation*}
$$

Suppose $u$ is a critcal point of $J$. We want to find the condition on $u$. Let $\beta$ be a variation of $u$ in $\mathcal{C}_{\gamma}(\Omega, \mathbb{R})$, and $h=\beta^{\prime}(0)$, i.e., $h(x, y)=\left.\frac{\partial}{\partial s}\right|_{s=0} \beta(s)(x, y)$. Since $\beta(s)|\mid \partial \Omega=\gamma$, we have $h| \partial \Omega=0$. Compute directly to get

$$
\begin{aligned}
& \left.\frac{d}{d s}\right|_{s=0} J(\beta(s))=\int_{\Omega} \frac{\partial L}{\partial u} h+\frac{\partial L}{\partial u_{x}} h_{x}+\frac{\partial L}{\partial u_{y}} h_{y} d x d y \\
& =\int_{\Omega} \frac{\partial L}{\partial u} h+\left(\frac{\partial L}{\partial u_{x}} h\right)_{x}-\left(\frac{\partial L}{\partial u_{x}}\right)_{x} h+\left(\frac{\partial L}{\partial u_{y}} h\right)_{y}-\left(\frac{\partial L}{\partial u_{y}}\right)_{y} h d x d y \\
& =\int_{\Omega}\left(\frac{\partial L}{\partial u}-\left(\frac{\partial L}{\partial u_{x}}\right)_{x}-\left(\frac{\partial L}{\partial u_{y}}\right)_{y}\right) h d x d y \\
& +\quad \int_{\Omega}\left(\frac{\partial L}{\partial u_{x}} h\right)_{x}+\left(\frac{\partial L}{\partial u_{y}} h\right)_{y} d x d y
\end{aligned}
$$

It follows from the Green's formula ( $(4.3 .2)$ that the second term is

$$
\int_{\partial \Omega} \frac{\partial L}{\partial u_{x}} h d y-\frac{\partial L}{\partial u_{y}} h d x .
$$

But $h \mid \partial \Omega=0$. So the line integral is zero. This shows that if $u$ is a critical point of $J$, then

$$
\int_{\Omega}\left(\frac{\partial L}{\partial u}-\left(\frac{\partial L}{\partial u_{x}}\right)_{x}-\left(\frac{\partial L}{\partial u_{y}}\right)_{y}\right) h d x d y=0
$$

for all smooth functions $h: \Omega \rightarrow \mathbb{R}$ that vanishes on $\partial \Omega$.
A similar proof gives the following analogue of Lemma ${ }_{4}^{\text {ag }} 0.15$ for $\mathbb{R}^{n}$.
Lemma 4.3.2. Let $\Omega \subset \mathbb{R}^{n}$ be an open subset, and $f: \bar{\Omega} \rightarrow \mathbb{R} a$ smooth map. If

$$
\int_{\Omega} f(x) g(x) d x_{1} \cdots d x_{n}=0
$$

for all smooth $g: \bar{\Omega} \rightarrow \mathbb{R}$ that vanishes on $\partial \Omega$, then $f=0$.
Hence we have proved
bq Theorem 4.3.3. If $u$ is a critical point of $J$ defined by ( ${ }_{4.3 \mathrm{k}}^{4.3 .1 \text { ), then }}$

$$
\begin{equation*}
\frac{\partial L}{\partial u}-\left(\frac{\partial L}{\partial u_{x}}\right)_{x}-\left(\frac{\partial L}{\partial u_{y}}\right)_{y}=0 \tag{4.3.3}
\end{equation*}
$$

(This is the Euler-Lagrange equation of $J$ defined by (ak
Example 4.3.4. Let $\Omega \subset \mathbb{R}^{2}, \alpha: \partial \Omega \rightarrow \mathbb{R}$ a smooth function, and $L(u)=u_{x}^{2}+u_{y}^{2}$ for $u \in C_{\gamma}^{\infty}(\Omega, \mathbb{R})$. Use Theorem $\frac{10}{4.3 .3 \text { to compute }}$ directly to see that $u$ is a critical point of the functional $J$ defined by $L$ if and only if

$$
u_{x x}+u_{y y}=0 .
$$

### 4.4. Minimal surfaces.

Let $M$ be an embedded surface in $\mathbb{R}^{3}$, and $f: \mathcal{O} \rightarrow U$ a local coordinate system on $M$, and $\bar{\Omega} \subset \mathcal{O}$ a compact domain. By Proposition ${ }^{\text {bm }} 3.3 .4$, the area of the parallelgram spanned by $f_{x_{1}}, f_{x_{2}}$ is $\sqrt{\operatorname{det}\left(g_{i j}\right)}$. To approximate the area of $f(\Omega)$, we first divide the domain $\Omega$ into small squares, and approximate the area of the image of $f$ of each square by the area of the parallelgram spanned by $f_{x_{1}} \triangle x_{1}$ and $f_{x_{2}} \triangle x_{2}$, then we sum up these approximations

$$
\sum \sqrt{\operatorname{det}\left(g_{i j}\right)} \triangle x_{1} \triangle x_{2}
$$

So the area of $f(\Omega)$ is

$$
\text { area }=\iint_{\Omega} \sqrt{\operatorname{det}\left(g_{i j}(x)\right)} d x_{1} d x_{2}=\iint_{\Omega} \sqrt{g_{11} g_{22}-g_{12}^{2}} d x_{1} d x_{2} .
$$

We want to show that if $f(\Omega)$ has minimum area among all surfaces in $\mathbb{R}^{3}$ with fixed boundary $f(\partial \Omega)$, then the mean curvature of $M$ must be zero. Since the condition for critical point of a variational functional is computed locally and locally a surface in $\mathbb{R}^{3}$ is a graph, we may assume that $f: \Omega \rightarrow U$ is a graph and the variations are done through graphs.

A surface in $\mathbb{R}^{3}$ is called minimal if its mean curvature $H=0$. We will show that if a surface $M$ in $\mathbb{R}^{3}$ is a critical point of the area functional then the mean cruvature of $M$ must be zero, i.e., $M$ is minimal. In fact, we will prove that the Euler-Lagrangian equation for the area functional is the equation $H=0$.

Let $u: \mathcal{O} \rightarrow \mathbb{R}$, and $f\left(x_{1}, x_{2}\right)=\left(x_{1}, x_{2}, u\left(x_{1}, x_{2}\right)\right)$ be the graph of $u$. Let $\bar{\Omega} \subset \mathcal{O}$. If $M=f(\Omega)$ has minimum area among all surfaces in $\mathbb{R}^{3}$ that have the same boundary as $f(\Omega)$, then $u$ must be a critical point of the area functional

$$
\mathcal{A}(u)=\int_{b 0} \int_{\Omega} \sqrt{\operatorname{det}\left(g_{i j}\right)} d x_{1} d x_{2}
$$

It is computed in Example 3.3.2 that

$$
g_{i i}=1+u_{x_{i}}^{2}, \quad g_{12}=u_{x_{1}} u_{x_{2}} .
$$

So $\operatorname{det}\left(g_{i j}\right)=\left(1+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right)$. The area functional is

$$
\mathcal{A}(u)=\iint_{\Omega}\left(1+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right)^{\frac{1}{2}} d x_{1} d x_{2}
$$

By Theorem ${ }_{4}^{\mathrm{bq}} 4.3 .3, u$ must satisfy the Euler-Lagrangian equation (4.3.3.3). A direct computation gives

$$
\left(\left(1+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right)^{-\frac{1}{2}} u_{x_{1}}\right)_{x_{1}}+\left(\left(1+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right)^{-\frac{1}{2}} u_{x_{2}}\right)_{x_{2}}=0
$$

So we have

$$
\frac{\left(1+u_{x_{2}}^{2}\right) u_{x_{1} x_{1}}-2 u_{x_{1}} u_{x_{2}} u_{x_{1} x_{2}}+\left(1+u_{x_{1}}^{2}\right) u_{x_{2} x_{2}}}{\left(1+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right)^{\frac{3}{2}}}=0 .
$$

But the left hand side is the mean curvature $H$ of the graph of $u$ (cf. Example $\stackrel{e}{3}_{3 .}^{3.6}$ formula ( 3.4 .2 ). So the mean curvature must be zero, i.e., $M$ is minimal. We have proved

Theorem 4.4.1. The Euler-Lagrangian equation of the area functional $\mathcal{A}$ is $H=0$.

## 5. Fundamental Theorem of surfaces in $\mathbb{R}^{3}$

Let $M$ be an embedded surface in $\mathbb{R}^{3}$, and $f: \mathcal{O} \rightarrow U \subset \mathbb{R}^{3}$ a local coordinate system on $M$. Our experience in curve theory tells us that we should find a moving frame on the surface and then differentiate the moving frame to get relations among the invariants. However, unlike the curves, there do not have natural local orthonormal frames on
general surfaces in $\mathbb{R}^{3}$. We will use several different moving frames $F=$ $\left(v_{1}, v_{2}, v_{3}\right)$ on the surface to derive the relations among local invariants. Express the $x$ and $y$ derivatives of the local frame $v_{i}$ in terms of $v_{1}, v_{2}, v_{3}$, then their coefficients can be written in terms of the two fundamental forms. Since $\left(v_{i}\right)_{x y}=\left(v_{i}\right)_{y x}$, we obtain a PDE relation for I and II. This is the Gauss-Codazzi equation of the surface. Conversely, given two symmetric bilinear forms $g, b$ on an open subset $\mathcal{O}$ of $\mathbb{R}^{2}$ such that $g$ is positive definite and $g, b$ satisfies the Gauss-Codazzi equation, then by the Frobenius Theorem there exists a unique surface in $\mathbb{R}^{3}$ having $g, b$ as the first and second fundamental forms respectively.

### 5.1. Gauss-Codazzi equation in local coordinates.

Let $f: \mathcal{O} \rightarrow U \subset M$ be a local coordinate system of an embedded surface $M$ in $\mathbb{R}^{3}$. We use the frame $\left(f_{x_{1}}, f_{x_{2}}, N\right)$, where

$$
N=\frac{f_{x_{1}} \times f_{x_{2}}}{\left\|f_{x_{1}} \times f_{x_{2}}\right\|}
$$

is the unit normal vector field. Since $f_{x_{1}}, f_{x_{2}}, N$ form a basis of $\mathbb{R}^{3}$, the partial derivatives of $f_{x_{i}}$ and $N$ can be written as linear combinations of $f_{x_{1}}, f_{x_{2}}$ and $N$. So we have

$$
\left\{\begin{array}{l}
\left(f_{x_{1}}, f_{x_{2}}, N\right)_{x_{1}}=\left(f_{x_{1}}, f_{x_{2}}, N\right) P  \tag{5.1.1}\\
\left(f_{x_{1}}, f_{x_{2}}, N\right)_{x_{2}}=\left(f_{x_{1}}, f_{x_{2}}, N\right) Q
\end{array}\right.
$$

where $P=\left(p_{i j}\right), Q=\left(q_{i j}\right)$ are $g l(3)$-valued maps. This means that

$$
\left\{\begin{array} { l } 
{ f _ { x _ { 1 } x _ { 1 } } = p _ { 1 1 } f _ { x _ { 1 } } + p _ { 2 1 } f _ { x _ { 2 } } + p _ { 3 1 } N , } \\
{ f _ { x _ { 2 } x _ { 1 } } = p _ { 1 2 } f _ { x _ { 1 } } + p _ { 2 2 } f _ { x _ { 2 } } + p _ { 3 2 } N , } \\
{ N _ { x _ { 1 } } = p _ { 1 3 } f _ { x _ { 1 } } + p _ { 2 3 } f _ { x _ { 3 } } + p _ { 3 3 } N , }
\end{array} \left\{\begin{array}{l}
f_{x_{1} x_{2}}=q_{11} f_{x_{1}}+q_{21} f_{x_{2}}+q_{31} N \\
f_{x_{2} x_{2}}=q_{12} f_{x_{1}}+q_{22} f_{x_{2}}+q_{32} N \\
N_{x_{2}}=q_{13} f_{x_{1}}+q_{23} f_{x_{3}}+q_{33} N
\end{array}\right.\right.
$$

Recall that the fundamental forms are given by

$$
g_{i j}=f_{x_{i}} \cdot f_{x_{j}}, \quad \ell_{i j}=-f_{x_{i}} \cdot N_{x_{j}}=f_{x_{i} x_{j}} \cdot N .
$$

We want to express $P$ and $Q$ in terms of $g_{i j}$ and $h_{i j}$. To do this, we need the following Propositions.
bs Proposition 5.1.1. Let $V$ be a vector space with an inner product ( , ), $v_{1}, \cdots, v_{n}$ a basis of $V$, and $g_{i j}=\left(v_{i}, v_{j}\right)$. Let $\xi \in V, \xi_{i}=\left(\xi, v_{i}\right)$, and $\xi=\sum_{i=1}^{n} x_{i} v_{i}$. Then

$$
\left(\begin{array}{c}
x_{1} \\
\cdot \\
\cdot \\
x_{n}
\end{array}\right)=G^{-1}\left(\begin{array}{c}
\xi_{1} \\
\cdot \\
\cdot \\
\xi_{n}
\end{array}\right)
$$

where $G=\left(g_{i j}\right)$.
Proof. Note that

$$
\xi_{i}=\left(\xi, v_{i}\right)=\left(\sum_{j=1}^{n} x_{j} v_{j}, v_{i}\right)=\sum_{j=1}^{n} x_{j}\left(v_{j}, v_{i}\right)=\sum_{j=1}^{n} g_{j i} x_{j} .
$$

So $\left(\xi_{1}, \cdots, \xi_{n}\right)^{t}=G\left(x_{1}, \cdots, x_{n}\right)^{t}$.
Proposition 5.1.2. The following statements are true:
(1) The glf(3) valued functions $P=\left(p_{i j}\right)$ and $Q=\left(q_{i j}\right)$ in equation (5.1.1) can be written in terms of $g_{i j}, \ell_{i j}$, and first partial derivatives of $g_{i j}$.
(2) The entries $\left\{p_{i j}, q_{i j} \mid 1 \leq i, j \leq 2\right\}$ can be computed from the first fundamental form.

Proof. We claim that

$$
f_{x_{i} x_{j}} \cdot f_{x_{k}}, \quad f_{x_{i} x_{j}} \cdot N, \quad N_{x_{i}} \cdot f_{x_{j}}, \quad N_{x_{i}} \cdot N
$$

can be expressed in terms of $g_{i j}, \ell_{i j}$ and first partial derivatives of $g_{i j}$. Then the Proposition follows from Proposition 5.1.1. To prove the claim, we proceed as follows:

$$
\begin{cases}f_{x_{i} x_{i}} \cdot f_{x_{i}}=\frac{1}{2}\left(g_{i i}\right)_{x_{i}}, & \text { if } i \neq j,  \tag{5.1.2}\\ f_{x_{i} x_{j}} \cdot f_{x_{i}}=\frac{1}{2}\left(g_{i i}\right)_{x_{j}}, & \text { if } i \neq j\end{cases}
$$

$$
\left\{\begin{array}{l}
f_{x_{i} x_{j}} \cdot N=\ell_{i j},  \tag{5.1.3}\\
N_{x_{i}} \cdot f_{x_{j}}=-\ell_{i j}, \\
N_{x_{i}} \cdot N=0
\end{array}\right.
$$

Let

$$
G=\left(\begin{array}{ccc}
g_{11} & g_{12} & 0  \tag{5.1.4}\\
g_{12} & g_{22} & 0 \\
0 & 0 & 1
\end{array}\right)
$$

By Proposition ${ }^{\text {bs }}$ 5.1.1, we have

$$
\left\{\begin{align*}
P & =\left(\begin{array}{ccc}
g^{11} & g^{12} & 0 \\
g^{12} & g^{22} & 0 \\
0 & 0 & 1
\end{array}\right)\left(\begin{array}{ccc}
\frac{1}{2}\left(g_{11}\right)_{x_{1}} & \frac{1}{2}\left(g_{11}\right)_{x_{2}} & -\ell_{11} \\
\left(g_{12}\right)_{x_{1}}-\frac{1}{2}\left(g_{11}\right)_{x_{2}} & \frac{1}{2}\left(g_{22}\right)_{x_{1}} & -\ell_{12} \\
\ell_{11} & \ell_{12} & 0
\end{array}\right)  \tag{5.1.5}\\
& =G^{-1} A_{1}, \\
Q & =\left(\begin{array}{ccc}
g^{11} & g^{12} & 0 \\
g^{12} & g^{22} & 0 \\
0 & 0 & 1
\end{array}\right)\left(\begin{array}{ccc}
\frac{1}{2}\left(g_{11}\right)_{x_{2}} & \left(g_{12}\right)_{x_{2}}-\frac{1}{2}\left(g_{22}\right)_{x_{1}} & -\ell_{12} \\
\frac{1}{2}\left(g_{22}\right)_{x_{1}} & \frac{1}{2}\left(g_{22}\right)_{x_{2}} & -\ell_{22} \\
\ell_{12} & \ell_{22} & 0
\end{array}\right) \\
& =G^{-1} A_{2} .
\end{align*}\right.
$$

This proves the Proposition.
Formula (5. ${ }^{\mathrm{bt}} 1.5$ ) gives explicit formulas for entries of $P$ and $Q$ in terms of $g_{i j}$ and $\ell_{i j}$. Moreover, they are related to the Christofell symbols $\Gamma_{j k}^{i}$ arise in the geodesic equation (4.2.6) in Theorem ${ }_{4}^{\text {aq. }}$.2.1. Recall that

$$
\Gamma_{i j}^{k}=\frac{1}{2} g^{k m}[i j, m],
$$

where $\left(g^{i j}\right)$ is the inverse matrix of $\left(g_{i j}\right),[i j, k]=g_{k i, j}+g_{j k, i}-g_{i j, k}$, and $g_{i j, k}=\frac{\partial g_{i j}}{\partial x_{k}}$.

Theorem 5.1.3. For $1 \leq i, j \leq 2$, we have

$$
\begin{equation*}
p_{j i}=\Gamma_{i 1}^{j}, \quad q_{j i}=\Gamma_{i 2}^{j} \tag{5.1.6}
\end{equation*}
$$

ey
Proof. Note that (b.1.5) implies

$$
\begin{cases}p_{11} & =\frac{1}{2} g^{11} g_{11,1}+g^{12}\left(g_{12,1}-\frac{1}{2} g_{11,2}\right)=\Gamma_{11}^{1}  \tag{5.1.7}\\ p_{12} & =\frac{1}{2} g^{11} g_{11,2}+\frac{1}{2} g^{12} g_{22,1}=\Gamma_{21}^{1} \\ p_{21} & =\frac{1}{2} g^{12} g_{11,1}+g^{22}\left(g_{12,1}-\frac{1}{2} g_{11,2}\right)=\Gamma_{11}^{2} \\ p_{22} & =\frac{1}{2} g^{12} g_{11,2}+\frac{1}{2} g^{22} g_{22,1}=\Gamma_{21}^{2}, \\ q_{11} & =\frac{1}{2} g^{11} g_{11,2}+\frac{1}{2} g^{12} g_{22,1}=\Gamma_{12}^{1} \\ q_{12} & =g^{11}\left(g_{12,2}-\frac{1}{2} g_{22,1}\right)+\frac{1}{2} g^{2} g_{22,2}=\Gamma_{22}^{1} \\ q_{21} & =\frac{1}{2} g^{12} g_{11,2}+\frac{1}{2} g^{22} g_{22,1}=\Gamma_{12}^{2} \\ q_{22} & =g^{12}\left(g_{12,2}-\frac{1}{2} g_{22,1}\right)+\frac{1}{2} g^{22} g_{22,2}=\Gamma_{22}^{2}\end{cases}
$$

Note that

$$
q_{11}=p_{12}, \quad q_{21}=p_{22}
$$

bu Theorem 5.1.4. The Fundamental Theorem of surfaces in $\mathbb{R}^{3}$. Suppose $M$ is an embedded surface in $\mathbb{R}^{3}$, and $f: \mathcal{O} \rightarrow U$ a local coordinate system on $M$, and $g_{i j}, \ell_{i j}$ are the coefficients of I, II. Let $P_{\mathrm{b}^{\circ} \mathrm{t}} Q$ be the smooth gl(3)-valued maps defined in terms of $g_{i j}$ and $\ell_{i j}$ by (5.1.5). Then $P, Q$ satisfy

$$
\begin{equation*}
P_{x_{2}}-Q_{x_{1}}=[P, Q] \tag{5.1.8}
\end{equation*}
$$

Conversely, let $\mathcal{O}$ be an open subset of $\mathbb{R}^{2},\left(g_{i j}\right),\left(\ell_{i j}\right): \mathcal{O} \rightarrow g l(2)$ smooth maps such that $\left(g_{i j}\right)$ is positive definite and $\left(\ell_{i j}\right)$ is symmetric, and $P, Q: U \rightarrow g l(3)$ the maps defined by (5.1.5). Suppose $P, Q$ satisfies the compatibility equation (5.1.8). Let $\left(x_{1}^{0}, x_{2}^{0}\right) \in \mathcal{O}, p_{0} \in \mathbb{R}^{3}$, and $u_{1}, u_{2}, u_{3}$ a basis of $\mathbb{R}^{3}$ so that $u_{i} \cdot u_{j}=g_{i j}\left(x_{1}^{0}, x_{2}^{0}\right)$ and $u_{i} \cdot u_{3}=0$ for $1 \leq i, j \leq 2$. Then there exists an open subset $\mathcal{O}_{0} \subset \mathcal{O}$ of $\left(x_{1}^{0}, x_{2}^{0}\right)$ and a unique immersion $f: \mathcal{O}_{0} \rightarrow \mathbb{R}^{3}$ so that $f$ maps $\mathcal{O}_{0}$ homeomorphically to $f\left(\mathcal{O}_{0}\right)$ such that
(1) the first and second fundamental forms of the embedded surface $f\left(\mathcal{O}_{0}\right)$ are given by $\left(g_{i j}\right)$ and $\left(\ell_{i j}\right)$ respectively,
(2) $f\left(x_{1}^{0}, x_{2}^{0}\right)=p_{0}$, and $f_{x_{i}}\left(x_{1}^{0}, x_{2}^{0}\right)=u_{i}$ for $i=1,2$.

Proof. We have proved the first half of the theorem, and it remains to prove the second half. We assume that $P, Q_{\mathrm{bf}}$ satisfy the compatibility condition (5.1.8). So Frobenius Theorem 2.4.1 implies that the following system has a unique local solution

$$
\left\{\begin{array}{l}
\left(v_{1}, v_{2}, v_{3}\right)_{x_{1}}=\left(v_{1}, v_{2}, v_{3}\right) P  \tag{5.1.9}\\
\left(v_{1}, v_{2}, v_{3}\right)_{x_{2}}=\left(v_{1}, v_{2}, v_{3}\right) Q \\
\left(v_{1}, v_{2}, v_{3}\right)\left(x_{1}^{0}, x_{2}^{0}\right)=\left(u_{1}, u_{2}, u_{3}\right)
\end{array}\right.
$$

Next we want to solve

$$
\left\{\begin{array}{l}
f_{x_{1}}=v_{1} \\
f_{x_{2}}=v_{2} \\
f\left(x_{1}^{0}, x_{2}^{0}\right)=p_{0}
\end{array}\right.
$$

The compatibility condition is $\left(v_{1}\right)_{x_{2}}=\left(v_{2}\right)_{x_{1}}$. But

$$
\left(v_{1}\right)_{x_{2}}=\sum_{j=1}^{3} q_{j 1} v_{j}, \quad\left(v_{2}\right)_{x_{1}}=\sum_{j=1}^{3} p_{j 2} v_{2}
$$

It follows from ( $5 . \frac{b t}{5 .} 1.5$ ) that the second column of $P$ is equal to the first column of $Q$. So $\left(v_{1}\right)_{x_{2}}=\left(v_{2}\right)_{x_{1}}$, and hence there exists a unique $f$.

We will prove below that $f$ is an immersion, $v_{3}$ is perpendicular to the surface $f,\left\|v_{3}\right\|=1, f_{x_{i}} \cdot f_{x_{j}}=g_{i j}$, and $\left(v_{3}\right)_{x_{i}} \cdot f_{x_{j}}=-\ell_{i j}$, i.e., $f$ is
a surface in $\mathbb{R}^{3}$ with

$$
\mathrm{I}=\sum_{i j} g_{i j} d x_{i} d x_{j}, \quad \mathrm{II}=\sum_{i j} \ell_{i j} d x_{i} d x_{j}
$$

as its first and second fundamental forms. The first step is to prove that the $3 \times_{b} 3$ matrix function $\Phi=\left(v_{i} \cdot v_{j}\right)$ is equal to the matrix $G$ defined by (5.1.5). To do this, we compute the first derivative of $\Phi$. Since $v_{1}, v_{2}, v_{3}$ satisfy (5.1.9), a direct computation gives

$$
\begin{aligned}
\left(v_{i} \cdot v_{j}\right)_{x_{1}} & =\left(v_{i}\right)_{x_{1}} \cdot v_{j}+v_{i} \cdot\left(v_{j}\right)_{x_{1}} \\
& =\sum_{k} p_{k i} v_{k} \cdot v_{j}+p_{k j} v_{k} \cdot v_{i}=\sum_{k} p_{k i} g_{j k}+g_{i k} p_{k j} \\
& =(G P)_{j i}+(G P)_{i j}=\left(G P+(G P)^{t}\right)_{i j}
\end{aligned}
$$

Formula (5.1.5) implies that $G P=G\left(G^{-1} A_{1}\right)=A_{1}$ and $A_{1}+A_{1}^{t}=G_{x_{1}}$. Hence $(G P)^{t}+G P=G_{x_{1}}$ and

$$
\Phi_{x_{1}}=G_{x_{1}}
$$

A similar computation implies that

$$
\Phi_{x_{2}}=G_{x_{2}}
$$

But the initial value $\Phi\left(x_{1}^{0}, x_{2}^{0}\right)=G\left(x_{1}^{0}, x_{2}^{0}\right)$. So $\Phi=G$. In other words, we have shown that

$$
f_{x_{i}} \cdot f_{x_{j}}=g_{i j}, \quad f_{x_{i}} \cdot v_{3}=0
$$

Thus
(1) $f_{x_{1}}, f_{x_{2}}$ are linearly independent, i.e., $f$ is an immersion,
(2) $v_{3}$ is the unit normal field to the surface $f$,
(3) the first fundamental form of $f$ is $\sum_{i j} g_{i j} d x_{i} d x_{j}$.

To compute the second fundamental form of $f$, we use (5.1.9) to compute

$$
\begin{aligned}
-\left(v_{3}\right)_{x_{1}} \cdot v_{j} & =\left(g^{11} \ell_{11}+g^{12} \ell_{12}\right) v_{1} \cdot v_{j}+\left(g^{12} \ell_{11}+g^{22} \ell_{12}\right) v_{2} \cdot v_{j} \\
& =\left(g^{11} \ell_{11}+g^{12} \ell_{12}\right) g_{1 j}+\left(g^{12} \ell_{11}+g^{22} \ell_{12}\right) g_{2 j} \\
& =\ell_{11}\left(g^{11} g_{1 j}+g^{12} g_{2 j}\right)+\ell_{12}\left(g^{21} g_{1 j}+g^{22} g_{2 j}\right) \\
& =\ell_{11} \delta_{1 j}+\ell_{12} \delta_{2 j} .
\end{aligned}
$$

So

$$
-\left(v_{3}\right)_{x_{1}} \cdot v_{1}=\ell_{11}, \quad-\left(v_{3}\right)_{x_{1}} \cdot v_{2}=\ell_{12} .
$$

Similar computations imply that

$$
-\left(v_{2}\right)_{x_{2}} \cdot v_{j}=\ell_{2 j} .
$$

This proves that $\sum_{i j} \ell_{i j} d x_{i} d x_{j}$ is the second fundamental form of $f$.

System ( ${ }^{\mathrm{bv}} 5.1 .8$ ) with $P, Q$ defined by ( ${ }^{\mathrm{bt}} 5.1 .5$ ) is called the GaussCodazzi equation for the surface $f(\mathcal{O})$, which is a second order PDE with 9 equations for six functions $g_{i j}$ and $\ell_{i j}$. Equaqtion (5.1.8) is too complicated to memorize. It is more useful and simpler to just remember how to derive the Gauss-Codazzi equation.

It follows from ( $5^{\mathrm{br}} 1.1$ ), ( ${ }^{\mathrm{bt}} .1 .5$ ), and ( ${ }^{\mathrm{ey}} 1.6$ ) that we have

$$
\begin{aligned}
& f_{x_{i} x_{1}}=\sum_{j=1}^{2} p_{j i} f_{x_{j}}+\ell_{i 1} N=\sum_{j=1}^{2} \Gamma_{i 1}^{j} f_{x_{j}}+\ell_{i 1} N \\
& f_{x_{i} x_{2}}=\sum_{j=1}^{2} q_{j 2=i} f_{x_{j}}+\ell_{i 2} N=\sum_{j=1}^{2} \Gamma_{i 2}^{j} f_{x_{j}}+\ell_{i 2} N
\end{aligned}
$$

where $p_{i j}$ and $q_{i j}$ are defined in (5.1.5).
So we have

$$
\begin{equation*}
f_{x_{i} x_{j}}=\Gamma_{i j}^{1} f_{x_{1}}+\Gamma_{i j}^{2} f_{x_{2}}+\ell_{i j} N, \tag{5.1.10}
\end{equation*}
$$

Proposition 5.1.5. Let $f: \mathcal{O} \rightarrow \mathbb{R}^{3}$ be a local coordinate system of an embedded surface $M$ i $\eta \mathbb{R}^{3}$, and $\alpha(t)=f\left(x_{1}(t), x_{2}(t)\right)$. Then $\alpha$ satisfies the geodesic equation (4.2.6) if and only if $\alpha^{\prime \prime}(t)$ is normal to $M$ at $\alpha(t)$ for all $t$.

Proof. Differentiate $\alpha^{\prime}$ to get $\alpha^{\prime}=\sum_{i=1}^{2} f_{x_{i}} x_{i}^{\prime}$. So

$$
\begin{aligned}
\alpha^{\prime \prime} & =\sum_{i, j=1}^{2} f_{x_{i} x_{j}} x_{i}^{\prime} x_{j}^{\prime}+f_{x_{i}} x_{i}^{\prime \prime} \\
& =\sum_{i, j, k=1}^{2} \Gamma_{i j}^{k} f_{x_{k}} x_{i}^{\prime} x_{j}^{\prime}+\ell_{i j} N+f_{x_{i}} x_{i}^{\prime \prime} \\
& =\sum_{i, j=1}^{2}\left(\Gamma_{i j}^{k} x_{i}^{\prime} x_{j}^{\prime}+x_{k}^{\prime \prime}\right) f_{x_{k}}+\ell_{i j} N=0+\ell_{i j} N=\ell_{i j} N .
\end{aligned}
$$

### 5.2. The Gauss Theorem.

Equation (5. ${ }^{\mathrm{bv}} .8$ ) is the Gauss-Codazzi equation for $M$.

The Gaussian curvature $K$ is defined to be the determinant of the shape operator $-d N$, which depends on both the first and second fundamental forms of the surface. In fact, by Proposition 3.4.2

$$
K=\frac{\ell_{11} \ell_{22}-\ell_{12}^{2}}{g_{11} g_{22}-g_{12}^{2}}
$$

We will show below that $K$ can be computed in terms of $g_{i j}$ alone. Equate the 12 entry of equation (5.1.8) to get

$$
\left(p_{12}\right)_{x_{2}}-\left(q_{12}\right)_{x_{1}}=\sum_{j=1}^{3} p_{1 j} q_{j 2}-q_{1 j} p_{j 2}
$$

Recall that formula ( $\stackrel{\mathrm{e} \mathrm{r}}{5 \mathrm{r}} 1.7$ ) gives $\left\{p_{i j}, q_{i j} \mid 1 \leq i, j \leq 2\right\}$ in terms of the first fundamental form I. We move terms involves $p_{i j}, q_{i j}$ with $1 \leq i, j \leq 2$ to one side to get

$$
\begin{equation*}
\left(p_{12}\right)_{x_{2}}-\left(q_{12}\right)_{x_{1}}-\sum_{j=1}^{2} p_{1 j} q_{j 2}-q_{1 j} p_{j 2}=p_{13} q_{32}-q_{13} p_{32} \tag{5.2.1}
\end{equation*}
$$

We claim that the right hand side of $\left(\frac{\mathrm{cu}}{5.2 .1}\right)$ is equal to

$$
-g^{11}\left(\ell_{11} \ell_{22}-\ell_{12}^{2}\right)=-g^{11}\left(g_{11} g_{22}-g_{12}^{2}\right) K
$$



$$
\begin{array}{ll}
p_{13}=-\left(g^{11} \ell_{11}+g^{12} \ell_{12}\right), & p_{32}=\ell_{12}, \\
q_{13}=-\left(g^{11} \ell_{12}+g^{12} \ell_{22}\right), & q_{32}=\ell_{22}
\end{array}
$$

So we get

$$
\begin{equation*}
\left(p_{12}\right)_{x_{2}}-\left(q_{12}\right)_{x_{1}}-\sum_{j=1}^{2} p_{1 j} q_{j 2}-q_{1 j} p_{j 3}=-g^{11}\left(g_{11} g_{22}-g_{12}^{2}\right) K \tag{5.2.2}
\end{equation*}
$$

Hence we have proved the claim and also obtained a formula of $K$ purely in terms of $g_{i j}$ and their derivatives:

$$
K=-\frac{\left(p_{12}\right)_{x_{2}}-\left(q_{12}\right)_{x_{1}}-\sum_{j=1}^{2} p_{1 j} q_{j 2}-q_{1 j} p_{j 3}}{g^{11}\left(g_{11} g_{22}-g_{12}^{2}\right)}
$$

This proves
ef Theorem 5.2.1. Gauss Theorem. The Gaussian curvature of a surface in $\mathbb{R}^{3}$ can be computed from the first fundamental form.

The equation ( $\stackrel{\mathrm{es}}{5.2 .2)}$ ), obtained by equating the 12-entry of ( ${ }_{5}^{\mathrm{bvv}} 1.8$ ), is the Gauss equation.

A geometric quantity on an embedded surface $M$ in $\mathbb{R}^{n}$ is called intrinsic if it only depends on the first fundamental form I. Otherwise, the property is called extrinsic, i.e., it depends on both I and II.

We have seen that the Gaussian curvature and geodesics are intrinsic quantities, and the mean curvature is extrinsic.

If $\phi: M_{1} \rightarrow M_{2}$ is a diffeomorphism and $f\left(x_{1}, x_{2}\right)$ is a local coordinates on $M_{1}$, then $\phi \circ f\left(x_{1}, x_{2}\right)$ is a local coordinate system of $M_{2}$. The diffeomorphism $\phi$ is an isometry if the first fundamental forms for $M_{1}, M_{2}$ are the same written in terms of $d x_{1}, d x_{2}$. In particular,
(i) $\phi$ preserves anlges and arc length, i.e., the arc length of the curve $\phi(\alpha)$ is the same as the curve $\alpha$ and the angle between the curves $\phi(\alpha)$ and $\phi(\beta)$ is the same as the angle between $\alpha$ and $\beta$,
(ii) $\phi$ maps geodesics to geodesics.

Euclidean plane geometry studies the geometry of triangles. Note that triangles can be viewed as a triangle in the plane with each side being a geodesic. So a natural definition of a triangle on an embedded surface $M$ is a piecewise smooth curve with three geodesic sides and any two sides meet at an angle lie in $(0, \pi)$. One important problem in geometric theory of $M$ is to understand the geometry of triangles on $M$. For example, what is the sum of interior angles of a triangle on an embedded surface $M$ ? This will be answered by the Gauss-Bonnet Theorem.

Note that the first fundamental forms for the plane

$$
f\left(x_{1}, x_{2}\right)=\left(x_{1}, x_{2}, 0\right)
$$

and the cylinder

$$
h\left(x_{1}, x_{2}\right)=\left(\cos x_{1}, \sin x_{1}, x_{2}\right)
$$

have the same and is equal to $\mathrm{I}=d x_{1}^{2}+d x_{2}^{2}$, and $\underset{\text { cs }}{\text { b }}$ both surfaces have constant zero Gaussian curvature (cf. Examples $\frac{10}{\text { Cs }} 4.3$ and $\frac{\mathrm{ct}}{3.4 .5 \text { ). We }}$ have also proved that geodesics are determined by I alone. So the geometry of triangles on the cylinder is the same as the geometry of triangles in the plane. For example, the sum of interior anlges of a triangle on the plane (and hence on the cylinder) must be $\pi$. In fact, let $\phi$ denote the map from $(0,2 \pi) \times \mathbb{R}$ to the cylinder minus the line ( $1,0, x_{2}$ ) defined by

$$
\phi\left(x_{1}, x_{2}, 0\right)=\left(\cos x_{1}, \sin x_{1}, x_{2}\right)
$$

Then $\phi$ is an isometry.

### 5.3. Gauss-Codazzi equation in orthogonal coordinates.

If the local coordinates $x_{1}, x_{v}$ are orthogonal, i.e., $g_{12}=0$, then the Gauss-Codazzi equation (5.1.8) becomes much simplier. Instead of putting $g_{12}=0$ to (5.1.8), we derive the Gauss-Codazzi equation directly using an o.n. moving frame. We write

$$
g_{11}=A_{1}^{2}, \quad g_{22}=A_{2}^{2}, \quad g_{12}=0
$$

Let

$$
e_{1}=\frac{f_{x_{1}}}{A_{1}}, \quad e_{2}=\frac{f_{x_{2}}}{A_{2}}, \quad e_{3}=N .
$$

Then $\left(e_{1}, e_{2}, e_{3}\right)$ is an o.n. moving frame on $M$. Write

$$
\left\{\begin{array}{l}
\left(e_{1}, e_{2}, e_{3}\right)_{x_{1}}=\left(e_{1}, e_{2}, e_{3}\right) \tilde{P} \\
\left(e_{1}, e_{2}, e_{3}\right)_{x_{2}}=\left(e_{1}, e_{2}, e_{3}\right) \tilde{Q}
\end{array}\right.
$$

Since $\left(e_{1}, e_{2}, e_{3}\right)$ is orthogonal, $\tilde{P}, \tilde{Q}$ are skew-symmetric. Moreover,

$$
\tilde{p}_{i j}=\left(e_{j}\right)_{x_{1}} \cdot e_{i}, \quad \tilde{q}_{i j}=\left(e_{j}\right)_{x_{2}} \cdot e_{i} .
$$

A direct computation gives

$$
\begin{aligned}
\left(e_{1}\right)_{x_{1}} \cdot e_{2} & =\left(\frac{f_{x_{1}}}{A_{1}}\right)_{x_{1}} \cdot \frac{f_{x_{2}}}{A_{2}}=\frac{f_{x_{1} x_{1}} \cdot f_{x_{2}}}{A_{1} A_{2}} \\
& =\frac{\left(f_{x_{1}} \cdot f_{x_{2}}\right)_{x_{1}}-f_{x_{1}} \cdot f_{x_{1} x_{2}}}{A_{1} A_{2}} \\
& =-\frac{\left(\frac{1}{2} A_{1}^{2}\right)_{x_{2}}}{A_{1} A_{2}}=-\frac{\left(A_{1}\right)_{x_{2}}}{A_{2}} .
\end{aligned}
$$

Similar computation gives the coefficients $\tilde{p}_{i j}$ and $\tilde{q}_{i j}$ :

$$
\tilde{P}=\left(\begin{array}{ccc}
0 & \frac{\left(A_{1}\right) x_{2}}{A_{2}} & -\frac{\ell_{11}}{A_{1}}  \tag{5.3.1}\\
-\frac{\left(A_{1}\right)_{x_{2}}}{A_{2}} & 0 & -\frac{\ell_{12}}{A_{2}} \\
\frac{\ell_{11}}{A_{1}} & \frac{\ell_{12}}{A_{2}} & 0
\end{array}\right), \quad \tilde{Q}=\left(\begin{array}{ccc}
0 & -\frac{\left(A_{2}\right) x_{1}}{A_{1}} & -\frac{\ell_{12}}{A_{1}} \\
\frac{\left(A_{2}\right)_{x_{1}}}{A_{1}} & 0 & -\frac{\ell_{22}}{A_{2}} \\
\frac{\ell_{12}}{A_{1}} & \frac{\ell_{22}}{A_{2}} & 0
\end{array}\right)
$$

To get the Gauss-Codazzi equation of the surface parametrized by an orthogonal coordinates we only need to compute the 21-th, 31-th, and 32 -the entry of the following equation

$$
(\tilde{P})_{x_{2}}-(\tilde{Q})_{x_{1}}=[\tilde{P}, \tilde{Q}]
$$

and we obtain

$$
\left\{\begin{array}{l}
-\left(\frac{\left(A_{1}\right)_{x_{2}}}{A_{2}}\right)_{x_{2}}-\left(\frac{\left(A_{2}\right)_{x_{1}}}{A_{1}}\right)_{x_{1}}=\frac{\ell_{11} \ell_{22}-\ell_{12}^{2}}{A_{1} A_{2}},  \tag{5.3.2}\\
\left(\frac{\ell_{11}}{A_{1}}\right)_{x_{2}}-\left(\frac{\ell_{12}}{A_{1}}\right)_{x_{1}}=\frac{\ell_{12}\left(A_{2}\right)_{x_{1}}}{A_{1} A_{2}}+\frac{\ell_{22}\left(A_{1}\right)_{2}}{A_{2}^{2}} \\
\left(\frac{\ell_{12}}{A_{2}}\right)_{x_{2}}-\left(\frac{\ell_{22}}{A_{2}}\right)_{x_{1}}=-\frac{\ell_{11}\left(A_{2}\right) x_{1}}{A_{1}^{2}}-\frac{\ell_{12}\left(A_{1}\right)_{x_{2}}}{A_{1} A_{2}} .
\end{array}\right.
$$

The first equation of (5.3.2) is called the Gauss equation. Note that the Gaussian curvature is

$$
K=\frac{\ell_{11} \ell_{22}-\ell_{12}^{2}}{\left(A_{1} A_{2}\right)^{2}}
$$

So we have

$$
\begin{equation*}
K=-\frac{\left(\frac{\left(A_{1}\right)_{x_{2}}}{A_{2}}\right)_{x_{2}}+\left(\frac{\left(A_{2}\right)_{x_{1}}}{A_{1}}\right)_{x_{1}}}{A_{1} A_{2}} \tag{5.3.3}
\end{equation*}
$$

We have seen that the Gauss-Codazzi equation becomes much simpler in orthogonal coordinates. Can we always find local orthogonal coordinates on a surface in $\mathbb{R}^{3}$ ? This question can be answered by the following theorem, which we state without a proof.
bx Theorem 5.3.1. Suppose $f: \mathcal{O} \rightarrow \mathbb{R}^{3}$ be a surface, $x_{0} \in \mathcal{O}$, and $Y_{1}, Y_{2}: \mathcal{O} \rightarrow \mathbb{R}^{3}$ smooth maps so that $Y_{1}\left(x_{0}\right), Y_{2}\left(x_{0}\right)$ are linearly independent and tangent to $M=f(\mathcal{O})$ at $f\left(x_{0}\right)$. Then there exist open subset $\mathcal{O}_{0}$ of $\mathcal{O}$ containing $x_{0}$, open subset $\mathcal{O}_{1}$ of $\mathbb{R}^{2}$, and a diffeomorphism $h: \mathcal{O}_{1} \rightarrow \mathcal{O}_{0}$ so that $(f \circ h)_{y_{1}}$ and $(f \circ h)_{y_{2}}$ are parallel to $Y_{1} \circ h$ and $Y_{2} \circ h$.

The above theorem says that if we have two linearly independent vector fields $Y_{1}, Y_{2}$ on a surface, then we can find a local coordinate system $\phi\left(y_{1}, y_{2}\right)$ so that $\phi_{y_{1}}, \phi_{y_{2}}$ are parallel to $Y_{1}, Y_{2}$ respectively.

Given an arbitrary local coordinate system $f\left(x_{1}, x_{2}\right)$ on $M$, we apply the Gram-Schmidt process to $f_{x_{1}}, f_{x_{2}}$ to construct smooth o.n. vector fields $e_{1}, e_{2}$ :

$$
\begin{aligned}
& e_{1}=\frac{f_{x_{1}}}{\sqrt{g_{11}}}, \\
& e_{2}=\frac{\sqrt{g_{11}}\left(f_{x_{2}}-\frac{g_{12}}{g_{11}} f_{x_{1}}\right)}{\sqrt{g_{11} g_{22}-g_{12}^{2}}},
\end{aligned}
$$

By Theorem ${ }^{\stackrel{\mathrm{bx}}{5}} \mathbf{5}$.3, there exists new local coordinate system $\tilde{f}\left(y_{1}, y_{2}\right)$ so that $\frac{\partial \tilde{f}}{\partial x_{1}}$ and $\frac{\partial \tilde{f}}{\partial x_{2}}$ are parallel to $e_{1}$ and $e_{2}$. So the first fundamental form written in this coordinate system has the form

$$
\tilde{g}_{11} d y_{1}^{2}+\tilde{g}_{22} d y_{2}^{2} .
$$

However, in general we can not find coordinate system $\hat{f}\left(y_{1}, y_{2}\right)$ so that $e_{1}$ and $e_{2}$ are coordinate vector fields $\frac{\partial \hat{f}}{\partial y_{1}}$ and $\frac{\partial \hat{f}}{\partial y_{2}}$ because if we can then the first fundamental form of the surface is $\mathrm{I}=d y_{1}^{2}+d y_{2}^{2}$, which implies that the Gaussian curvature of the surface must be zero.

### 5.4. Line of curvature coordinates.

Suppose $M$ is an embedded surface in $\mathbb{R}^{3}, f: \mathcal{O} \rightarrow U \subset M$ a local coordinate system, and $p_{0}=f\left(x^{0}\right)$ is not an umbilic point. This means that the shape operator $d N_{p_{0}}$ has two distinct eigenvalues. Since $d N$ is smooth, there exists an open subset $\mathcal{O}_{0}$ of $\mathcal{O}$ containing $x^{0}$ so that $f(x)$ is not umbilic for all $x \in \mathcal{O}_{0}$. We can use linear algebra to write down a formula for two eigenvectors $v_{1}, v_{2}$ of $d N$ on $\mathcal{O}_{0}$ and see that $v_{1}, v_{2}$ are smooth on $\mathcal{O}_{0}$. Since $d N_{p}$ is self-adjoint with distinct eigenvalues, $v_{1}$ and $v_{2}$ are perpendicular. By Theorem 5.3.1, we can change coordinates so that the coordinate vector fields are parallel to $v_{1}$ and $v_{2}$. This means that we may change coordinates $\tilde{f}: \mathcal{O}_{1} \rightarrow U_{1} \subset M \subset \mathbb{R}^{3}$ so that $\frac{\partial \tilde{f}}{\partial y_{1}}$ and $\frac{\partial \tilde{f}}{\partial y_{2}}$ are parallel to $v_{1}, v_{2}$ respectively. In this new coordinates,

$$
\tilde{g}_{12}=\tilde{f}_{y_{1}} \cdot \tilde{f}_{y_{2}}=0
$$

because $v_{1} \cdot v_{2}=0$. But since $(\tilde{f})_{y_{1}}$ and $(\tilde{f})_{y_{2}}$ are in the principal directions, the shape operator is diagonalized with respect to this eigenbasis. Hence

$$
\tilde{\ell}_{12}=0 .
$$

In other words, both I, II are diagonalized. We call such coordinates line of curvature coordinates.

Therefore we have proved
cm Theorem 5.4.1. If $M \subset \mathbb{R}^{3}$ is an embedded surface and $p_{0} \in M$ is not umbilic, then there exists a local line of curvature coordinate system $f: \mathcal{O} \rightarrow U \subset M \subset$ near $p_{0}$, i.e.,

$$
g_{12}=f_{x_{1}} \cdot f_{x_{2}}=0, \quad \ell_{12}=f_{x_{1} x_{2}} \cdot N=0
$$

or equivalently,

$$
\mathrm{I}=A_{1}^{2} d x_{1}^{2}+A_{2}^{2} d x_{2}^{2}, \quad \mathrm{II}=\ell_{11} d x_{1}^{2}+\ell_{22} d x_{2}^{2} .
$$

Theorem 5.4.2. Suppose $f\left(x_{1}, x_{2}\right)$ is a local line of curvature coordinate system on an embedded surface $M$ in $\mathbb{R}^{3}$ with I, II as in Theorem

1cm 5.4.1. Then the Gauss-Codazzi equation is

$$
\begin{align*}
& \left(\frac{\left(A_{1}\right)_{x_{2}}}{A_{2}}\right)_{x_{2}}+\left(\frac{\left(A_{2}\right)_{x_{1}}}{A_{1}}\right)_{x_{1}}=-\frac{\ell_{11} \ell_{22}}{A_{1} A_{2}}  \tag{5.4.1}\\
& \left(\frac{\ell_{11}}{A_{1}}\right)_{x_{2}}=\frac{\ell_{22}\left(A_{1}\right)_{x_{2}}}{A_{2}^{2}}  \tag{5.4.2}\\
& \left(\frac{\ell_{22}}{A_{2}}\right)_{x_{1}}=\frac{\ell_{11}\left(A_{2}\right)_{x_{1}}}{A_{1}^{2}} \tag{5.4.3}
\end{align*}
$$

by: 2
by:3
$P_{\mathrm{b}} \mathrm{b}_{\mathrm{W}}$ of. This theorem can be proved either by substituting $\ell_{12}=0$ into (5.3.2) or derive directly. But we will derive it directly. Let $e_{1}, e_{2}, e_{3}$ be the o.n. basis so that

$$
e_{1}=\frac{f_{x_{1}}}{A_{1}}, \quad e_{2}=\frac{f_{x_{2}}}{A_{2}}, \quad e_{3}=N
$$

and $\left(e_{i}\right)_{x_{1}} \cdot e_{j}=p_{j i},\left(e_{i}\right)_{x_{2}} \cdot e_{j}=q_{i j}$. A direct computation gives

$$
\left\{\begin{array}{c}
\left(e_{1}, e_{2}, e_{3}\right)_{x_{1}}=\left(e_{1}, e_{2}, e_{3}\right)\left(\begin{array}{ccc}
0 & \frac{\left(A_{1}\right)_{x_{2}}}{A_{2}} & -\frac{\ell_{11}}{A_{1}} \\
-\frac{\left(A_{1}\right) x_{2}}{A_{2}} & 0 & 0 \\
\frac{\ell_{11}}{A_{1}} & 0 & 0
\end{array}\right)  \tag{5.4.4}\\
\left(e_{1}, e_{2}, e_{3}\right)_{x_{2}}=\left(e_{1}, e_{2}, e_{3}\right)\left(\begin{array}{ccc}
0 & -\frac{\left(A_{2}\right)_{x_{1}}}{A_{1}} & 0 \\
\frac{\left(A_{2}\right)_{x_{1}}}{A_{1}} & 0 & -\frac{\ell_{22}}{A_{2}} \\
0 & \frac{\ell_{22}}{A_{2}} & 0
\end{array}\right)
\end{array}\right.
$$

The Gauss-Codazzi equation is $P_{x_{2}}-Q_{x_{1}}=[P, Q]$. Since $P, Q$ are skew symmetric, we get 3 equations (the 12 , 13, and 23 entries) as stated in the theorem.

## 6. Surfaces in $\mathbb{R}^{3}$ with $K=-1$

We will show that there exists line of curvature coordinates $f(x, y)$ so that the angle $u$ between asymptotic lines satisfies the sine-Gordon equation (SGE):

$$
u_{x x}-u_{y y}=\sin u
$$

In fact, we prove that there is a one to one correspondence between local solutions $u$ of the SGE with $\operatorname{Im}(u) \subset(0, \pi)$ and local surfaces in $\mathbb{R}^{3}$ with $K=-1$ up to rigid motions. We also prove
(1) the classical Bäcklund transformations that generate infinitely many surfaces of $K=-1$ in $\mathbb{R}^{3}$ from a given one by solving a system of compatible ODEs.
(2) the Hilbert Theorem that there is no complete $K=-1$ surfaces in $\mathbb{R}^{3}$.

## 6.1. $K=-1$ surfaces in $\mathbb{R}^{3}$ and the sine-Gordon equation.

Suppose $M$ be a surface in $\mathbb{R}^{3}$ with $K=-1$, and $\lambda_{1}, \lambda_{2}$ are the two principal curvatures. Since $K=\lambda_{1} \lambda_{2}=-1, \lambda_{1} \neq \lambda_{2}$, i.e., there are no umbilic points. Hence we can find local line of curvature coordinates $f\left(x_{1}, x_{2}\right)$ on $M$. Assume the fundamental forms are

$$
\mathrm{I}=A_{1}^{2} d x_{1}^{2}+A_{2}^{2} d x_{2}^{2}, \quad \mathrm{II}=\ell_{11} d x_{1}^{2}+\ell_{22} d x_{2}^{2}
$$

Then the prinicpal curvatures are

$$
\lambda_{1}=\frac{\ell_{11}}{A_{1}^{2}}, \quad \lambda_{2}=\frac{\ell_{22}}{A_{2}^{2}} .
$$

But $\lambda_{1} \lambda_{2}=-1$, so we may assume that there exists a smooth function $q$ so that

$$
\lambda_{1}=\frac{\ell_{11}}{A_{1}^{2}}=\tan q, \quad \lambda_{2}=\frac{\ell_{22}}{A_{2}^{2}}=-\cot q,
$$

i.e.,

$$
\begin{equation*}
\frac{\ell_{11}}{A_{1}}=A_{1} \tan q, \quad \frac{\ell_{22}}{A_{2}}=-A_{2} \cot q . \tag{6.1.1}
\end{equation*}
$$

Co
Since $A_{1}, A_{2}, \ell_{11}, \ell_{22}$ satisfy the Gauss-Codazzi equation, substitue ( ${ }_{6}^{\text {čo }} \mathrm{f} .1 .1$ ) to the Codazzi equations $\left(\frac{b y .2}{5.2}\right)$ and (5.4.3) to get

$$
\begin{aligned}
& \left(A_{1} \tan q\right)_{x_{2}}=-\cot q\left(A_{1}\right)_{x_{2}}, \\
& \left(-A_{2} \cot q\right)_{x_{1}}=\tan q\left(A_{2}\right)_{x_{1}} .
\end{aligned}
$$

Compute directly to get

$$
\left(A_{1}\right)_{x_{2}} \tan q+A_{1} \sec ^{2} q q_{x_{2}}=-\cot q\left(A_{1}\right)_{x_{2}},
$$

which implies that

$$
(\tan q+\cot q)\left(A_{1}\right)_{x_{2}}=-A_{1}\left(\sec ^{2} q\right) q_{x_{2}}
$$

So we obtain

$$
\frac{\left(A_{1}\right)_{x_{2}}}{A_{1}}=-\frac{\sin q}{\cos q} q_{x_{2}} .
$$

Use a similar computation to get

$$
\frac{\left(A_{2}\right)_{x_{1}}}{A_{2}}=\frac{\cos q}{\sin q} q_{x_{1}} .
$$

In other words, we have

$$
\left(\log A_{1}\right)_{x_{2}}=(\log \cos q)_{x_{2}}, \quad\left(\log A_{2}\right)_{x_{1}}=(\log \sin q)_{x_{1}}
$$

Hence there exist $c_{1}\left(x_{1}\right)$ and $c_{2}\left(x_{2}\right)$ so that

$$
\log A_{1}=\log \cos q+c_{1}\left(x_{1}\right), \quad \log A_{2}=\log \sin q+c_{2}\left(x_{2}\right)
$$

i.e.,

$$
A_{1}=e^{c_{1}\left(x_{1}\right)} \cos q, \quad A_{2}=e^{c_{2}\left(x_{2}\right)} \sin q .
$$

Because I is positive definite, $A_{1}, A_{2}$ never vanishes. So we may assume both $\sin q$ and $\cos q$ are positive, i.e., $q \in\left(0, \frac{\pi}{2}\right)$. Now change coordinates to ( $\tilde{x}_{1}\left(x_{1}\right), \tilde{x}_{2}\left(x_{2}\right)$ ) so that

$$
\frac{d \tilde{x}_{1}}{d x_{1}}=e^{c_{1}\left(x_{1}\right)}, \quad \frac{d \tilde{x}_{2}}{d x_{2}}=e^{c_{2}\left(x_{2}\right)}
$$

Since

$$
\frac{\partial f}{\partial \tilde{x}_{1}}=\frac{\partial f}{\partial x_{1}} \frac{\partial x_{1}}{\partial \tilde{x}_{1}}=\frac{\partial f}{\partial x_{1}} e^{-c_{1}\left(x_{1}\right)}
$$

$\left\|f_{\tilde{x}_{1}}\right\|=\cos q$. Similar calculation implies that $\left\|f_{\tilde{x}_{2}}\right\|=\sin q$. Since $\tilde{x}_{i}$ is a function of $\tilde{x}_{i}$ alone, $f_{\tilde{x}_{i}}$ is parallel to $f_{x_{i}}$. So $\left(\tilde{x}_{1}, \tilde{x}_{2}\right)$ is also a line of curvature coordinate system and the coefficients of II in ( $\left.\tilde{x}_{1}, \tilde{x}_{2}\right)$ coordinate system is

$$
\tilde{\ell}_{11}=\tan q \cos ^{2} q=\sin q \cos q, \quad \tilde{\ell}_{22}=-\cot q \sin ^{2} q=-\sin q \cos q .
$$

Therefore, we have proved part of the following Theorem:
cq Theorem 6.1.1. Let $M^{2}$ be a surface in $\mathbb{R}^{3}$ with $K=-1$. Then locally there exists line of curvature coordinates $x_{1}, x_{2}$ so that

$$
\begin{equation*}
\mathrm{I}=\cos ^{2} q d x_{1}^{2}+\sin ^{2} q d x_{2}^{2}, \quad \mathrm{II}=\sin q \cos q\left(d x_{1}^{2}-d x_{2}^{2}\right) \tag{6.1.2}
\end{equation*}
$$

where $2 q$ is the angle between two asymptotic directions. Moreover, the Gauss-Codazzi equation is the sine-Gordon equation

$$
\begin{equation*}
q_{x_{1} x_{1}}-q_{x_{2} x_{2}}=\sin q \cos q . \tag{6.1.3}
\end{equation*}
$$

Proof. It remains to compute the Gauss-Codazzi equation. Let

$$
e_{1}=\frac{f_{x_{1}}}{\cos q}, \quad e_{2}=\frac{f_{x_{2}}}{\sin q}, \quad e_{3}=N=\frac{f_{x_{1}} \times f_{x_{2}}}{\left\|f_{x_{1}} \times f_{x_{2}}\right\|}
$$

Write

$$
\left(e_{1}, e_{2}, e_{3}\right)_{x_{1}}=\left(e_{1}, e_{2}, e_{3}\right) P, \quad\left(e_{1}, e_{2}, e_{3}\right)_{x_{2}}=\left(e_{1}, e_{2}, e_{3}\right) Q
$$

We have given formulas for $P, Q$ in $\left(\frac{\mathrm{cg}}{5.4 .4)}\right.$ when $\left(x_{1}, x_{2}\right)$ is a line of curvature coordinate system. Since $A_{1}=\cos q, A_{2}=\sin q$, and $\ell_{11}=$ $-\ell_{22}=\sin q \cos q$, we have

$$
P=\left(\begin{array}{ccc}
0 & -q_{x_{2}} & -\sin q  \tag{6.1.4}\\
q_{x_{2}} & 0 & 0 \\
\sin q & 0 & 0
\end{array}\right), \quad Q=\left(\begin{array}{ccc}
0 & -q_{x_{1}} & 0 \\
q_{x_{1}} & 0 & \cos q \\
0 & -\cos q & 0
\end{array}\right)
$$

A direct computation shows that Codazzi equations are satified automatically, i.e., the 13 and 23 entries of $P_{x_{2}}-Q_{x_{1}}$ are equal to the 13 and 23 entries of $[P, Q]$ respectively. The Gauss equation (equating the 12 entry) gives

$$
-q_{x_{2} x_{2}}+q_{x_{1} x_{1}}=\sin q \cos q
$$

Since $I I=\sin q \cos q\left(d x_{1}^{2}-d x_{2}^{2}\right), f_{x_{1}} \pm f_{x_{2}}$ are asymptotic directions. Use I to see that $f_{x_{1}} \pm f_{x_{2}}$ are unit vectors. Since

$$
\left(f_{x_{1}}+f_{x_{2}}\right) \cdot\left(f_{x_{1}}-f_{x_{2}}\right)=\cos ^{2} q-\sin ^{2} q=\cos (2 q),
$$

the angle between the asymptotic directions $f_{x_{1}}+f_{x_{2}}$ and $f_{x_{1}}-f_{x_{2}}$ is $2 q$.

A consequence of the proof of the above theorem is
cr Corollary 6.1.2. Let $P, Q$ be as in ( $\left.{ }^{\mathrm{c} p} 1.4\right)$. Then system

$$
\left(e_{1}, e_{2}, e_{3}\right)_{x_{1}}=\left(e_{1}, e_{2}, e_{3}\right) P, \quad\left(e_{1}, e_{2}, e_{3}\right)_{x_{2}}=\left(e_{1}, e_{2}, e_{3}\right) Q
$$

is solvable if and only if $q$ satisfies the $S G E$ ( ${ }^{\mathrm{bz}} 1.3$ ).
Let $O(3)$ denote the space of all orthogonal $3 \times 3$ matrices, and $o(3)$ the space of all skew-symmetric $3 \times 3$ matrices. A map $g: \mathbb{R}^{2} \rightarrow O(3)$ is smooth if $i \circ g: \mathbb{R}^{2} \rightarrow g l(3)$ is smooth, where $i: O(3) \rightarrow g l(3)$ is the inclusion.

It follows from the Fundamental Theorem of surfaces and Corollary ${ }_{6}^{\mathrm{cr}} \mathrm{b} 1.2$ that The converse of Theorem $\frac{{ }^{\mathrm{Cq}} \mathrm{Cq}}{6.1 .1}$ is true :
Theorem 6.1.3. Let $q: \mathcal{O} \rightarrow \mathbb{R}$ be a solution of the $S G E$ ( ${ }_{6}^{\text {bz. }} 1.3$ ), $p_{0} \in \mathbb{R}^{3},\left(x_{1}^{0}, x_{2}^{0}\right) \in \mathcal{O}$, and $u_{\mathrm{cP}}, u_{2}, u_{3}$ an o.n. basis. Let $P, Q: \mathcal{O} \rightarrow o(3)$ be the maps defined by (6.1.4). Then there exists an open subset $\mathcal{O}_{1}$ of $\left(x_{1}^{0}, x_{2}^{0}\right)$ in $\mathcal{O}$ and a unique solution $\left(f, e_{1}, e_{2}, e_{3}\right): \mathcal{O}_{1} \rightarrow \mathbb{R}^{3} \times O(3)$ for the following system

$$
\left\{\begin{array}{l}
\left(e_{1}, e_{2}, e_{3}\right)_{x_{1}}=\left(e_{1}, e_{2}, e_{3}\right) P  \tag{6.1.5}\\
\left(e_{1}, e_{2}, e_{3}\right)_{x_{2}}=\left(e_{1}, e_{2}, e_{3}\right) Q \\
f_{x_{1}}=\cos q e_{1}, \\
f_{x_{2}}=\sin q e_{2}, \\
f\left(x_{1}^{0}, x_{2}^{0}\right)=p_{0}, \quad e_{1}\left(x_{1}^{0}, x_{2}^{0}\right)=u_{1}, \quad e_{2}\left(x_{1}^{0}, x_{2}^{0}\right)=u_{2}
\end{array}\right.
$$

Moreover, if $\sin q \cos q>0$ on $\mathcal{O}_{1}$, then $f\left(\mathcal{O}_{1}\right)$ is an immersed surface with $K=-1$ and the two fundamental forms are of the form (бै.1.2).

In other words, there is a $1-1$ correspondence between solutions $q$ of the $\operatorname{SGE}(6.1 .3)$ with $\operatorname{Im}(q) \subset\left(0, \frac{\pi}{2}\right)$ and local surfaces of $\mathbb{R}^{3}$ with $K=-1$ up to rigid motion.

Let $f\left(x_{1} x_{2}\right)$ be a local line of curvature coordinate system given in Theorem 6.1 .1 of an embedded surface $M_{\mathrm{bz}}$ in $\mathbb{R}^{3}$ with $K=-1$, and $q$ the corresponding solution of the SGE (6.1.3). The SGE is a nonlinear wave equation, and $\left(x_{1}, x_{2}\right)$ is the space-time coordinate. We have proved that $f_{x_{1}} \pm f_{x_{2}}$ are asymptotic directions. If we make a change of coordinates $x_{1}=s+t, x_{2}=s-t$, then $f_{s}=f_{x_{1}}+f_{x_{2}}$ and $f_{t}=f_{x_{1}}-f_{x_{2}}$. A direct computation shows that the two fundamental forms written in $(s, t)$ coordinates are

$$
\left\{\begin{array}{l}
\mathrm{I}=d s^{2}+2 \cos 2 q d s d t+d t^{2}  \tag{6.1.6}\\
\mathrm{II}=\sin 2 q d s d t
\end{array}\right.
$$

and the SGE ( $\left(\frac{\mathrm{bz}}{6.1} 1.3\right)$ becomes

$$
\begin{equation*}
2 q_{s t}=\sin (2 q) . \tag{6.1.7}
\end{equation*}
$$

A local coordinate system $(x, y)$ on a surface $f(x, y)$ in $\mathbb{R}^{3}$ is called an asymptotic coordinate system if $f_{x}, f_{y}$ are parallel to the asymptotic direction, i.e., $\ell_{11}=\ell_{22}=0$. Note that the $(s, t)$ coordinate constructed above for $K=-1$ surfaces in $\mathbb{R}^{3}$ is an asymptotic coordinate system and $s, t$ are arc length parameter. This coordinate system is called the Tchbyshef coordinate system for $K=-1$ surfaces in $\mathbb{R}^{3}$.

### 6.2. Moving frame method.

If $P, Q$ are skew-symmetric, then so is $[P, Q]=P Q-Q P$ because $(P Q-Q P)^{t}=Q^{t} P^{t}-P^{t} Q^{t}=(-Q)(-P)-(-P)(-Q)=-(P Q-Q P)$.

We use Cartan's method of moving frames to derive the local theory of surfaces in $\mathbb{R}^{3}$. This method uses differential forms instead of vector fields to derive the Gauss-Codazzi equation.
ds Proposition 6.2.1. Let $\mathcal{O}$ be an open subset of $\mathbb{R}^{2}$, and $P=\left(p_{i j}\right), Q=$ $\left(q_{i j}\right)$ smooth maps from $\mathcal{O}$ to o(3). Let $w_{i j}=p_{i j} d x_{1}+q_{i j} d x_{2}$. Then the following statements are equivalent:
(i) $\left\{\begin{array}{l}g_{x_{1}}=g P, \\ g_{x_{2}}=g Q\end{array} \quad\right.$ is solvable,
(ii) $-P_{x_{2}}+Q_{x_{1}}=[P, Q]$,
(iii)

$$
\begin{equation*}
d w_{i j}=-\sum_{k=1}^{3} w_{i k} \wedge w_{k j} \tag{6.2.1}
\end{equation*}
$$

Proof. The equivalence of (i) and (ii) is given by Corollary ${ }_{2}^{\mathrm{bc}} 4.2$. It remains to prove (ii) if and only if (iii). A direct computation gives

$$
\begin{aligned}
& d w_{i j}=d\left(p_{i j} d x_{1}+q_{i j} d x_{2}\right)=\left(-\left(p_{i j}\right)_{x_{2}}+\left(q_{i j}\right)_{x_{1}}\right) d x_{1} \wedge d x_{2}, \quad \text { and } \\
& -\sum_{k} w_{i k} \wedge w_{k j}=-\sum_{k}\left(p_{i k} d x_{1}+q_{i k} d x_{2}\right) \wedge\left(p_{k j} d x_{1}+q_{k j} d x_{2}\right) \\
& \quad=-\sum_{k}\left(p_{i k} q_{k j}-q_{i k} p_{k j}\right) d x_{1} \wedge d x_{2} \\
& \quad=-[P, Q]_{i j} d x_{1} \wedge d x_{2} .
\end{aligned}
$$

So (ii) and (iii) are equivalent.
We introduce some notations. Let $w_{i j}$ be 1 -forms for $1 \leq i, j \leq 3$. Then $w=\left(w_{i j}\right)$ is called a $g l(3)$-valued 1 -form. Let $\tau=\left(\tau_{i j}\right)$ be a $g l(3)$ valued 1-form. Let $\tau \wedge w$ denote the $g l(3)$-valued 1 -form whose $i j$-th entry is $\sum_{k=1}^{3} \tau_{i k} \wedge w_{k j}$. Use matrix notation, ( $\frac{\mathrm{du} .2 .1 \text { ) can be written as }}{}$

$$
\begin{equation*}
d w=-w \wedge w \tag{6.2.2}
\end{equation*}
$$

which is called the Maurer-Cartan equation.
Once we get used to computations involving matrix valued maps and differential forms, we can obtain the Maurer-Cartan equation easily as follows: Because $w=g^{-1} d g$, we have

$$
d w=d\left(g^{-1}\right) \wedge d g+g^{-1} d(d g)=-g^{-1} d g g^{-1} \wedge d g+0=-w \wedge w
$$

i.e., $w$ satisfies $\left(\frac{\mathrm{dt}}{6.2 .2}\right)$.

Since $w_{i j}=-w_{j i}, w_{i i}=0$. So $\left(\frac{\mathrm{du}}{\hbar .2 .1}\right)$ becomes

$$
\left\{\begin{array}{l}
d w_{12}=-w_{13} \wedge w_{32}  \tag{6.2.3}\\
d w_{13}=-w_{12} \wedge w_{23} \\
d w_{23}=-w_{21} \wedge w_{13}
\end{array}\right.
$$

dy Lemma 6.2.2. Cartan's Lemma If $\theta_{1}, \theta_{2}$ are linearly independent 1forms on an open subset $\mathcal{O}$ of $\mathbb{R}^{2}$ at each point of $\mathcal{O}$, then there exists a unique $\theta_{12}$ satisfies

$$
\left\{\begin{array}{l}
d \theta_{1}=-\theta_{12} \wedge \theta_{2},  \tag{6.2.4}\\
d \theta_{2}=-\theta_{21} \wedge \theta_{1},
\end{array} \quad \text { where } \theta_{12}+\theta_{21}=0\right.
$$

Proof. Suppose $\theta_{i}=A_{i} d x+B_{i} d y$ with $i=1,2$, and $\theta_{12}=p d x+q d y$. Then (6.2.4) implies that

$$
\left\{\begin{array}{l}
-p B_{2}+q A_{2}=-\left(A_{1}\right)_{y}+\left(B_{1}\right)_{x}  \tag{6.2.5}\\
p B_{1}-q A_{1}=-\left(A_{2}\right)_{y}+\left(B_{2}\right)_{x}
\end{array}\right.
$$

db

Since $\theta_{1}, \theta_{2}$ are linearly independent at every point, $A_{1} B_{2}-A_{2} B_{1} \neq 0$. So the linear system (6.2.5) has a unique solution $p, q$.

Let $f: \mathcal{O} \rightarrow U \subset M$ be a local coordinate system of the embedded surface $M$ in $\mathbb{R}^{3}$. When we view $f$ as a map from $\mathcal{O}$ to $\mathbb{R}^{3}$, the differential of $f$ is

$$
d f=f_{x_{1}} d x_{1}+f_{x_{2}} d x_{2} .
$$

Let $\pi_{i}: M \rightarrow \mathbb{R}^{3}$ be the function defined by $\pi_{i}\left(y_{1}, y_{2}, y_{3}\right)=y_{i}$. The differential

$$
\left(d \pi_{i}\right)_{f(q)}=d\left(f_{i}\right)_{q} .
$$

Let $\pi=\left(\pi_{1}, \pi_{2}, \pi_{3}\right): M \rightarrow \mathbb{R}^{3}$. Then $\pi$ is the inclusion map, and

$$
d \pi_{f(q)}=d f_{q}=f_{x_{1}} d x_{1}+f_{x_{2}} d x_{2}
$$

Since $d x_{1}, d x_{2}$ are 1-forms dual to the basis $f_{x_{1}}, f_{x_{2}}$ of $T M$, the restriction of $d \pi$ to each $T M_{p}$ is the identity map. Suppose $e_{1}, e_{2}$ are smooth orthonormal vector fields of $M$ defined on $f(\mathcal{O})$, and $w_{1}, w_{2}$ are dual 1-form on $M$, i. $e_{\text {dp }} w_{i}\left(e_{j}\right)=\delta_{i j}$. Since $d f_{q}=d \pi_{f(q)}$ is the identiy map, by Corollary 3.1.7 we have

$$
d f=\theta_{1} e_{1}+\theta_{2} e_{2} .
$$

Let $e_{3}=N$ be the unit normal vector field. Set

$$
\begin{equation*}
w_{i j}=d e_{i} \cdot e_{j}, \quad 1 \leq i, j \leq 3 \tag{6.2.6}
\end{equation*}
$$

dw
Let $g=\left(e_{1}, e_{2}, e_{3}\right)$. Rewrite $(\stackrel{\mathrm{d} \mathbf{w}}{6.2 .6})$ as

$$
d g=g w .
$$

Differentiate $e_{i} \cdot e_{j}=\delta_{i j}$ to get

$$
w_{i j}+w_{j i}=0
$$

i.e., $w=\left(w_{i j}\right)$ is a $o(3)$-valued 1-form. The Gauss-Codazzi equation is the compatibility condition for $d g=g w$. So it follows from Proposition
$\frac{\mathrm{ds}}{6.2 .1}$ that the Gauss-Codazzi equation for $M$ is the Maurer-Cartan equation $d w=-w \wedge w$. Since $w_{i i}=0$, the Maurer-Cartan equation is

$$
\left\{\begin{array}{l}
d w_{12}=-w_{13} \wedge w_{32}  \tag{6.2.7}\\
d w_{13}=-w_{12} \wedge w_{23} \\
d w_{23}=-w_{21} \wedge w_{13}
\end{array}\right.
$$

$$
\begin{aligned}
& \text { Differentiate } d f=w_{1} e_{1}+w_{2} e_{2} \text { to get } \\
& 0=d(d f)=d w_{1} e_{1}-w_{1} \wedge d e_{1}+d w_{2} e_{2}-w_{2} \wedge d e_{2} \\
& =d w_{1} e_{1}-w_{1} \wedge\left(w_{21} e_{2}+w_{31} e_{3}\right)+d w_{2} e_{2}-w_{2} \wedge\left(w_{12} e_{1}+w_{32} e_{3}\right) \\
& =\left(d w_{1}-w_{2} \wedge w_{12}\right) e_{1}+\left(d w_{2}-w_{1} \wedge w_{21}\right) e_{2}-\left(w_{1} \wedge w_{31}+w_{2} \wedge w_{32}\right) e_{3} .
\end{aligned}
$$

So we have

$$
\begin{equation*}
\left(d w_{1}-w_{2} \wedge w_{12}\right) e_{1}+\left(d w_{2}-w_{1} \wedge w_{21}\right) e_{2}-\left(w_{1} \wedge w_{31}+w_{2} \wedge w_{32}\right) e_{3}=0 \tag{6.2.8}
\end{equation*}
$$

The coefficients of $e_{1}, e_{2}$ are zero imply that

$$
\left\{\begin{array}{l}
d w_{1}-w_{2} \wedge w_{12}=0  \tag{6.2.9}\\
d w_{2}-w_{1} \wedge w_{21}=0
\end{array}\right.
$$

where $w_{12}=-w_{21}$. Equation $\binom{(6 y y}{6.9}$ is called the structure equation. It follows from the Cartan Lemma $\sqrt{6.2 .2}$ that $w_{12}$ can be computed in terms of $w_{1}, w_{2}$. But the first fundamental form is

$$
\mathrm{I}=w_{1} \otimes w_{1}+w_{2} \otimes w_{2}
$$

So this proves that $w_{12}$ only depend dx on I .
The coefficient of $e_{3}$ in equation (6.2.8) is zero implies that

$$
\begin{equation*}
w_{1} \wedge w_{31}+w_{2} \wedge w_{32}=0 \tag{6.2.10}
\end{equation*}
$$

Write

$$
w_{31}=h_{11} w_{1}+h_{12} w_{2}, \quad w_{32}=h_{21} w_{1}+h_{22} w_{2}
$$

Then ( ${ }_{6}^{\text {čZ }} .2 .10$ ) implies that

$$
h_{12} w_{1} \wedge w_{2}+h_{21} w_{2} \wedge w_{1}=0
$$

So $h_{12}=h_{21}$. The shape operator is

$$
-d N=-d e_{3}=-\left(w_{13} e_{1}+w_{23} e_{2}\right)=w_{31} e_{1}+w_{32} e_{2}=\sum_{i j} h_{i j} w_{i} e_{j}
$$

Because $e_{1}, e_{2}$ are orthonormal, $h_{1 \mathrm{a}_{\mathrm{a}}}=h_{21}$ means that $d N$ is self-adjoint. This gives a proof of Proposition 3.3 .1 via differential forms.

By definition of II, we see that

$$
\mathrm{II}=\sum_{i j} h_{i j} w_{i} \otimes w_{j}
$$

The mean and the Gaussian curvature are

$$
H=h_{11}+h_{22}, \quad K=h_{11} h_{22}-h_{12}^{2}
$$

The first equation of the Gauss-Codazzi equation $\left(\frac{\mathrm{da}}{6.2 .3)}\right.$ is

$$
\begin{aligned}
d w_{12} & =-w_{13} \wedge w_{32}=w_{31} \wedge w_{32}=\left(h_{11} w_{1}+h_{12} w_{2}\right) \wedge\left(h_{12} w_{1}+h_{22} w_{2}\right) \\
& =\left(h_{11} h_{22}-h_{12}^{2}\right) w_{1} \wedge w_{2}=K w_{1} \wedge w_{2} .
\end{aligned}
$$

So we have

$$
d w_{12}=K w_{1} \wedge w_{2}
$$

But we have proved that $w_{12}$ can be computed from $w_{1}, w_{2}$. So $K$ deppends only on I. This gives another proof of the Gauss Theorem 5.2.1 using differential forms.

We summerize Cartan's moving frame method. Let $M$ be an embedded surface in $\mathbb{R}^{3}, e_{1}, e_{2}$ be a local o.n. tangent frame on an open subset $U$ of $M$, and $w_{1}, w_{2}$ the dual 1-form. Let $e_{3}$ be the unit normal field on $M$, and

$$
w_{i j}=d e_{j} \cdot e_{i}, \quad 1 \leq i, j \leq 3
$$

Then:
(1) The first and second fundametnal forms of $M$ are

$$
\left\{\begin{array}{l}
\mathrm{I}=w_{1}^{2}+w_{2}^{2} \\
\mathrm{II}=w_{1} \otimes w_{31}+w_{2} \otimes w_{32}
\end{array}\right.
$$

(2) $w_{1}, w_{2}, w_{i j}$ satisfy the structure equation

$$
\left\{\begin{array}{l}
d w_{1}=-w_{12} \wedge w_{2}  \tag{6.2.11}\\
d w_{2}=-w_{21} \wedge w_{1}
\end{array}\right.
$$

and the Gauss and Codazzi equation

$$
\left\{\begin{array}{l}
d w_{12}=-w_{13} \wedge w_{32}  \tag{6.2.12}\\
d w_{13}=-w_{12} \wedge w_{23} \\
d w_{23}=-w_{21} \wedge w_{23}
\end{array}\right.
$$

(3) Write $w_{3 i}=h_{i 1} w_{1}+h_{i 2} w_{2}$. Then $h_{i j}=h_{j i}$, hence II is symmetric.

$$
\begin{equation*}
d w_{12}=-w_{13} \wedge w_{32}=K w_{1} \wedge w_{2} \tag{4}
\end{equation*}
$$

ev
(5) $w_{12}$ only depend on I and can be solved using the structure equation $\left(\frac{\mathrm{eg}}{6.2 .11)}\right.$ ) and the Gaussian curvature $K$ can be computed using I, namely

$$
d w_{12}=K w_{1} \wedge w_{2}
$$

Next we state the Fundamental Theorem of surfaces in $\mathbb{R}^{3}$ in terms of differential forms. Given two symmetric bilinar forms

$$
\mathrm{I}=\sum_{i j} g_{i j} d x_{i} d x_{j}, \quad \mathrm{II}=\sum_{i j} \ell_{i j} d x_{i} d x_{j}
$$

on an open subset $\mathcal{O}$ of $\mathbb{R}^{2}$ such that $I$ is positive definite. First, we complete the square for I to get an orthonormal dual 1-form:

$$
\mathrm{I}=g_{11}\left(d x_{1}+\frac{g_{12}}{g_{11}} d x_{2}\right)^{2}+\frac{g_{11} g_{22}-g_{12}^{2}}{g_{11}} d x_{2}^{2}
$$

Then $\mathrm{I}=w_{1}^{2}+w_{2}^{2}$, where

$$
w_{1}=\sqrt{g_{11}}\left(d x_{1}+\frac{g_{12}}{g_{11}} d x_{2}\right), \quad w_{2}=\sqrt{\frac{g_{11} g_{22}-g_{12}^{2}}{g_{11}}} d x_{2} .
$$

Let $w_{12}$ be the unique 1-form satisfies the structure equation (use the proof of Cartan's lemma to find the formula of $w_{12}$ ). Note that $w_{31}$ and $w_{32}$ can be solved algebraically by writing

$$
\mathrm{II}=\sum_{i j} \ell_{i j} d x_{i} d x_{j}=w_{1} \otimes w_{31}+w_{2} \otimes w_{32}
$$

Suppose $w_{\mathrm{dh}} w_{2}, w_{i j}$ satisfy the structure $\left({ }^{\mathrm{eg}} \mathrm{f} .2 .11\right)$ and Gauss-Codazzi equation (6.2.12). Then given $q_{0} \in \mathcal{O}, p_{0} \in \mathbb{R}^{3}$, and an o.n. basis $u_{1}, u_{2}, u_{3}$ of $\mathbb{R}^{3}$, we claim that there exists a unique solution $\left(f, e_{1}, e_{2}, e_{3}\right)$ for

$$
\begin{cases}d f=w_{1} e_{1}+w_{2} e_{2}  \tag{6.2.14}\\ d e_{i}=\sum_{j=1}^{3} w_{j i} e_{j}, & 1 \leq i, j \leq 3 \\ \left(f, e_{1}, e_{2}, e_{3}\right)\left(q_{o}\right)=\left(p_{0}, u_{1}, u_{2}, u_{3}\right)\end{cases}
$$

```
\(\square\)
```

with initial data $\left(p_{0}, u_{1}, u_{2}, u_{3}\right)$. To prove the claim, we only need to check the compatibility condition. Note that solvability of the first equation of $\left(\frac{e_{2}}{6.2 .14)}\right.$ is $d\left(w_{1} e_{1}+w_{2} e_{2}\right)=0$, which is the structure equation (6.2.9). The solvability of the second equation of (6.2.14) is the Gauss-Codazzi equation. This proves the claim.

We give an example using Cartan's moving frame method below:
Theorem 6.2.3. Hilbert Theorem for $K=-1$ surfaces in $\mathbb{R}^{3}$ Suppose $M$ is a surface in $\mathbb{R}^{3}$ with $K=-1$ given by an entire Tchebyshef asymptotic coordinate system $(s, t)$, i.e., there is a coordinate system $f: \mathbb{R}^{2} \rightarrow \mathbb{R}^{3}$ with $f\left(\mathbb{R}^{2}\right)_{e} \overline{\bar{j}} M$ so that the first and second fundamental forms are of the form (6.1.6). Then the area of $M$ is less than $2 \pi$.

Proof. We have proved that $\mathrm{I}=d s^{2}+2 \cos 2 q d s d t+d t^{2}$. Complete square to get an o.n. dual 1-forms:

$$
\mathrm{I}=(d s+\cos 2 q d t)^{2}+(\sin 2 q d t)^{2}
$$

Let

$$
w_{1}=d s+\cos 2 q d t, \quad w_{2}=\sin 2 q d t .
$$

We use the structure equation ( ${ }^{(\mathrm{eg}} .2 .11$ ) to compute $w_{12}$. Let

$$
w_{12}=A d s+B d t
$$

Structure equation ( ${ }^{\frac{\mathrm{eg}}{6.2}} 2.11$ ) implies that

$$
\left\{\begin{array}{l}
-2 \sin 2 q q_{s}=-A \sin 2 q  \tag{6.2.15}\\
2 \cos 2 q q_{s}=-B+A \cos 2 q
\end{array}\right.
$$

So $A=2 q_{s}$ and $B=0$, i.e.,

$$
w_{12}=2 q_{s} d s
$$

The Gauss equation is $d w_{12}=K w_{1} \wedge w_{2}$, so we have

$$
\begin{equation*}
d w_{12}=-w_{1} \wedge w_{2}=-\sin q \cos q d s \wedge d t \tag{6.2.16}
\end{equation*}
$$

ei
The area of $M$ is

$$
\begin{aligned}
\operatorname{Area}(M) & =\iint_{\mathbb{R}^{2}} w_{1} \wedge w_{2}=\iint_{\mathbb{R}^{2}}(d s+\cos 2 q d t) \wedge \sin 2 q d t \\
& =\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \sin 2 q d s d t \quad \text { by }\left(\frac{\mathrm{ei}}{6} .2 .16\right) \\
& =\int_{-\infty}^{\infty} \int_{-\infty}^{\infty}-d w_{12}, \quad \text { by the Green's formula (这.3.2) } \\
& =\lim _{r \rightarrow \infty} \int_{-r}^{r} \int_{-r}^{r}-w_{12}=-\lim _{r \rightarrow \infty} \oint 2 q_{s} d s \\
& =-2 \lim _{r \rightarrow \infty} \int_{-r}^{r} q_{s}(s,-r) d s+\int_{r}^{-r} q_{s}(s, r) d s \\
& =-2 \lim _{r \rightarrow \infty}(q(r,-r)-q(-r,-r))+(q(-r, r)-q(r, r))
\end{aligned}
$$

Since I is positive definite, $1-\cos ^{2} 2 q=\sin ^{2} 2 q>0$. We may assume $2 q \in(0, \pi)$. So the above equation implies that the area of $M$ is less than $2 \pi$.

### 6.3. Bäcklund Theorem.

The idea of Bäcklund transformations comes from the study of line congruences. A line congruence in $R^{3}$ is a two-parameter family of lines

$$
L(u, v): \quad x(u, v)+\tau \xi(u, v), \quad-\infty<\tau<\infty
$$

i.e., $L(u, v)$ is the line through $x(u, v)$ and parallel to $\xi(u, v)$. A surface $M$ given by

$$
Y(u, v)=x(u, v)+t(u, v) \xi(u, v)
$$

for some smooth function $t$ is called a focal surface of the line congruence if the line $L(u, v)$ is tangent to $M$ at $Y(u, v)$ for all $(u, v)$. Hence $\xi(u, v)$ lies in the tangent plane of $M$ at $Y(u, v)$, which is spanned by

$$
Y_{u}=x_{u}+t_{u} \xi+t \xi_{u}, \quad Y_{v}=x_{v}+t_{v} \xi+t \xi_{v}
$$

Note that $v_{1}, v_{2}, v_{3}$ are linearly dependent if and only if $\operatorname{det}\left(v_{1}, v_{2}, v_{3}\right)=$ 0 . So $t$ satisfies the following quadratic equation:

$$
\operatorname{det}\left(\xi, x_{u}+t_{u} \xi+t \xi_{u}, x_{v}+t_{v} \xi+t \xi_{v}\right)=0
$$

This implies

$$
\begin{aligned}
0 & =\operatorname{det}\left(\xi, x_{u}+t \xi_{u}, x_{v}+t \xi_{v}\right) \\
& =t^{2} \operatorname{det}\left(\xi, \xi_{u}, \xi_{v}\right)+t\left(\operatorname{det}\left(\xi, \xi_{u}, x_{v}\right)+\operatorname{det}\left(\xi, x_{u}, \xi_{v}\right)\right)+\operatorname{det}\left(\xi, x_{u}, x_{v}\right)
\end{aligned}
$$

Note that $x, \xi$ are given and the above equation is a quadratic equation in $t$. In general, this quadratic equation has two distinct solutions for $t$. Hence generically each line congruence has two focal surfaces, $M$ and $M^{*}$. This results in a diffeomorphism $\ell: M \rightarrow M^{*}$ such that the line joining $p$ and $p^{*}=\ell(p)$ is tangent to both $M$ and $M^{*}$. We will call $\ell$ also a line congruence.

A line congruence $\ell: M \rightarrow M^{*}$ is called a Bäcklund transformation with constant $\theta$ if the angle between the normal of $M$ at $p$ and the normal of $M^{*}$ at $p^{*}=\ell(p)$ is $\theta$ and the distance between $p$ and $p^{*}$ is $\sin \theta$ for all $p \in M$. We restate the definition below:
Definition 6.3.1. Let $M, M^{*}$ be two surfaces in $\mathbb{R}^{3}$. A diffeomorphism $\ell: M \rightarrow M^{*}$ is called a Bäcklund transformation with constant $\theta$ if for any $P \in M$,
(1) the line joing $P$ and $P^{*}=\ell(P)$ is tangent to $M$ at $P$ and to $M^{*}$ at $P^{*}$ respectively,
(2) the distance between $P$ and $P^{*}$ is $\sin \theta$,
(3) the angle between the normal $N_{P}$ of $M$ at $P$ and $N_{P^{*}}^{*}$ of $M^{*}$ at $P^{*}$ is equal to $\theta$.

In 1883, A. Bäcklund proved:
Theorem 6.3.2. Bäcklund Theorem. Let $M$ and $M^{*}$ be two surfaces in $\mathbb{R}^{3}$. If $\ell: M \rightarrow M^{*}$ is a Bäcklund transformation with constant $\theta$, then both $M$ and $M^{*}$ have Gaussian curvature -1 .

Proof. Let $e_{1}$ denote the unit vector field on $M$ so that

$$
\ell(p)=p+\sin \theta e_{1}(p)
$$

for all $p \in M$. Let $e_{2}$ be the unit vector field tangent to $M$ that is perpendicular to $e_{1}$, and $e_{3}=N$ the unit normal vector field on $M$. Then $e_{1}, e_{2}, e_{3}$ is a local o.n. frame on $M$. Let $e_{1}^{*}, e_{2}^{*}$ be the tangent o.n. frame on $M^{*}$ so that $e_{1}^{*}=-e_{1}$ (the direction of $\overrightarrow{\ell(p) p)}$, and $e_{3}^{*}$ is the unit normal of $M^{*}$. By assumption, the angle between $e_{3}$ and $e_{3}^{*}$ is $\theta$. So the two o.n. frames are related by

$$
\left\{\begin{array}{l}
e_{1}^{*}=-e_{1}, \\
e_{2}^{*}=\cos \theta e_{2}+\sin \theta e_{3} \\
e_{3}^{*}=-\sin \theta e_{2}+\cos \theta e_{3}
\end{array}\right.
$$

Let $\left(w_{i j}\right)$ and $\left(w_{i j}^{*}\right)$ be the so(3)-valued 1-form computed from

$$
\begin{equation*}
d e_{i}=\sum_{j} w_{j i} e_{j}, \quad d e_{i}^{*}=\sum_{j} w_{j i}^{*} e_{j}^{*} . \tag{6.3.1}
\end{equation*}
$$

dd

Let $f: \mathcal{O} \rightarrow \mathbb{R}^{3}$ be a local parametrization of $M$, then

$$
f^{*}=f+\sin \theta e_{1}
$$

is a local parametrization of $M^{*}$. Let $w_{1}, w_{2}$ be the dual coframe of $e_{1}, e_{2}$, and $w_{1}^{*}, w_{2}^{*}$ the dual coframe of $e_{1}^{*}, e_{2}^{*}$. So we have

$$
\begin{equation*}
d f=w_{1} e_{1}+w_{2} e_{2}, \quad d f^{*}=w_{1}^{*} e_{1}^{*}+w_{2}^{*} e_{2}^{*} \tag{6.3.2}
\end{equation*}
$$

eu
The relation of these dual coframes can be computed by differentiating $f$ and $f^{*}$ :

$$
\begin{aligned}
d f^{*} & =w_{1}^{*} e_{1}^{*}+w_{2}^{*} e_{2}^{*} \\
& =-w_{1}^{*} e_{1}+w_{2}^{*}\left(\cos \theta e_{2}+\sin \theta e_{3}\right)=d\left(f+\sin \theta e_{1}\right) \\
& =d f+\sin \theta d e_{1} \quad \text { by }(\text { eu. } 6.3 .2) \text { and }\left(\frac{\mathrm{dad}}{6.3 .1)}\right. \\
& =w_{1} e_{1}+w_{2} e_{2}+\sin \theta\left(w_{21} e_{2}+w_{31} e_{3}\right) .
\end{aligned}
$$

Compare coefficients of $e_{1}, e_{2}, e_{3}$ of the above equation to get

$$
\begin{align*}
w_{1}^{*} & =-w_{1},  \tag{6.3.3}\\
w_{2}^{*} \cos \theta & =w_{2}+\sin \theta w_{21},  \tag{6.3.4}\\
w_{2}^{*} & =w_{31} . \tag{6.3.5}
\end{align*}
$$

df: 1
df:2
df:3

So we have

$$
\begin{equation*}
w_{31} \cos \theta=w_{2}+\sin \theta w_{21} . \tag{6.3.6}
\end{equation*}
$$

Use $d e_{i}^{*} \cdot e_{j}^{*}=w_{i j}^{*}$ to compute
$w_{31}^{*}=d e_{1}^{*} \cdot e_{3}^{*}=-d e_{1} \cdot\left(-\sin \theta e_{2}+\cos \theta e_{3}\right)=\sin \theta w_{21}-\cos \theta w_{31}=-w_{2}$.
Here we use ( ${ }^{\mathrm{dg}}$.3.6). We also have

$$
\begin{align*}
w_{32}^{*} & =d e_{2}^{*} \cdot e_{3}^{*}=\left(\cos \theta d e_{2}+\sin \theta d e_{3}\right) \cdot\left(-\sin \theta e_{2}+\cos \theta e_{3}\right) \\
& =\cos ^{2} \theta w_{32}-\sin ^{2} \theta w_{23}=w_{32} . \tag{6.3.8}
\end{align*}
$$

Let $w_{3 i}=\sum_{j} h_{i j} w_{\dot{e}_{\mathrm{v}}}$ and $K_{i d z} K^{*}$ the Gaussian curvature of $M$ and $M^{*}$


$$
\begin{aligned}
K^{*} w_{1}^{*} \wedge w_{2}^{*} & =w_{31}^{*} \wedge w_{32}^{*}=-w_{2} \wedge w_{32}=h_{12} w_{1} \wedge w_{2} \\
& =-K^{*} w_{1} \wedge w_{31}=-K^{*} h_{12} w_{1} \wedge w_{2} .
\end{aligned}
$$

This implies that $-K^{*} h_{12}=h_{12}$. Hence $K^{*}=-1$. Since the situation is symmetric, $K=-1$ too.

Theorem 6.3.3. Given a surface $M$ in $\mathbb{R}^{3}$, a constant $0<\theta<\pi$, and a unit vector $v_{0} \in T M_{p_{0}}$ not a principal direction, then there exists a unique surface $M^{*}$ and a Bäcklund transformation $\ell: M \rightarrow M^{*}$ with constant $\theta$ such that $\overrightarrow{p_{0} \ell\left(p_{0}\right)}=(\sin \theta) v_{0}$.

Proof. To construct a Bäcklund transformation $\ell$ with constant $\theta$ is the same as to find a unit vector field $e_{1}$ on $M_{d s}$ so that the corresponding frames $w_{i}$ and $\left(w_{i j}\right)$ satisfy the equation ( 6.3 .6 ). We can rewrite this differential form equation as a system of first order PDE as follows. Let $v_{1}, v_{2}, v_{3}$ be a local o.n. frame on $M$ so that $v_{1}, v_{2}$ are principal directions. Let $\left(x_{1}, x_{2}\right)_{\mathrm{q}}$ be the local line of curvature coordinates constructed in Theorem 6.1.1, and $q\left(x_{1}, x_{2}\right)$ the corresponding solution of the SGE ( ${ }_{6}^{62} .1 .3$ ). We know from the proof of Theorem $\frac{\mathrm{cq}}{6.1 .1}$ that

$$
\theta_{1}=\cos q d x_{1}, \quad \theta_{2}=\sin q d x_{2}
$$

are the 1 -forms dual to $v_{1}, v_{2}$, and $\theta_{i j}=d v_{j} \cdot v_{i}=p_{i j} d x_{1}+q_{i j} d x_{2}$ is given as follows:

$$
\begin{aligned}
\theta_{12} & =-q_{x_{2}} d x_{1}-q_{x_{1}} d x_{2}, \\
\theta_{13} & =-\sin q d x_{1}, \\
\theta_{23} & =\cos q d x_{2} .
\end{aligned}
$$

Let $q^{*}$ denote the angle between $v_{1}$ and $e_{1}$, and $e_{1}, e_{2}, e_{3}$ an o.n. frame on $M$ so that $e_{3}=v_{3}$, i.e.,

$$
\left\{\begin{array}{l}
e_{1}=\cos q^{*} v_{1}+\sin q^{*} v_{2},  \tag{6.3.9}\\
e_{2}=-\sin q^{*} v_{1}+\cos q^{*} v_{2}, \\
e_{3}=v_{3} .
\end{array}\right.
$$

Let $w_{i j}=d e_{j} \cdot e_{i}$. Use $\left(\frac{d j_{j}}{6.3 .9)}\right.$ to get

$$
\begin{aligned}
w_{31} & =d e_{1} \cdot e_{3}=\left(\cos q^{*} d v_{1}+\sin q^{*} d v_{2}\right) \cdot v_{3}=\cos q^{*} \theta_{31}+\sin q^{*} \theta_{32} \\
& =\cos q^{*} \sin q d x_{1}-\sin q^{*} \cos q d x_{2}, \\
w_{2} & =-\sin q^{*} \theta_{1}+\cos q^{*} \theta_{2}=-\sin q^{*} \cos q d x_{1}+\cos q^{*} \sin q d x_{2}, \\
w_{21} & =d e_{1} \cdot e_{2}=d q^{*}+\theta_{21}=d q^{*}+q_{x_{2}} d x_{1}+q_{x_{1}} d x_{2} .
\end{aligned}
$$

Equate the coefficients of $d x_{1}$ and $d x_{2}$ of equation $\left(\frac{\mathrm{dg}}{6.3 .6)}\right.$ ) to get

$$
\left\{\begin{array}{l}
-\sin q^{*} \cos q+\sin \theta\left(\left(q^{*}\right)_{x_{1}}+q_{x_{2}}\right)=\cos \theta \cos q^{*} \sin q  \tag{6.3.10}\\
\cos q^{*} \sin q+\sin \theta\left(\left(q^{*}\right)_{x_{2}}+q_{x_{1}}\right)=-\cos \theta \sin q^{*} \cos q
\end{array}\right.
$$

Now we make a change of coordinates:

$$
x_{1}=s+t, \quad x_{2}=s-t
$$

i.e., $(s, t)$ is the Tchebyshef asymptotic coordinate system. Note that

$$
u_{s}=u_{x_{1}}+u_{x_{2}}, \quad u_{t}=u_{x_{1}}-u_{x_{2}} .
$$

Add and substract the two equations of $\left(\frac{\mathrm{dk}}{6.3 .10)}\right.$ ) and use $(s, t)$ coordinates, we see that ( 6.3 .10 ) gives

$$
\left\{\begin{array}{l}
\sin \theta\left(q^{*}+q\right)_{s}=(1-\cos \theta) \sin \left(q^{*}-q\right) \\
\sin \theta\left(q^{*}-q\right)_{t}=(1+\cos \theta) \sin \left(q^{*}+q\right)
\end{array}\right.
$$

So we have

$$
\left\{\begin{array}{l}
\left(q^{*}+q\right)_{s}=\frac{(1-\cos \theta)}{\sin \theta} \sin \left(q^{*}-q\right), \\
\left(q^{*}-q\right)_{t}=\frac{(1+\cos \theta)}{\sin \theta} \sin \left(q^{*}+q\right)
\end{array}\right.
$$

Let $\mu=\frac{(1-\cos \theta)}{\sin \theta}$. Then $\frac{1}{\mu}=\frac{1+\cos \theta}{\sin \theta}$. So the above system becomes

$$
\left\{\begin{array}{l}
\left(q^{*}+q\right)_{s}=\mu \sin \left(q^{*}-q\right),  \tag{6.3.11}\\
\left(q^{*}-q\right)_{t}=\frac{1}{\mu} \sin \left(q^{*}+q\right),
\end{array}\right.
$$

i.e.,

$$
\left\{\begin{array}{l}
\left(q^{*}\right)_{s}=-q_{s}+\mu \sin \left(q^{*}-q\right),  \tag{6.3.12}\\
\left(q^{*}\right)_{t}=q_{t}+\frac{1}{\mu} \sin \left(q^{*}+q\right) .
\end{array}\right.
$$

To find the unit vector field $e_{1}$ so that $\ell=f+\sin \theta e_{1}$ is a Bäcklund transformation is equivalent to solve the differential frorm equation


By the Frobenius Theorem 2.4 .1 , the condition for system $\left(6.3_{d 1}\right)$ to be solvable is to equate the $t$-derivative of the first equation of ( $\sqrt{6.3 .12 \text { ) }}$ and the $s$-derivative of the second equation of ( $\sqrt{6.3 .12})$ :

$$
\begin{aligned}
\left(\left(q^{*}\right)_{s}\right)_{t} & =-q_{s t}+\mu \cos \left(q^{*}-q\right)\left(q^{*}-q\right)_{t} \\
& =-q_{s t}+\cos \left(q^{*}-q\right) \sin \left(q^{*}+q\right), \\
=\left(\left(q^{*}\right)_{t}\right)_{s} & =q_{s t}+\frac{1}{\mu} \cos \left(q^{*}+q\right)\left(q^{*}+q\right)_{s} \\
& =q_{s t}+\cos \left(q^{*}+q\right) \sin \left(q^{*}-q\right) .
\end{aligned}
$$

So the compatibility condition is

$$
2 q_{s t}=\cos \left(q^{*}-q\right) \sin \left(q^{*}+q\right)-\cos \left(q^{*}+q\right) \sin \left(q^{*}-q\right)=\sin (2 q)
$$ i.e., $q$ satisfies the SGE ( (hn $^{\text {d. }} 7$ ).

We use $\mathrm{BT}_{q, \mu}$ to denote the system ( $\frac{\mathrm{di}}{6.3 .11 \text { ) (or }\left(\frac{\mathrm{dl}}{6.3 .12))} \text { ) for given } q . \square . ~\right.}$ and constant $\mu$. The proof of the above theorem gives a result in PDE.
Corollary 6.3.4. Let $\mu$ be a non-zero real constant and $q(s, t) a$ smooth function. Then the ${ }_{\text {dhy }}$ (6.3.11) for $q^{*}$ is solvable if and only if $q$ satisfies the $S_{i} G E$ ( 6.1 .7 ), i.e., $2 q_{s t}=\sin (2 q)$. Moreover, if $q^{*}$ is a solution of ( $\overline{6} .3 .11$ ), then $q^{*}$ is also a solution of the $S G E$ ( $\overline{6.1} .7$ ).

The above Corollary also implies that if $q$ is a solution of the SGE and a real constant $\mu$, then the first order system $\mathrm{BT}_{q, \mu}(6.3 .11)$ is solvable for $q^{*}$ and the solution $q^{*}$ is again a solution of the SGE. So we can solve the system $\mathrm{BT}_{q, \mu}$ of two compatible ODE to get a family of solutions of SGE. If we apply this method again, we get a second family of solutions. This gives an inifintely many families of solutions from one given solution of SGE.

For example, the constant function $q=0$ is a trivial solution of the SGE. The system $\mathrm{BT}_{0, \mu}$ is

$$
\left\{\begin{array}{l}
\alpha_{s}=\mu \sin \alpha \\
\alpha_{t}=\frac{1}{\mu} \sin \alpha
\end{array}\right.
$$

It has explicit solution

$$
\alpha(s, t)=2 \tan ^{-1}\left(e^{\mu s+\frac{1}{\mu} t}\right) .
$$

We can apply Bäcklund transformation $\mathrm{BT}_{\alpha, \mu_{1}}$ to this new solution $\alpha$ to get another family of solutions. Note that this two system of ODEs are not as easy to solve as $\mathrm{BT}_{0 . \mu}$. But instead of solving $\mathrm{BT}_{\alpha, \mu_{1}}$ we
can use the following classical Bianchi's theorem, which says that the second family of solutions can be constructed from the first family of solutions by an explicit algebraic formula:

## eb Theorem 6.3.5. Bianchi Permutability Theorem.

Let $0<\theta_{1}, \theta_{2}<\pi$ be constants so that $\sin \theta_{1}^{2} \neq \sin \theta_{2}^{2}$, and $\ell_{i}: M_{0} \rightarrow M_{i}$ Bäcklund transformations with constant $\theta_{i}$ for $i=1,2$. Then there exists a unique surface $M_{3}$ and Bäcklund transformations $\tilde{\ell}_{1}: M_{2} \rightarrow$ $M_{3}$ and $\tilde{\ell}_{2}: M_{1} \rightarrow M_{3}$ with constant $\theta_{1}, \theta_{2}$ respectively so that $\tilde{\ell}_{1} \circ \ell_{2}=$ $\tilde{\ell}_{2} \circ \ell_{1}$. Moreover, if $q_{i}$ is the solution of the $S G E$ corresponding to $M_{i}$ for $0 \leq i \leq 3$, then

$$
\begin{equation*}
\tan \left(\frac{q_{3}-q_{0}}{2}\right)=\frac{\mu_{1}+\mu_{2}}{\mu_{1}-\mu_{2}} \tan \left(\frac{q_{1}-q_{2}}{2}\right) \tag{6.3.13}
\end{equation*}
$$

where $\mu_{i}=\frac{1-\cos \theta_{i}}{\sin \theta_{i}}$.
The proof of the geometric part of the above theorem is rather long and tidious (cf. []). But to prove that the function $q_{3}$ defined by (6.3.13) solves the system $\mathrm{BT}_{\text {dad }, \mu_{2}}$ and $\mathrm{BT}_{q_{2}, \mu_{1}}$ can be done by a direct computation. Equation (6.3.13) gives a formula of $q_{3}$ in terms of $q_{0}, q_{1}, q_{2}$.

Whena $q_{0}=0$, we have solved $q_{i}=2 \tan ^{-1} e^{\mu_{i} s+\frac{1}{\mu_{i}} t}$ for $i=1,2$. So apply ( $\overline{6} .3 .13$ ) get a second family of solutions:

$$
\begin{equation*}
\tan \frac{q_{12}}{2}=\frac{\mu_{1}+\mu_{2}}{\mu_{1}-\mu_{2}}\left(\frac{e^{\mu_{1} s+\frac{1}{\mu_{1}} t}-e^{\mu_{2} s+\frac{1}{\mu_{2}} t}}{1+e^{\left(\mu_{1}+\mu_{2}\right) s+\left(\frac{1}{\mu_{1}}+\frac{1}{\mu_{2}}\right) t}}\right) \tag{6.3.14}
\end{equation*}
$$

Apply ( $(6.3 .13)$ again, we get a third family of solutions. To be more precise, let $\mu_{1}, \mu_{2}, \mu_{3}$ be real numbers so that $\mu_{1}^{2}, \mu_{2}^{2}, \mu_{3}^{2}$ are distinct, $q_{12}$ the solution defined by ( 6.3 .14 ), and $q_{23}$ the solution given by the same formula ( f .3 .14 ) replacing $\mu_{1}, \mu_{2}$ by $\mu_{2}, \mu_{3}$. Then

$$
\tan \frac{q_{123}}{2}=\frac{\mu_{1}+\mu_{3}}{\mu_{1}-\mu_{3}} \tan \left(\frac{q_{12}-q_{23}}{2}\right) .
$$

We can continue this process to construct an infinitely many families of explicit solutions of the SGE.

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[^0]:    ${ }^{1}$ Research supported in part by NSF Grant NSF DMS-9972172 and NCTS of Taiwan.

