# Example: Laplace Equation Problem 

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We would like to find the steady-state temperature of the first quadrant when we keep the axes at the following temperatures:

$$
\begin{array}{r}
u(x, 0)=1 \text { for } 0<x<1 \\
u(x, 0)=0 \text { for } x>1 \\
u(0, y)=0 \text { for all } y>0
\end{array}
$$

So we need to solve the boundary value problem:

$$
\begin{array}{r}
\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}}=0 \\
u(x, 0)=1 \text { for } 0<x<1 \\
u(x, 0)=0 \text { for } x>1 \\
u(0, y)=0 \text { for all } y>0 \tag{3}
\end{array}
$$

We shall use the Fourier transform. But since we have only half the real line as our domain (for $x$ ), we need to use the sine or cosine Fourier transform. When we apply the cosine or sine Fourier transform to the equation, we want to get a simpler differential equation for $U_{c}=\mathcal{F}_{c}\{u(x, y)\}$ (or $U_{s}=\mathcal{F}_{s}\{u(x, y)\}$ if we are taking the sine transform); where the transform is taken with respect to $x$. To this end, we need to see what the Fourier sine transform of the second derivative of $u$ with respect to $x$ is in terms of $U_{c}$ (or $\left.U_{s}\right)$. But this is easy: we just need to use integration by parts. If $f$ is a function that is absolutely integrable and that converges to zero as $x$ goes to $\infty$ (we assume that $f$ is absolutely integrable so that the integrals we take below exist; in general we assume all our functions are absolutely integrable - recall that absolutely integrable means that the integral of the absolute value of $f$ on the whole interval we are concerned with exists):

$$
\begin{array}{r}
\mathcal{F}_{c}\left\{f^{\prime}\right\}=\int_{0}^{\infty} \cos \alpha x f^{\prime}(x) d x=\left.(f(x) \cdot \cos \alpha x)\right|_{0} ^{\infty}+\alpha  \tag{4}\\
\int_{0}^{\infty} f(x) \sin \alpha x d x \\
=f(0)+\alpha \mathcal{F}_{s}\{f\}
\end{array}
$$

Similarly for the sine transform:

$$
\begin{align*}
& \mathcal{F}_{s}\left\{f^{\prime}\right\}=\int_{0}^{\infty} \sin \alpha x f^{\prime}(x) d x=\left.(f(x) \cdot \sin \alpha x)\right|_{0} ^{\infty}-\alpha \int_{0}^{\infty} f(x) \cos \alpha x d x  \tag{5}\\
& =-\alpha \mathcal{F}_{c}\{f\}
\end{align*}
$$

Now we can combine these two to get:

$$
\begin{array}{r}
\mathcal{F}_{\mathcal{F}}\left\{f^{\prime \prime}\right\}=f^{\prime}(0)-\alpha^{2} \mathcal{F}_{c}\{f\}  \tag{6}\\
\text { and } \mathcal{F}_{s}\left\{f^{\prime \prime}\right\}=-\alpha f(0)-\alpha^{2} \mathcal{F}_{s}\{f\}
\end{array}
$$

If we used the cosine transform, i.e. we applied the cosine transform to the equation, we would get:

$$
\begin{equation*}
\mathcal{F}_{c}\left\{\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}}\right\}=\mathcal{F}_{c}\left\{\frac{\partial^{2} u}{\partial x^{2}}\right\}+\mathcal{F}_{c}\left\{\frac{\partial^{2} u}{\partial y^{2}}\right\}=\frac{\partial u}{\partial x}(0, y)-\alpha^{2} U_{c}+\frac{\partial^{2} U_{c}}{\partial y^{2}}=0 \tag{7}
\end{equation*}
$$

which is not very nice because we don't know what to do with the first derivative of $u$. Whereas:

$$
\begin{equation*}
\mathcal{F}_{s}\left\{\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}}\right\}=\mathcal{F}_{s}\left\{\frac{\partial^{2} u}{\partial x^{2}}\right\}+\mathcal{F}_{s}\left\{\frac{\partial^{2} u}{\partial y^{2}}\right\}=u(0, y)-\alpha^{2} U_{s}+\frac{\partial^{2} U_{s}}{\partial y^{2}}=0 \tag{8}
\end{equation*}
$$

which is much better because one of our assumptions for $u$ in the boundary value problem is that $u(0, y)=0$. So we should prefer the sine transform in this case. (Note that instead of this boundary problem, we looked at the problem with the $y$ axis being isolated, we would prefer the cosine transform). Transforming also the boundary consitions using the sine transform, we have a new boundary value problem for $U=U_{s}$.

$$
\begin{align*}
\frac{\partial^{2} U}{\partial y^{2}}-\alpha^{2} U & =0  \tag{9}\\
U(0, y) & =0  \tag{10}\\
U(\alpha, 0) & =\int_{0}^{\infty} u(x, 0) \sin \alpha x d x=\int_{0}^{1} 50 \sin \alpha x d x=50 \frac{-1}{\alpha} \sin \alpha \tag{11}
\end{align*}
$$

We know the solution to the above differential equation. It must be of the form:

$$
\begin{equation*}
U=c_{1}(\alpha) \cosh \alpha y+c_{2}(\alpha) \sinh \alpha y \tag{12}
\end{equation*}
$$

Plugging in the boundary condition, we see that:

$$
\begin{equation*}
U(\alpha, 0)=c_{1}(\alpha)=-50 \frac{\sin \alpha}{\alpha} \tag{13}
\end{equation*}
$$

We cannot say anything about $c_{2}$ by looking at the boundary conditions. However, we know that the function $U$ must be bounded as $\alpha \rightarrow \infty$. We can argue that this should be true because $U$ must physically make sense. But we can also see this mathematically. Indeed, $U$ is the transform of an absolutely integrable function, so it must be bounded as $\alpha \rightarrow \infty$. Now, recalling the definitions:

$$
\begin{equation*}
\cosh (\alpha y)=\frac{e^{\alpha y}+e^{-\alpha y}}{2} \text { and } \sinh \alpha y=\frac{e^{\alpha y}-e^{-\alpha y}}{2} \tag{14}
\end{equation*}
$$

we see that the only way to have $U$ bounded is to have the $e^{\alpha y}$ 's cancel. So we must have $c_{2}(\alpha)=-c_{1}(\alpha)$. Therefore we have:

$$
\begin{equation*}
U(\alpha, y)=-50 \frac{\sin \alpha}{\alpha} e^{-\alpha y} \tag{15}
\end{equation*}
$$

Finally, taking the inverse Fourier sine transform of $U$ to get $u$, we have:

$$
\begin{equation*}
u(x, y)=-\frac{2}{\pi} \int_{0}^{\infty} 50 \frac{\sin \alpha}{\alpha} e^{-\alpha y} \sin (x \alpha) d \alpha \tag{16}
\end{equation*}
$$

