

## 16.7 Surface Integrals

**Definition.** Suppose a surface  $S$  is parametrized by  $\mathbf{r}(u, v)$  where  $(u, v) \in D \subseteq \mathbb{R}^2$ , and that  $f : S \rightarrow \mathbb{R}$  is continuous on  $S$ . The integral of  $f$  over  $S$  is defined to be

$$\iint_S f \, dS = \iint_D f(\mathbf{r}(u, v)) |\mathbf{r}_u \times \mathbf{r}_v| \, du \, dv$$

This is the same  $dS$  from previously, so any formulae from special cases can be utilized directly.

**Example 1** Compute the integral of  $f(x, y, z) = x^2$  over the part of the plane  $3x + 2y + 6z = 6$  lying in the positive octant.

View  $S$  as the graph of the function  $z = g(x, y) = 1 - \frac{1}{2}x - \frac{1}{3}y$  above the triangle

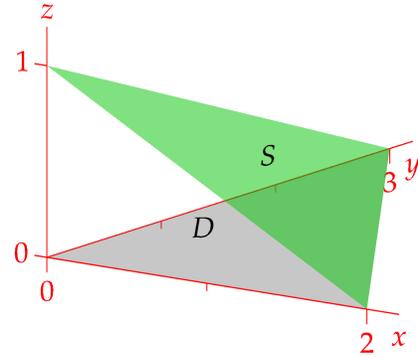
$$D : 0 \leq x \leq 2; \quad 0 \leq y \leq 3 - \frac{3}{2}x$$

The area element using standard rectangular co-ordinates is therefore

$$dS = \sqrt{g_x^2 + g_y^2 + 1} \, dx \, dy = \sqrt{\frac{1}{4} + \frac{1}{9} + 1} \, dx \, dy = \frac{7}{6} \, dx \, dy$$

and the required integral is

$$\iint_S x^2 \, dS = \frac{7}{6} \int_0^2 \int_0^{3-\frac{3}{2}x} x^2 \, dy \, dx = \frac{7}{6} \int_0^2 3x^2 - \frac{3}{2}x^3 \, dx = \frac{7}{3}$$



**Example 2** Calculate the average height of a point on the surface of the cone  $(1 - z)^2 = x^2 + y^2$  lying above the  $x, y$ -plane.

We parametrize the surface using polar co-ordinates  $(r, \theta)$  on the base disk in the  $(x, y)$ -plane:

$$\mathbf{r}(r, \theta) = \begin{pmatrix} r \cos \theta \\ r \sin \theta \\ 1 - r \end{pmatrix}, \quad 0 \leq r \leq 1, \quad 0 \leq \theta < 2\pi$$

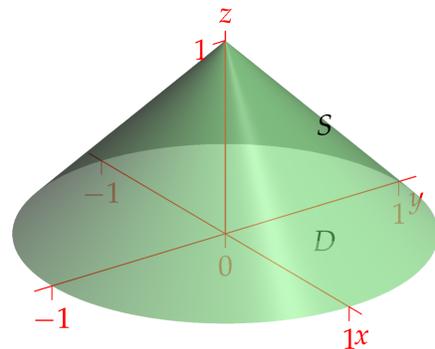
Therefore

$$\mathbf{r}_r \times \mathbf{r}_\theta = \begin{pmatrix} \cos \theta \\ \sin \theta \\ -1 \end{pmatrix} \times \begin{pmatrix} -r \sin \theta \\ r \cos \theta \\ 0 \end{pmatrix} = \begin{pmatrix} r \cos \theta \\ r \sin \theta \\ r \end{pmatrix}$$

$$\Rightarrow dS = \sqrt{2}r \, dr \, d\theta$$

The average height is therefore

$$z_{\text{av}} = \frac{\iint_S z \, dS}{\iint_S dS} = \frac{\int_0^{2\pi} \int_0^1 (1 - r) \sqrt{2}r \, dr \, d\theta}{\int_0^{2\pi} \int_0^1 \sqrt{2}r \, dr \, d\theta} = \frac{1}{3}$$



## Oriented Surfaces and Flux Integrals

To develop integrals of vector fields over surfaces, we need to describe which side of a surface is 'up.'

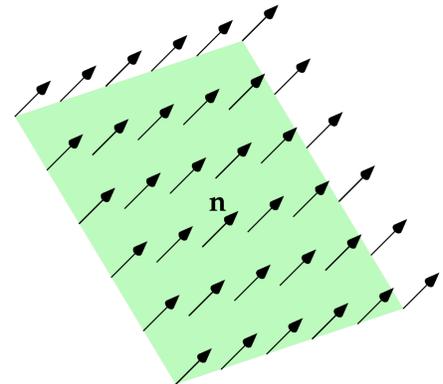
**Definition.** A connected surface  $S$  is *orientable* if a smooth choice of unit normal vector field  $\mathbf{n}$  can be made: this choice is called an *orientation*.

Suppose  $\mathbf{r}(u, v)$  parametrizes  $S$ . We say that  $(u, v)$  are *oriented co-ordinates* if  $\mathbf{r}_u \times \mathbf{r}_v$  always points in the same direction as the orientation  $\mathbf{n}$ .

Loosely speaking, an orientable surface  $S$  has two faces. More precisely, it has exactly two orientations: reverse the normal field to obtain the other ( $\mathbf{n} \mapsto -\mathbf{n}$ ).

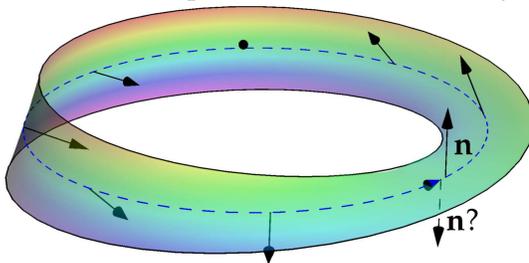
There are several conventions when choosing an orientation.

1. Graphs of functions are typically have  $\mathbf{n}$  pointing upwards. Equivalently,  $(x, y)$  or  $(r, \theta)$  are oriented co-ordinates.
2. The *positive orientation* of a closed surface (such as a sphere) is when  $\mathbf{n}$  points *outwards*.
3. If a surface is piecewise smooth, the orientation must be preserved across edges.



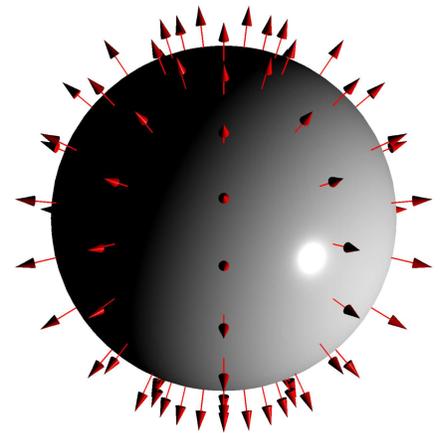
A choice of orientation of a plane

Not all surfaces are orientable. If we travel round the Möbius strip (below), the notion of 'up' reverses: it is impossible to smoothly choose  $\mathbf{n}$  on the entire strip because the surface only has one face!



A Möbius strip (non-orientable)

The Klein Bottle is another famous non-orientable surface.



A positively oriented sphere

**Definition.** If  $\mathbf{F}$  is a continuous vector field defined on an oriented surface  $S$  with unit normal vector field  $\mathbf{n}$ , then the *flux integral* of  $\mathbf{F}$  over  $S$  (or *surface integral*, or *flux* of  $\mathbf{F}$  across  $S$ ) is

$$\iint_S \mathbf{F} \cdot \mathbf{n} \, dS = \iint_S \mathbf{F} \cdot d\mathbf{S}$$

where  $\mathbf{n} \, dS$  is the *vector area element* of the surface. If  $(u, v) \in D$  are oriented co-ordinates, then

$$d\mathbf{S} = \mathbf{n} \, dS = \frac{\mathbf{r}_u \times \mathbf{r}_v}{|\mathbf{r}_u \times \mathbf{r}_v|} |\mathbf{r}_u \times \mathbf{r}_v| \, du \, dv = \mathbf{r}_u \times \mathbf{r}_v \, du \, dv \implies \iint_S \mathbf{F} \cdot d\mathbf{S} = \iint_D \mathbf{F} \cdot (\mathbf{r}_u \times \mathbf{r}_v) \, du \, dv$$

If we reverse the orientation ( $\mathbf{n} \mapsto -\mathbf{n}$ ), the flux changes sign.

You should compare all this with the 2D-flux encountered earlier:  $\int_C \mathbf{F} \cdot \mathbf{n} \, ds = \int_a^b \mathbf{F}(\mathbf{r}(t)) \cdot \mathbf{r}'(t) \, dt$ .

**Example 1** Compare the fluxes of the vector fields  $\mathbf{F}_1 = x^2 \mathbf{i} - y \mathbf{j} + z \mathbf{k}$  and  $\mathbf{F}_2 = x^2 \mathbf{i} + y \mathbf{j} + z \mathbf{k}$  across the upward-oriented plane  $x + y + z = 1$  in the positive octant.

The plane is a graph, so we use rectangular co-ordinates:

$$\mathbf{r}(x, y) = \begin{pmatrix} x \\ y \\ 1-x-y \end{pmatrix} \quad 0 \leq x \leq 1, \quad 0 \leq y \leq 1-x$$

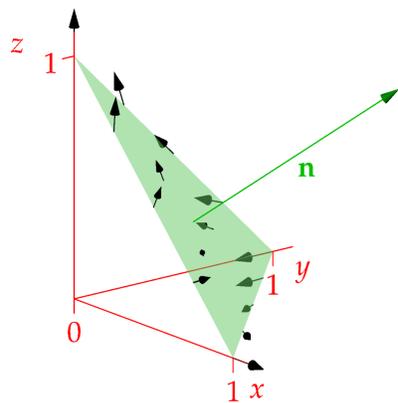
quickly gives the correct normal vector field  $\mathbf{n} = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$  (points upward!). The flux is therefore

$$\begin{aligned} \iint_S \mathbf{F}_1 \cdot d\mathbf{S} &= \int_0^1 \int_0^{1-x} \begin{pmatrix} x^2 \\ -y \\ 1-x-y \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} dy dx \\ &= \int_0^1 \int_0^{1-x} x^2 - x - 2y + 1 dy dx = \frac{1}{12} \end{aligned}$$

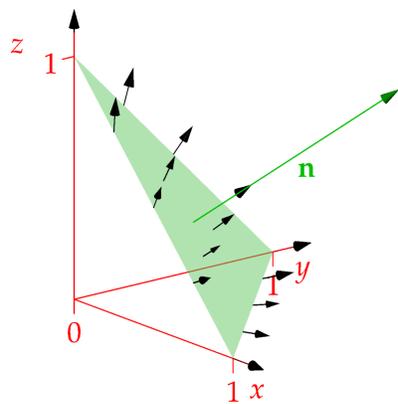
For the second flux,

$$\begin{aligned} \iint_S \mathbf{F}_2 \cdot d\mathbf{S} &= \int_0^1 \int_0^{1-x} \begin{pmatrix} x^2 \\ y \\ 1-x-y \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} dy dx \\ &= \int_0^1 \int_0^{1-x} x^2 - x + 1 dy dx = \frac{5}{12} \end{aligned}$$

Consider how the directions of the arrows help explain the *fivefold* increase in flux. The second field generally has a larger component in the direction of the normal field  $\mathbf{n}$ , so the flux is larger.



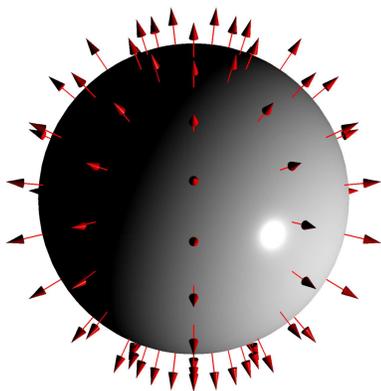
Flux of  $\mathbf{F}_1$  across  $S$



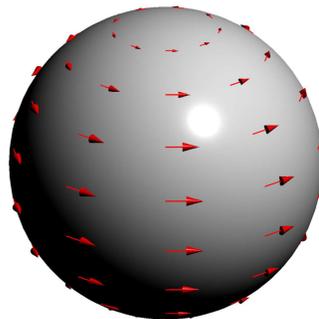
Flux of  $\mathbf{F}_2$  across  $S$

**Example 2** We compare the fluxes of the radial and rotational fields  $\mathbf{F}_1 = \mathbf{r} = \begin{pmatrix} x \\ y \\ z \end{pmatrix}$  and  $\mathbf{F}_2 = \begin{pmatrix} -y \\ x \\ 0 \end{pmatrix}$  across the sphere of radius 1. Here  $\mathbf{n} = \begin{pmatrix} x \\ y \\ z \end{pmatrix} = \mathbf{r}$  (write  $r = |\mathbf{r}|$ ) and so

$$\begin{aligned} \iint_S \mathbf{F}_1 \cdot d\mathbf{S} &= \iint_S \mathbf{r} \cdot \mathbf{r} dS = \iint_S r^2 dS = \iint_S dS = \text{Area}(S) = 4\pi \\ \iint_S \mathbf{F}_2 \cdot d\mathbf{S} &= \iint_S \begin{pmatrix} -y \\ x \\ 0 \end{pmatrix} \cdot \begin{pmatrix} x \\ y \\ z \end{pmatrix} dS = \iint_S dS = 0 \end{aligned}$$



Arrows normal to  $S$ : flux  $> 0$



Arrows tangent to  $S$ : flux  $= 0$

**Example 3: Net Flux** We compute the flux of  $\mathbf{F} = x \mathbf{i} + (5 - z) \mathbf{k}$  out of  $S$ , which comprises two sub-surfaces:

$S_1$ : the paraboloid  $z = 4 - (x^2 + y^2)$  for  $z \geq 0$

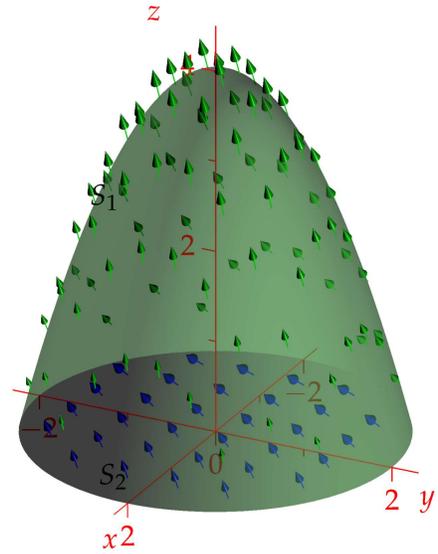
$S_2$ : the elliptical disk  $x^2 + y^2 \leq 4$  in the  $xy$ -plane

Parametrize  $S_1$  using polar co-ordinates on the base:

$$\mathbf{r}_1(r, \theta) = \begin{pmatrix} r \cos \theta \\ r \sin \theta \\ 4 - r^2 \end{pmatrix}, \quad 0 \leq r \leq 2, \quad 0 \leq \theta \leq 2\pi$$

we see that the normal vector field  $\mathbf{n}$  points upwards. Indeed

$$\begin{aligned} \iint_{S_1} \mathbf{F} \cdot d\mathbf{S} &= \int_0^{2\pi} \int_0^2 \begin{pmatrix} r \cos \theta \\ 0 \\ 1+r^2 \end{pmatrix} \cdot \begin{pmatrix} 2r^2 \cos \theta \\ 2r^2 \sin \theta \\ r \end{pmatrix} dr d\theta \\ &= \int_0^{2\pi} \int_0^2 2r^3 \cos^2 \theta + r^3 + r dr d\theta = 20\pi \end{aligned}$$



Over  $S_2$ , the outward pointing normal vector is  $\mathbf{n} = -\mathbf{k}$  (downward!) and so

$$\iint_{S_2} \mathbf{F} \cdot d\mathbf{S} = - \iint_S \begin{pmatrix} x \\ 0 \\ 5-z \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix} dS = -5 \iint_S dS = -20\pi$$

since  $S_2$  is a disk of radius 2. The net flux is therefore zero: there is as much flux into the region bounded by  $S$  as there is leaving it.<sup>1</sup>

### Application: Heat Flow

If  $T(x, y, z)$  measures temperature, the *heat flow* is the vector field  $\mathbf{F} = -k\nabla T$ , where  $k$  (the *thermal conductivity*) is a constant of the material. The *heat flux* out of a region with boundary surface  $S$  is

$$\iint_S \mathbf{F} \cdot d\mathbf{S} = -k \iint_S \nabla T \cdot \mathbf{n} dS$$

(Units:  $[T] = \text{K}$ ,  $[\nabla T] = \text{Km}^{-1}$ ,  $[k] = \text{Wm}^{-1}\text{K}^{-1}$ ,  $[\mathbf{F}] = \text{Wm}^{-2}$ )

**Example** At a given moment, a hot 10cm-radius sphere of gold ( $k = 318 \text{ Wm}^{-1}\text{K}^{-1}$ ) has temperature  $T(r) = \frac{500}{1+r} \text{ K}$ , measured a distance  $r$  meters from its center. The heat flow is therefore

$$\mathbf{F} = -k\nabla T = \frac{159}{r(1+r)^2} \mathbf{r} \quad \text{kWm}^{-2} \quad (\nabla f(r) = \frac{f'(r)}{r} \mathbf{r})$$

Since the unit normal vector field on the surface is  $\frac{1}{r} \mathbf{r}$ , the total heat flux out of the sphere at the given moment is

$$\iint_S \mathbf{F} \cdot \mathbf{n} dS = \iint_{r=0.1} \frac{159}{(1+r)^2} dS = \frac{159}{1.1^2} \iint_{r=0.1} dS = \frac{159}{1.1^2} \cdot 4\pi(0.1)^2 = \frac{636\pi}{121} \approx 16.5 \text{ kW}$$

<sup>1</sup>As we'll see later, the cheap reason for this is because the divergence of  $\mathbf{F}$  is zero...