

## 16.7 Surface Integrals

Recall: if  $D$  is a 2D-region, then  $\iint_D f \, dA = \text{Area}(D) \cdot f_{\text{av}}$

For surfaces: approximate  $S$  by parallelograms  $\Delta S_{ij}$  located at points  $P_{ij}$ . On each parallelogram,  $f_{\text{av}} \approx f(P_{ij})$ , and so it seems sensible to have

$$\iint_S f \, dS \approx \sum_{i,j} f(P_{ij}) \cdot \text{Area}(\Delta S_{ij})$$

where the approximation improves with more parallelograms. If  $S$  is parameterized then we know how to compute the areas:

**Theorem.** Suppose the surface  $S$  is parameterized by  $\mathbf{r}(u, v)$  where  $(u, v) \in D \subseteq \mathbb{R}^2$ , and that  $f$  is continuous on  $S$ . The integral of  $f$  over  $S$  is defined to be

$$\iint_S f \, dS = \iint_D f(\mathbf{r}(u, v)) |\mathbf{r}_u \times \mathbf{r}_v| \, du \, dv$$

**Similar Notations** Observe how the notation is similar to that for a line integral:

$$\text{Surface Integral: } \iint_S f \, dS = \iint_D f(\mathbf{r}(u, v)) |\mathbf{r}_u \times \mathbf{r}_v| \, du \, dv$$

$$\text{Line Integral: } \int_C f \, ds = \int_a^b f(\mathbf{r}(t)) |\mathbf{r}'(t)| \, dt$$

**Example** Find the integral of the function  $f(x, y, z) = x^2$  over the part of the plane  $x + y + z = 2$  above the triangle with corners  $(0, 0)$ ,  $(1, 0)$ ,  $(0, 1)$

The surface  $S$  is the graph of

$$z = g(x, y) = 2 - x - y$$

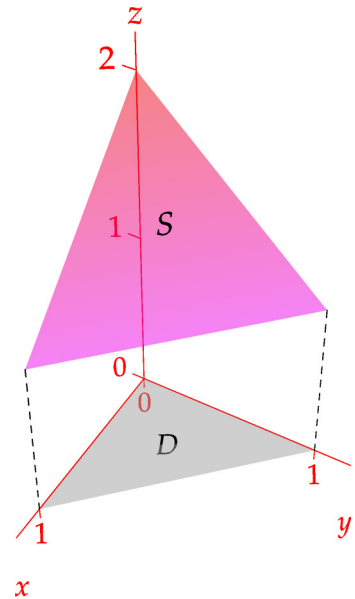
hence the area element is

$$dS = \sqrt{g_x^2 + g_y^2 + 1} \, dx \, dy = \sqrt{3} \, dx \, dy$$

The integral is therefore

$$\begin{aligned} \iint_S x^2 \, dS &= \iint_D x^2 \sqrt{3} \, dx \, dy = \int_0^1 \int_0^{1-x} x^2 \sqrt{3} \, dy \, dx \\ &= \sqrt{3} \int_0^1 x^2 - x^3 \, dx = \frac{\sqrt{3}}{12} \end{aligned}$$

This example is easy because  $S$  is a *graph* and we can therefore see the domain  $D$  as part of the  $xy$ -plane. The parameterization  $\mathbf{r}(u, v) = \begin{pmatrix} u \\ v \\ 2-u-v \end{pmatrix}$  and the computation of  $dS = |\mathbf{r}_u \times \mathbf{r}_v| \, du \, dv$  does not need to be made explicit.



**Example** Calculate  $\iint_S z \, dS$  where  $S$  is the region of the cone  $z = 1 - r$  in cylindrical polar co-ordinates lying above the  $x, y$ -plane

This is also a graph, although it is easiest to parameterize the cone using polar co-ordinates  $(u, v) = (r, \theta)$ :

$$\mathbf{r}(u, v) = \begin{pmatrix} u \cos v \\ u \sin v \\ 1 - u \end{pmatrix}, \quad 0 \leq u \leq 1, \quad 0 \leq v < 2\pi$$

Therefore

$$\begin{aligned} \mathbf{r}_u \times \mathbf{r}_v &= \begin{pmatrix} \cos v \\ \sin v \\ -1 \end{pmatrix} \times \begin{pmatrix} -u \sin v \\ u \cos v \\ 0 \end{pmatrix} = \begin{pmatrix} u \cos v \\ u \sin v \\ u \end{pmatrix} \\ \implies dS &= \sqrt{2} u \, du \, dv \end{aligned}$$

The integral is therefore

$$\iint_S z \, dS = \iint_D (1 - u) \sqrt{2} u \, du \, dv = \sqrt{2} \int_0^{2\pi} \int_0^1 u - u^2 \, du \, dv = \frac{\sqrt{2}\pi}{3}$$

**Piecewise Surfaces** Surfaces can consist of multiple parts: we must compute a separate surface integral, with separate parameterization, for each piece.

For example, we compute the integral  $\iint_S y^2 - z \, dS$  where  $S$  comprises parts of the sphere  $x^2 + y^2 + z^2 = 4$  and the cylinder  $x^2 + y^2 = 1$  as shown. The intersections happen at  $z = \pm\sqrt{3}$

Take  $S_1$  to be the upper part of the sphere,  $S_2$  the lower, and  $S_3$  the cylinder

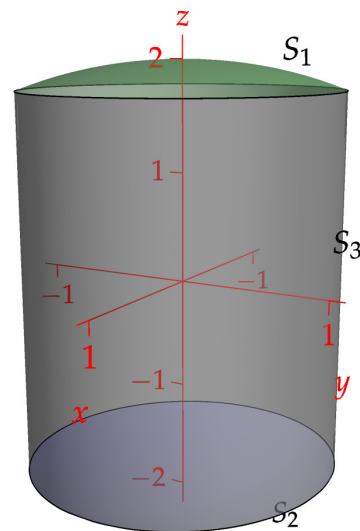
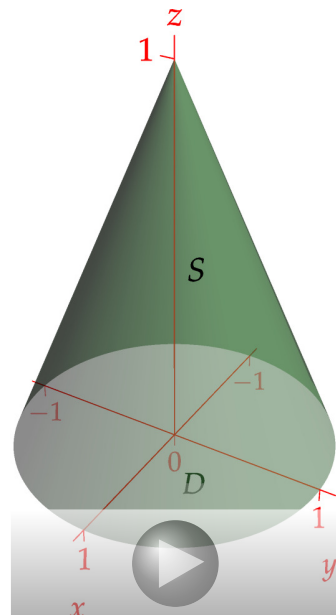
There are several choices for how to parameterize each surface. We will use spherical co-ordinates for the parts of the spheres, and cylindrical co-ordinates for the cylinder.

It is straightforward to check that the maximum angle  $\phi$  for the upper cap  $S_1$  is  $\phi = \tan^{-1} \frac{1}{\sqrt{3}} = \frac{\pi}{6}$ . We therefore have the parameterizations:

$$\mathbf{r}_1(u, v) = \begin{pmatrix} 2 \sin u \cos v \\ 2 \sin u \sin v \\ 2 \cos u \end{pmatrix} \quad 0 \leq u \leq \frac{\pi}{6}, \quad 0 \leq v < 2\pi \quad ((u, v) = (\phi, \theta))$$

$$\mathbf{r}_2(u, v) = \mathbf{r}_1(u, v) \quad \frac{5\pi}{6} \leq u \leq \pi, \quad 0 \leq v < 2\pi$$

$$\mathbf{r}_3(u, v) = \begin{pmatrix} \cos u \\ \sin u \\ v \end{pmatrix} \quad 0 \leq u < 2\pi, \quad -\sqrt{3} \leq v \leq \sqrt{3} \quad ((u, v) = (\theta, z))$$



Calculate the area elements:

$$\begin{aligned} dS_1 = dS_2 &= \left| \begin{pmatrix} 2 \cos u \cos v \\ 2 \cos u \sin v \\ -2 \sin u \end{pmatrix} \times \begin{pmatrix} -2 \sin u \sin v \\ 2 \sin u \cos v \\ 0 \end{pmatrix} \right| du dv = 4 \sin u \left| \begin{pmatrix} \sin u \cos v \\ \sin u \sin v \\ \cos u \end{pmatrix} \right| du dv \\ &= 4 \sin u du dv \end{aligned}$$

$$dS_3 = \left| \begin{pmatrix} -\sin u \\ \cos u \\ 0 \end{pmatrix} \times \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \right| du dv = \left| \begin{pmatrix} \cos u \\ \sin u \\ 0 \end{pmatrix} \right| du dv = du dv$$

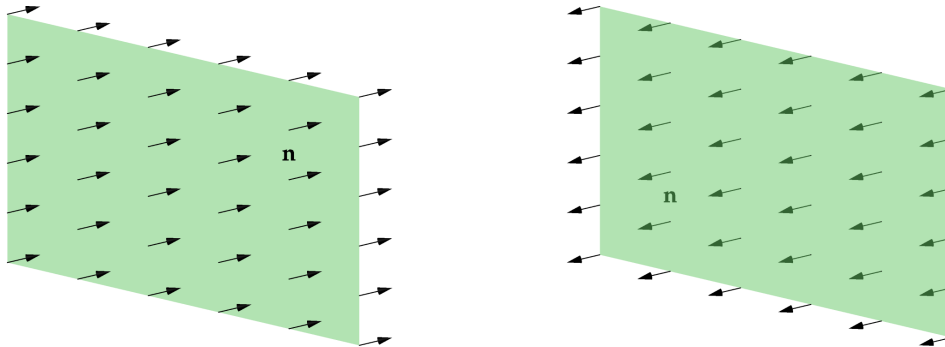
To compute  $\iint_S z^2 dS$  we sum the three integrals:

$$\begin{aligned} \iint_S z^2 dS &= \iint_{S_1} z^2 dS_1 + \iint_{S_2} z^2 dS_2 + \iint_{S_3} z^2 dS_2 \\ &= \left( \int_0^{2\pi} \int_0^{\frac{\pi}{6}} + \int_0^{2\pi} \int_{\frac{5\pi}{6}}^{\pi} \right) (2 \cos u)^2 \cdot 4 \sin u du dv + \int_{-\sqrt{3}}^{\sqrt{3}} \int_0^{2\pi} v^2 du dv \\ &= 2\pi \left( \int_0^{\frac{\pi}{6}} + \int_{\frac{5\pi}{6}}^{\pi} \right) 16 \cos^2 u \sin u du + 2\pi \int_{-\sqrt{3}}^{\sqrt{3}} v^2 dv \\ &= \frac{32\pi}{3} [-\cos^3 u]_0^{\frac{\pi}{6}} + \frac{32\pi}{3} [-\cos^3 u]_{\frac{5\pi}{6}}^{\pi} + \frac{4\pi}{3} v^3 \Big|_0^{\sqrt{3}} \\ &= \frac{32\pi}{3} \left( -\left(\frac{\sqrt{3}}{2}\right)^3 + 1 - (-1) + \left(-\frac{\sqrt{3}}{2}\right)^3 \right) + 4\pi\sqrt{3} = \frac{4(16 - 3\sqrt{3})}{3} \pi \end{aligned}$$

## Orientation

An inhabitant of a surface needs to know which way is 'up'.

**Definition.** A connected surface  $S$  is orientable if it has exactly two faces. An orientation is a smooth choice of unit normal vector field<sup>1</sup>  $\mathbf{n}$ .

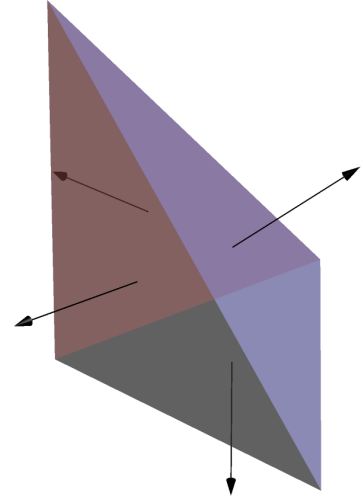


Suppose that  $\mathbf{r}(u, v)$  parameterizes  $S$ . We say that  $(u, v)$  are oriented co-ordinates if  $\mathbf{r}_u \times \mathbf{r}_v$  always points in the same direction as the orientation  $\mathbf{n}$ .

<sup>1</sup>At each point we have a choice of two unit normal vectors  $\mathbf{n}$ . An orientable surface therefore has two opposite possible orientations.

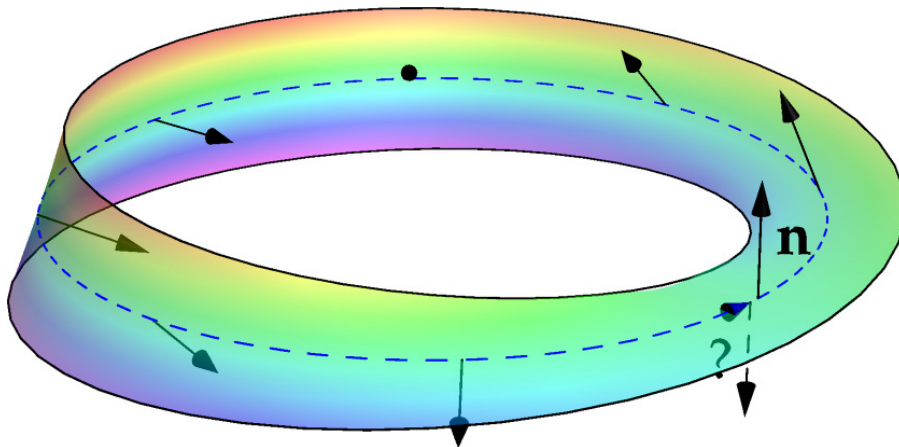
**Choice of Orientation** For most orientable surfaces the choice is entirely arbitrary, although there are several conventions.

1. Graphs of functions  $z = f(x, y)$  or  $z = f(r, \theta)$  are typically orientated so that the normal vectors point upwards. This is equivalent to having  $(x, y)$  and  $(r, \theta)$  be oriented co-ordinates.
2. Closed surfaces (such as spheres) tend to be oriented so that  $\mathbf{n}$  points outwards. We refer to this as *positive orientation* of a closed surface.
3. If a surface is piecewise smooth, the orientation *must* be preserved across edges.



A positively oriented tetrahedron is drawn. Notice that the orientation is preserved across edges, so that all faces are oriented outwards.

**A non-orientable surface** Not all surfaces are orientable, the Möbius strip being the classic example.<sup>2</sup>



Traveling round the dashed curve changes the notion of ‘up’: it is impossible to smoothly choose  $\mathbf{n}$  on the entire surface because the surface only has one face. Small parts of a Möbius strip can be oriented, but not the entire surface.

<sup>2</sup>The Klein Bottle is even weirder...

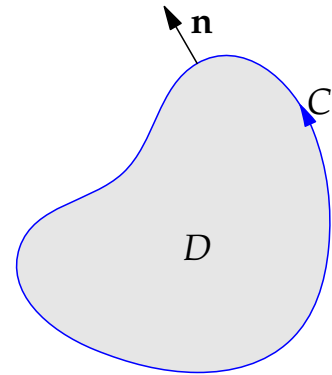
## Flux Integrals

Recall an earlier physical interpretation of Green's Theorem.

If  $C$  is the boundary of a simply-connected region  $D$ , then  $\mathbf{F} \cdot \mathbf{n}$  is the component of  $\mathbf{F}$  pointing out of  $D$ : indeed

$$\oint_C \mathbf{F} \cdot \mathbf{n} \, ds = \text{total flow of } \mathbf{F} \text{ out of } D$$

This notion can be extended to 3-dimensions.



**Definition.** If  $\mathbf{F}$  is a continuous vector field defined on an oriented surface  $S$  with normal vector field  $\mathbf{n}$ , then the flux integral<sup>3</sup> of  $\mathbf{F}$  over  $S$  is defined as

$$\iint_S \mathbf{F} \cdot \mathbf{n} \, dS$$

**Notation** If  $(u, v) \in D$  are oriented co-ordinates on  $S$ , then  $\mathbf{r}_u \times \mathbf{r}_v$  points in the same direction as  $\mathbf{n}$ , whence

$$\mathbf{n} \, dS = \frac{\mathbf{r}_u \times \mathbf{r}_v}{|\mathbf{r}_u \times \mathbf{r}_v|} |\mathbf{r}_u \times \mathbf{r}_v| \, du \, dv = \mathbf{r}_u \times \mathbf{r}_v \, du \, dv$$

Because of this, it is conventional to define the *vector area element*<sup>4</sup>

$$d\mathbf{S} = \mathbf{n} \, dS = \mathbf{r}_u \times \mathbf{r}_v \, du \, dv$$

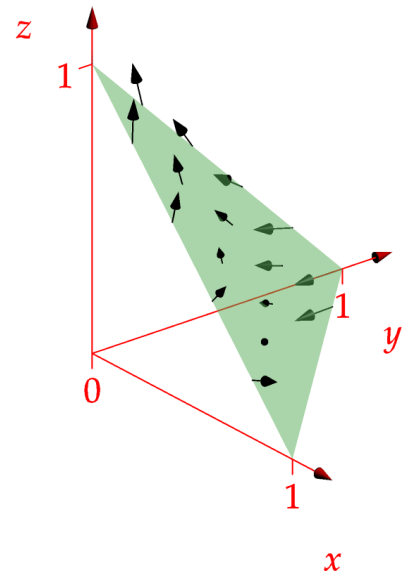
and write the flux integral of  $\mathbf{F}$  as

$$\iint_S \mathbf{F} \cdot d\mathbf{S} = \iint_S \mathbf{F} \cdot \mathbf{n} \, dS = \iint_D \mathbf{F} \cdot (\mathbf{r}_u \times \mathbf{r}_v) \, du \, dv$$

**Example** Compute the flux of the vector field  $\mathbf{F} = x^2\mathbf{i} - y\mathbf{j} + z\mathbf{k}$  across the upward-oriented plane  $x + y + z = 1$  in the positive octant.

The plane is a graph, whence we may use co-ordinates  $(u, v) = (x, y)$ . The parameterization by  $\mathbf{r}(u, v) = \begin{pmatrix} u \\ v \\ 1-u-v \end{pmatrix}$  where  $0 \leq u \leq 1$  and  $0 \leq v \leq 1-u$ , quickly gives the correct normal vector field  $\mathbf{n} = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$ . The flux is then

$$\begin{aligned} \iint_S \mathbf{F} \cdot d\mathbf{S} &= \int_0^1 \int_0^{1-u} \begin{pmatrix} u^2 \\ -v \\ 1-u-v \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} \, dv \, du \\ &= \int_0^1 \int_0^{1-u} u^2 - u - 2v + 1 \, dv \, du = \frac{1}{12} \end{aligned}$$



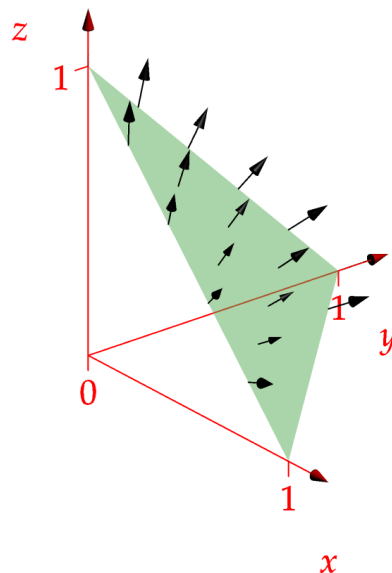
<sup>3</sup>Or *surface integral* or simply *flux*

<sup>4</sup>The distinction between  $ds = |\mathbf{r}'(t)| \, dt$  and  $d\mathbf{r} = \mathbf{r}'(t) \, dt$  is similar. The advantage of  $d\mathbf{S}$  and  $d\mathbf{r}$  are that they require no square-roots to compute.

Now we compute the same flux as before, but with a subtle change in the vector field  $\mathbf{F} = x^2\mathbf{i} + y\mathbf{j} + z\mathbf{k}$ .

$$\begin{aligned} \iint_S \mathbf{F} \cdot d\mathbf{S} &= \int_0^1 \int_0^{1-u} \begin{pmatrix} u^2 \\ v \\ 1-u-v \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} dv du \\ &= \int_0^1 \int_0^{1-u} u^2 - u + 1 dv du = \frac{5}{12} \end{aligned}$$

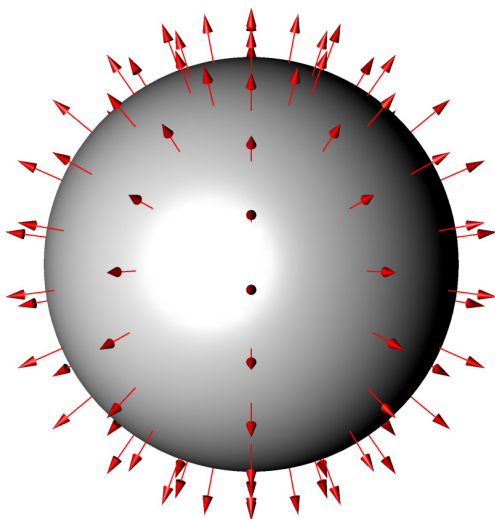
Note how the directions of the arrows explain the *fivefold* increase in flux. The vector field  $\mathbf{F}$  is mostly pointing in the same direction as the normal vector  $\mathbf{n}$ , hence a positive flux.



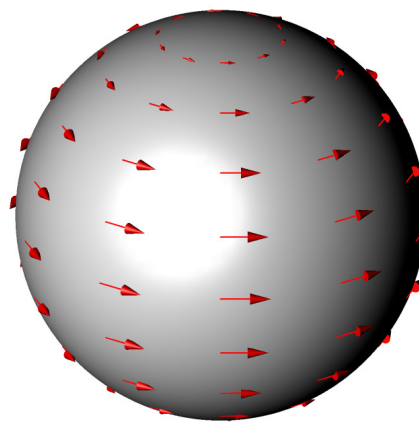
**Example** Compare the flux of the radial and rotational vector fields  $\mathbf{F}_1 = \mathbf{r} = \begin{pmatrix} x \\ y \\ z \end{pmatrix}$ ,  $\mathbf{F}_2 = \begin{pmatrix} -y \\ x \\ 0 \end{pmatrix}$  across the sphere of radius 1.

Here  $\mathbf{n} = \begin{pmatrix} x \\ y \\ z \end{pmatrix} = \mathbf{r}$  and so

$$\begin{aligned} \iint_S \mathbf{F}_1 \cdot d\mathbf{S} &= \iint_S \mathbf{r} \cdot \mathbf{r} dS = \iint_S r^2 dS = \iint_S dS = \text{Area}(S) = 4\pi \\ \iint_S \mathbf{F}_2 \cdot d\mathbf{S} &= \iint_S \begin{pmatrix} -y \\ x \\ 0 \end{pmatrix} \cdot \mathbf{r} dS = 0 \end{aligned}$$



Arrows normal to  $S$ : flux  $> 0$



Arrows tangent: flux  $= 0$

**Net Flux** Flux can be both in and out of a region

For example, let  $S$  comprise two subsurfaces,

$S_1$ : the **paraboloid**  $z = 4 - x^2 - 4y^2$  for  $z \geq 0$

$S_2$  the **elliptical disk**  $x^2 + 4y^2 \leq 4$

Let  $\mathbf{F} = (z - x)\mathbf{i} - \mathbf{j} + (z + 3)\mathbf{k}$

Parameterizing the paraboloid  $S_1$  as a graph  $(u, v) = (x, y)$  oriented upwards, we quickly compute the normal vector field  $\mathbf{n} = \mathbf{r}_u \times \mathbf{r}_v = \begin{pmatrix} 2u \\ 8v \\ 1 \end{pmatrix}$ , whence

$$\begin{aligned} \iint_{S_1} \mathbf{F} \cdot d\mathbf{S} &= \iint_D \begin{pmatrix} (4 - u^2 - 4v^2) - u \\ -1 \\ 4 - u^2 - 4v^2 + 3 \end{pmatrix} \cdot \begin{pmatrix} 2u \\ 8v \\ 1 \end{pmatrix} du dv \\ &= \iint_D (2u + 1)(4 - u^2 - 4v^2) - 2u^2 - 8v + 3 du dv \end{aligned}$$

where  $D$  is the inside of the ellipse  $u^2 + 4v^2 = 4$ .

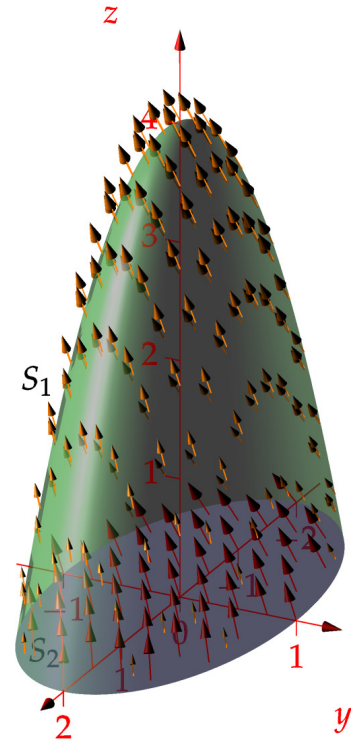
To compute this integral, change variables to modified polar co-ordinates  $u = 2r \cos \theta$ ,  $v = r \sin \theta$ , so that  $du dv = 2r dr d\theta$ . With a little substitution and cancelling, we obtain

$$\begin{aligned} \iint_{S_1} \mathbf{F} \cdot d\mathbf{S} &= \int_0^{2\pi} \int_0^1 ((4r \cos \theta + 1)(4 - 4r^2) - 8r^2 \cos^2 \theta - 8r \sin \theta + 3) 2r dr d\theta \\ &= \int_0^{2\pi} \int_0^1 ((4 - 4r^2 - 4r^2(1 + \cos 2\theta) + 3) 2r dr d\theta \\ &= 4\pi \int_0^1 (7 - 8r^2) r dr = 6\pi \end{aligned}$$

Over  $S_2$ , the outward pointing normal vector is  $\mathbf{n} = -\mathbf{k}$  (*downward!*) and so

$$\iint_{S_2} \mathbf{F} \cdot d\mathbf{S} = - \iint_S \begin{pmatrix} -x \\ -1 \\ 3 \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix} dS = -3 \iint_S dS = -6\pi$$

since  $S_2$  is an ellipse. The net flux is therefore zero: there is as much flux into the region bounded by  $S$  as there is leaving it.



### Non-examinable Application: Heat Flow

If the temperature at a point  $(x, y, z)$  in a solid body is  $T(x, y, z)$ , then its gradient vector field records the direction and rate of greatest increase in the temperature.

The *heat flow* is the vector field  $\mathbf{F} = -k\nabla T$ , where  $k$  (the *conductivity*) is a constant depending on the material.

The *heat flux* (= power output) out of a solid with boundary surface  $S$  is the flux of the heat flow:

$$\iint_S \mathbf{F} \cdot d\mathbf{S} = -k \iint_S \nabla T \cdot \mathbf{n} \, dS$$

(Units:  $[T] = \text{K}$ ,  $[\nabla T] = \text{K m}^{-1}$ ,  $[k] = \text{W m}^{-1} \text{K}^{-1}$ ,  $[\mathbf{F}] = \text{W m}^{-2}$ )

**Example** A hot sphere of gold ( $k = 318 \text{ W m}^{-1} \text{K}^{-1}$ ) of radius 10 cm has temperature

$$T(r) = \frac{1000}{1 + 10r} \text{ K}$$

at a distance of  $r = |\mathbf{r}|$  m from its center. The heat flow is then<sup>5</sup>

$$\mathbf{F} = \frac{3180000}{r(1 + 10r)^2} \mathbf{r} \text{ W m}^{-2}$$

On the surface of the sphere ( $r = 0.1$  m) this is  $\mathbf{F} = 7950000\mathbf{r} \text{ W m}^{-2}$ . Since the unit normal vector field is  $10\mathbf{r}$ , we find that the total heat flux of the gold sphere is

$$\iint_S \mathbf{F} \cdot \mathbf{n} \, dS = \iint_{r=0.1} 7950000\mathbf{r} \cdot 10\mathbf{r} \, dS = 795000 \iint_{r=0.1} dS = 795000 \cdot \frac{4\pi}{10^2} = 31800\pi \text{ W}$$

or approximately 100 kW.

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<sup>5</sup>Using  $\nabla f(r) = \frac{f'(r)}{r} \mathbf{r}$  for any radial function  $f$ .